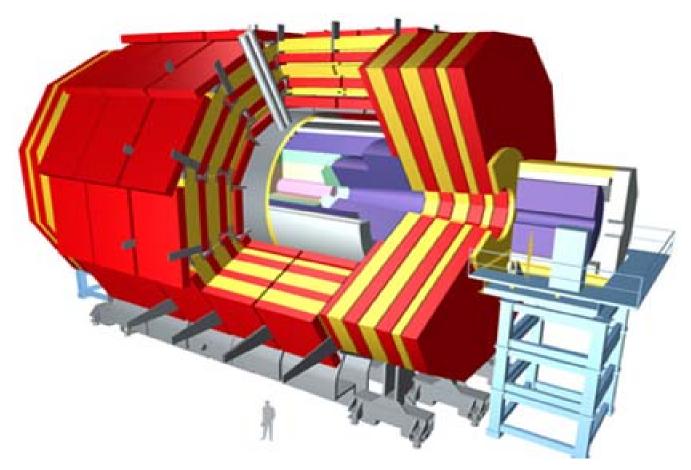
Early searches for supersymmetry at the LHC in the all-hadronic channel

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The CMS detector

Introduction



How do we look for supersymmetry at the LHC?

• What are the challenges of the all-hadronic channel?

How can we guard against mismeasurement?

• Can we use kinematics to constrain fake missing E_T ?

How do we account for Standard Model backgrounds?

• What data-driven tools exist to estimate real missing E_T ?

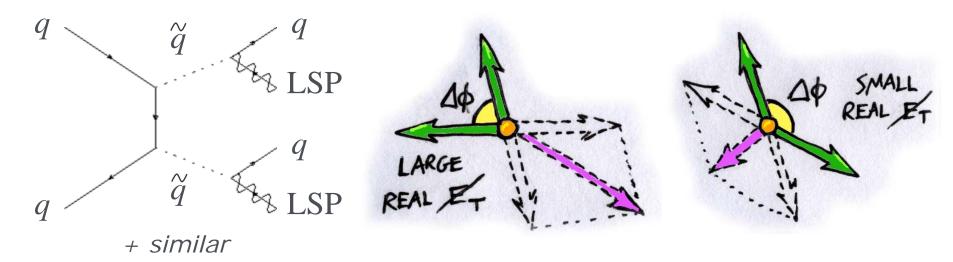


Supersymmetry at the LHC

Supersymmetry is an extension to the Standard Model (SM) that predicts massive, undetectable superpartners to SM particles. These may be produced in LHC proton-proton collisions.

Typical experimental signature:

• Large missing transverse energy plus final state objects.





The all-hadronic channel

We consider events with only jets in the final state:

- "Jet": clustered energy deposits in the calorimeters from the hadronisation of partons.
- Tracking information may also be used for identification and/or correction of measured jet energy.

Advantage: no isolated leptons, which can indicate SM processes with *real* missing E_T (typically featuring $W \rightarrow lv$)

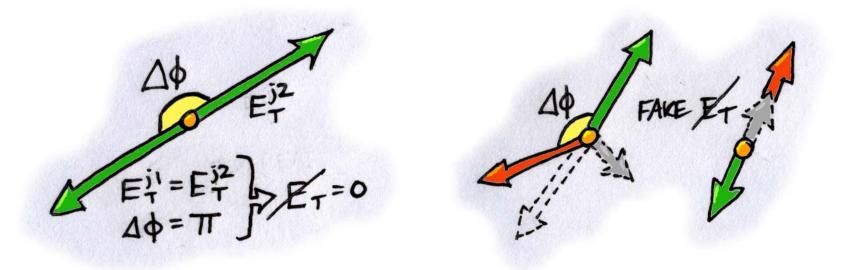
• Still leaves $Z \rightarrow \nu\nu$, $W \rightarrow \tau\nu$ (the τ decaying hadronically)

Disadvantage: Large QCD background. Statistically unlikely detector mismeasurements, that produce *fake* missing $E_{T'}$ start to overwhelm any non-SM signal.

Allowing for mismeasurement



QCD-like events will generally conserve transverse momentum.



Mismeasurement leads to the observation of "fake" missing E_T . The search described here takes the following approach:

• Use only "trusted" physics objects as input to the missing E_T calculations, ignoring unclustered energy and events with "anomalies".

• Use event observables that compensate for object mismeasurements.

