





The GS09 Double Parton Distribution Functions

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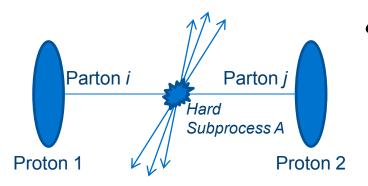
IOP 2010, 29-31 March 2010, University College London

Work performed in collaboration with W.J. Stirling (arXiv:0910.4347), and C.H. Kom,

A. Kulesza and W.J. Stirling (arXiv:1003.3953).

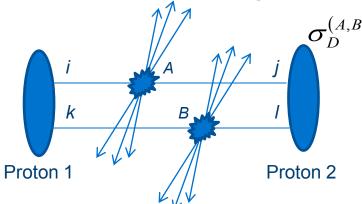
Single vs. Double Parton Scattering

Single Parton Scattering (SPS):



$$\sigma_{S}^{(A)} = \sum_{i,j} D_{h}^{i}(x_{1}; Q_{A}) D_{h}^{j}(x'_{1}; Q_{A}) \hat{\sigma}_{ij}^{A}(x_{1}, x'_{1}) \mathrm{d}x_{1} \mathrm{d}x'_{1}$$
(Single) Parton
$$Distribution Functions$$
Hard subprocess
$$Cross section$$

Double Parton Scattering (DPS):



Symmetry factor

Double Parton Distribution
Functions (dPDFs)

$$\sigma_D^{(A,B)} = \frac{m}{2\sigma_{eff}} \sum_{i,j,k,l} \int D_h^{ik}(x_1,x_2;Q_A,Q_B) D_h^{jl}(x_1',x_2';Q_A,Q_B) \times \hat{\sigma}_{ij}^A(x_1,x_1') \hat{\sigma}_{kl}^B(x_2,x_2') \mathrm{d}x_1 \mathrm{d}x_1' \mathrm{d}x_2 \mathrm{d}x_2'$$
Factor related to correlations of partons in transverse space

Why should we care about double parton scattering at the LHC?

Crudest approximation for dPDFs:

$$D_h^{ij}(x_1, x_2; Q_A, Q_B) \approx D_h^i(x_1; Q_A) D_h^j(x_2; Q_B)$$

$$\Rightarrow \sigma_D^{(A,B)} \approx \frac{m}{2} \frac{\sigma_S^{(A)} \sigma_S^{(B)}}{\sigma_{eff}}$$

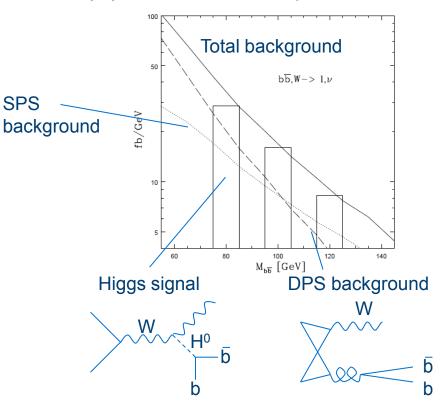
⇒DPS cross sections go like the product of SPS ones!

 \Rightarrow DPS cross sections **grow faster** with energy than SPS σ .

DPS processes...

- provide significant backgrounds to Higgs and new physics signals.
- reveal information about the structure of the proton.

DPS background to Higgs + W production (Del Fabbro and Treleani, hep-ph/9911358,1999):

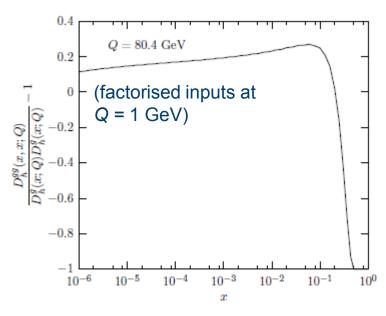


What information do we have on the dPDFs?

Experimental: CDF/D0 investigations of DPS. Their findings are consistent with the factorised approximation for dPDFs (but low statistics & looking at sea quarks at low *x* only).

Theoretical: 'double DGLAP equation' describing the change in the dPDFs with factorisation scale, for the dPDFs with $Q_A = Q_B = Q$. (Kirschner, Phys.Lett.B84:266, 1979 and Shelest, Snigirev, and Zinovjev, Phys.Lett.B113:325,1982).

Crucial prediction of this equation: **pQCD** evolution causes dPDFs to deviate from factorised forms!



JG and Stirling, 0910.4347, 2009



Input dPDFs

The most accurate approach to modelling the (equal scale) dPDFs is to use the double DGLAP equation along with some suitably chosen inputs at a low scale Q_0 .

But what should the inputs look like? Can we get any theoretical insight?

First reaction - **NO**! A dPDF at any particular scale receives contributions from **non-perturbative physics**.



