

Dark Matter, Low-Background Physics



1. Evidence (Astrophysical Detection)

2. Candidates, Properties

3. Direct Detection (Particle Physics)

1st Observation: 1930s



Fritz Zwicky



Virial Theorem: kinetic energy \propto potential energy
implies 400x more mass than visible!

Confirmation: 1980s

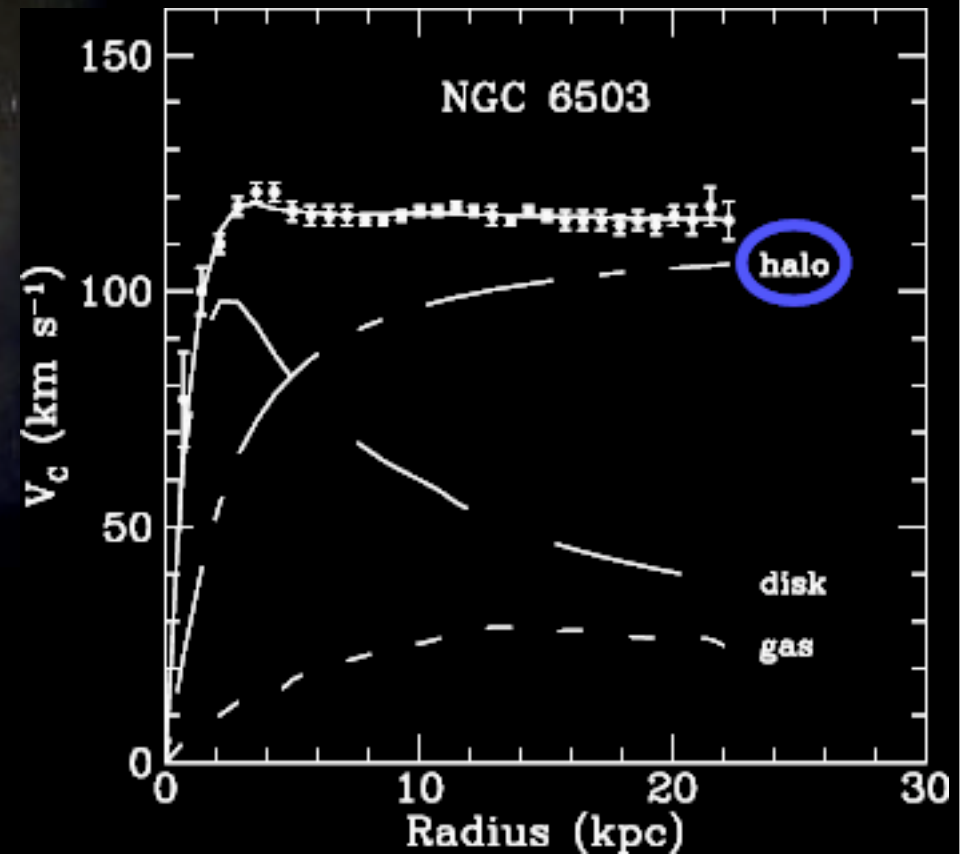
Vera Rubin



Rotation velocity $v(r)$ of
spiral galaxies



(M. Kamionkowski, astro-ph/9809214)