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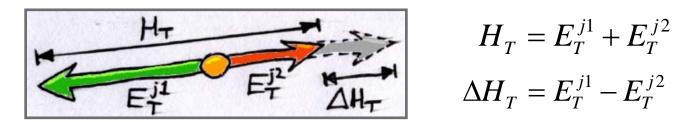
The α_T observable



Dijets: Randall (2008) proposed $\alpha = E^{j_2} / M(\text{dijet})$. We use $\alpha_T = E^{j_2} / M_T(\text{dijet})$:

Denominator:
$$M_T = \sqrt{2E_T^{j1}E_T^{j2}(1-\cos\Delta\phi)} = \sqrt{H_T^2 - \mathbf{p}_T^2}$$

Numerator: $E_T^{j2} = \frac{1}{2}(H_T - \Delta H_T)$



Conserved event, perfect measurement $\Rightarrow \alpha_T = 1/2$ Real missing E_{τ} : Small denom. $\Rightarrow \alpha_{\tau} > 1/2$ $\alpha_{T(N)} = \frac{1}{2} \frac{H_T - \min(\Delta H_T)}{\sqrt{H_T^2 - \mathbf{p}_T^2}}$ *Mismeasurement: Small num.* $\Rightarrow \alpha_T < 1/2$ Extension to *n*-jets: Generalise ΔH_T

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Event selection

Triggers:

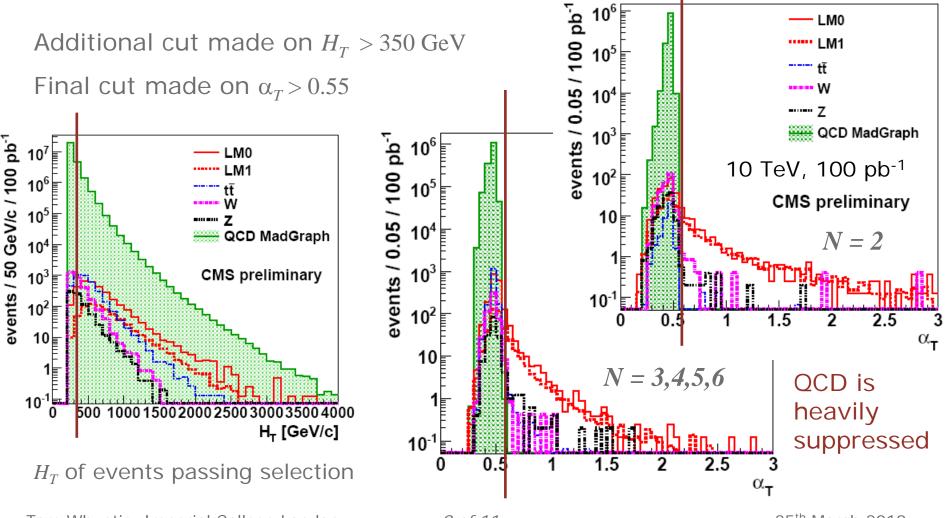
• High-Level Trigger (HLT) Single 110 GeV jet.

Pre-selection:

- Jet requirements (also defines event jet multiplicity, N):
 - » At least two jets with $p_T > 50 \text{ GeV}$, $|\eta| < 3.0$, EM fraction < 0.9;
 - » Leading jet $p_T > 100 \text{ GeV}, |\eta| < 2.0;$
 - » Second jet $p_T > 100 \text{ GeV}$;
- Lepton veto: Reject events with isolated *e* or μ , $p_T > 10 \text{ GeV}$;
- Photon veto: Reject events with isolated γ , $p_T > 25$ GeV;
- "Bad" jet veto: Reject events with $p_T > 50$ GeV jets that fail $|\eta|$, EM fraction requirements.



Applying the kinematic cuts



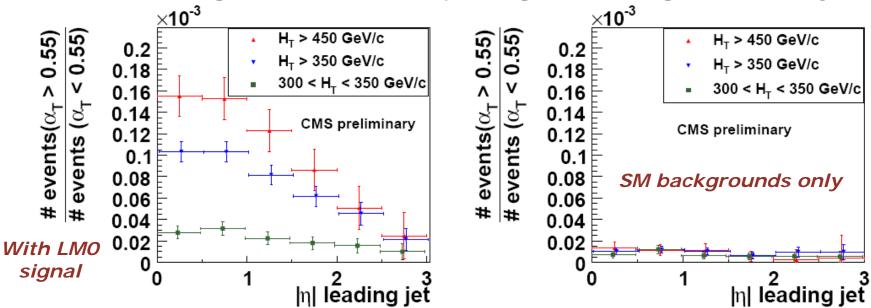
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Estimating SM contributions

1) Treat the backgrounds as one, exploiting non-SM signal centrality.



2) Estimate individual SM background contributions, e.g.

 $\frac{Z \rightarrow vv + jets}{Z \rightarrow \mu\mu}$ is statistics limited; use W+jets, γ +jets.

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d; tt + jet(s), W + jet(s), etc.Replace $W \rightarrow \mu \nu$ with $W \rightarrow \tau \nu$ in data using τ template.

Conclusions



We can look for SUSY in the all-hadronic channel:

• Signature: Large missing E_T + final state jets.

Non-SUSY backgrounds can be controlled with kinematics:

- Mismeasured QCD events are the dominant SM background;
- Compensating observables, e.g. α_T , can suppress these.

Tools exist for estimating real missing E_T SM backgrounds.