The dPDF Sum Rules

Actually – **YES**, we can! We have shown that the following equalities (**sum rule equalities**) are preserved by double DGLAP:

$$\sum_{j_1} \int_0^{1-x_2} dx_1 x_1 D_h^{j_1 j_2}(x_1, x_2; t) = (1-x_2) D_h^{j_2}(x_2; t)$$

$$\int_0^{1-x_2} dx_1 D_h^{j_1 v j_2}(x_1, x_2; t) = \begin{cases} N_{j_1 v} D_h^{j_2}(x_2; t) & \text{when } j_2 \neq j_1 \text{ or } \overline{j}_1 \\ (N_{j_1 v} - 1) D_h^{j_2}(x_2; t) & \text{when } j_2 = j_1 \\ (N_{j_1 v} + 1) D_h^{j_2}(x_2; t) & \text{when } j_2 = \overline{j}_1 \end{cases}$$

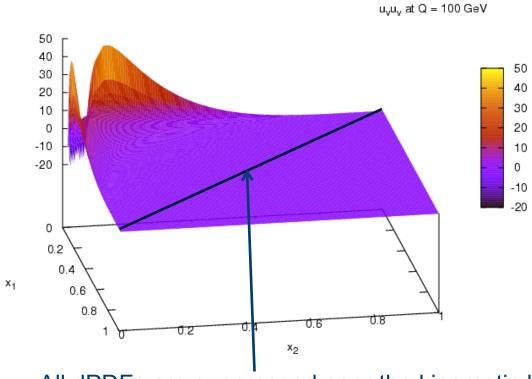
These equalities are no more than the statements of conservation of momentum and quark number for the dPDFs, and have an interpretation in terms of conditional probabilities.

The sum rules impose important constraints on the type of input dPDFs that are allowable.

First set of publicly available LO equal-scale dPDFs (available from **HepForge**). Grid of dPDF values obtained by applying numerical double DGLAP evolution to certain inputs.

Inputs used at Q_0 = 1 GeV are based on products of MSTW2008LO single PDFs, but contain a number of **key features** to ensure that they approximately satisfy the sum rules.

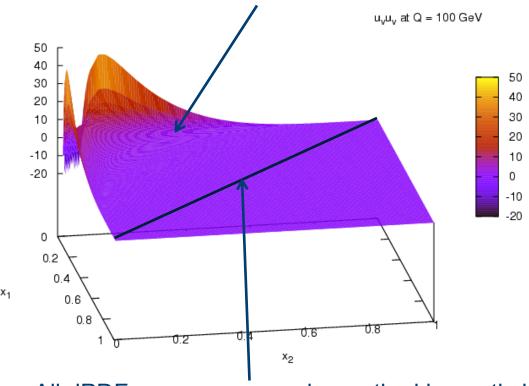




All dPDFs are suppressed near the kinematic boundary x_1+x_2 = 1 to take account of **phase space considerations**.



Terms have been added/subtracted from dPDFs to take account of **number effects**.



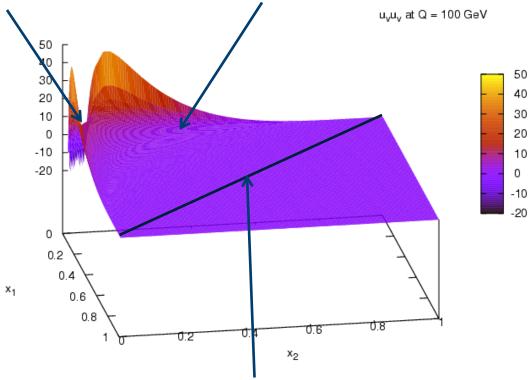
All dPDFs are suppressed near the kinematic boundary x_1+x_2 = 1 to take account of **phase space considerations**.



Terms have been added to the $j\bar{j}$ distributions to take account of $j\bar{j}$

correlations.

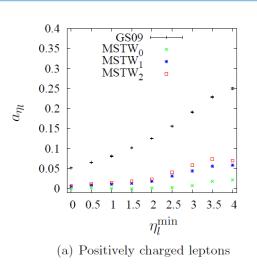
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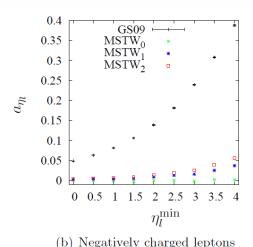


All dPDFs are suppressed near the kinematic boundary x_1+x_2 = 1 to take account of **phase space considerations**.



Comparison of GS09 with factorised dPDFs





JG, Kom, Kulesza, Stirling, 1003.3953, 2010

Comparison in the context of a particular process – equal sign W pair production.

$$a_{\eta_l} = \frac{\sigma(\eta_{l_1} \times \eta_{l_2} < 0) - \sigma(\eta_{l_1} \times \eta_{l_2} > 0)}{\sigma(\eta_{l_1} \times \eta_{l_2} < 0) + \sigma(\eta_{l_1} \times \eta_{l_2} > 0)}.$$

 $a_{\eta l}$ larger for GS09 due to number effect subtractions, especially for large η_l^{min} (i.e. large x, where number effect subtractions have the largest impact).



Future Work

Extend treatment to NLO!

- Need to compute 1→2 splitting functions at NLO (trivial at LO).
- Will need NLO coefficient functions for certain benchmark processes (e.g. equal sign WW production).



Summary

- Important to understand DPS will produce significant backgrounds and interesting signals at the LHC.
- For DPS predictions, require dPDFs. A 'double DGLAP' equation exists dictating the evolution of the equal-scale dPDFs, and we have derived the number and momentum sum rules for these quantities.
- We have produced the first publicly available set of LO equal-scale dDPFs. Sum rules used to guide construction of inputs at $Q_0 = 1$ GeV, and double DGLAP equation used to obtain dPDF values at other scales.
- Number and momentum correlations in GS09 dPDFs affect the signatures of DPS processes – but may be difficult to see this at LHC due to SPS background.