*implies 100x more mass
than visible!*

THE BIG BANG

INFLATION

COSMIC MICROWAVE
BACKGROUND
400,000 YEARS AFTER
BIG BANG

THE DARK AGES

FIRST STARS
400,000,000 YEARS
AFTER BIG BANG

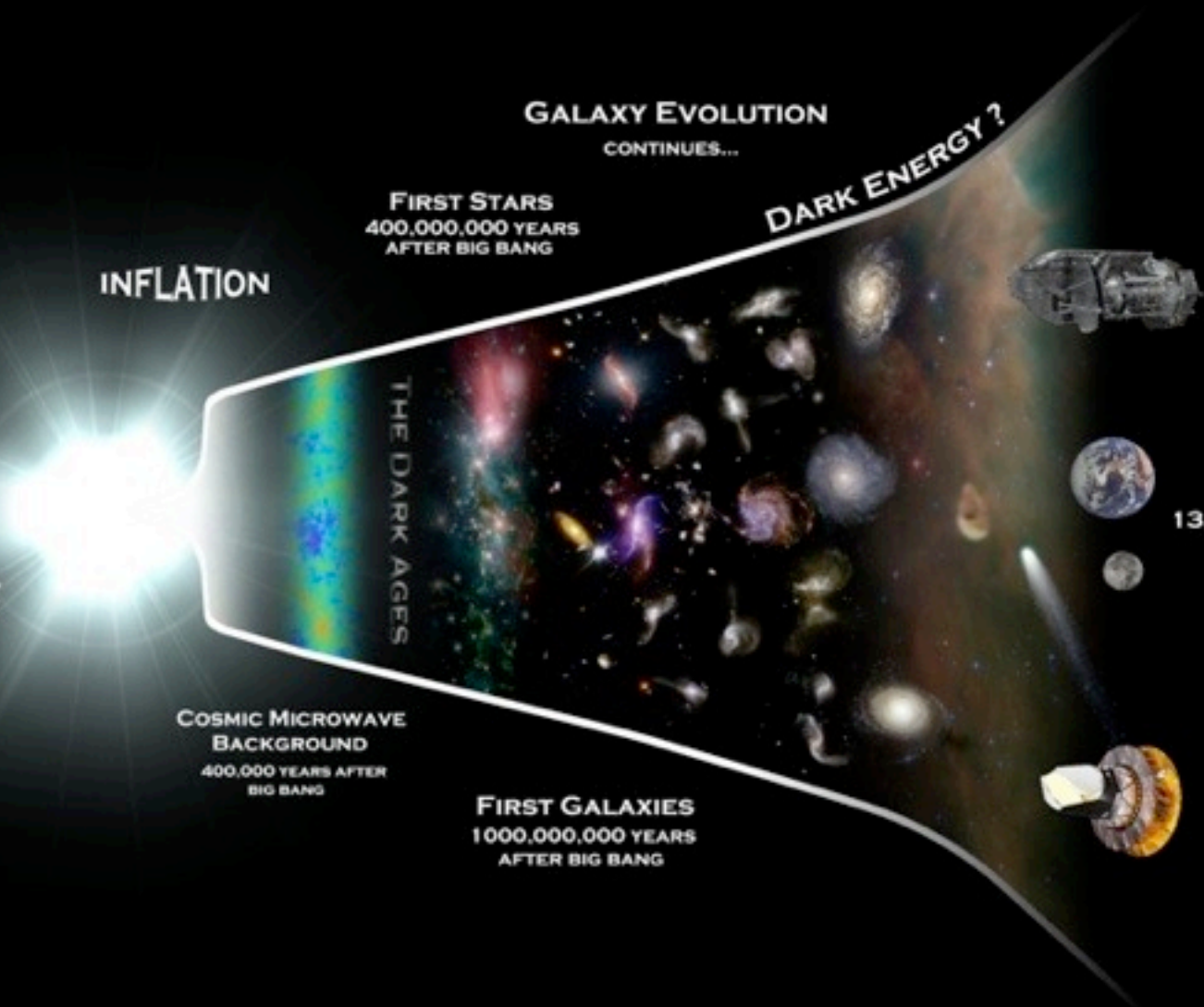
FIRST GALAXIES
1,000,000,000 YEARS
AFTER BIG BANG

GALAXY EVOLUTION
CONTINUES...

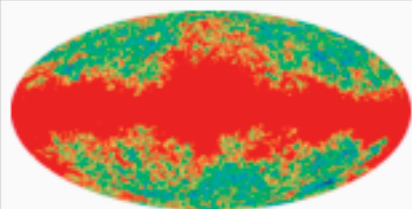
DARK ENERGY ?

FORMATION OF
THE SOLAR SYSTEM
8,700,000,000 YEARS
AFTER BIG BANG

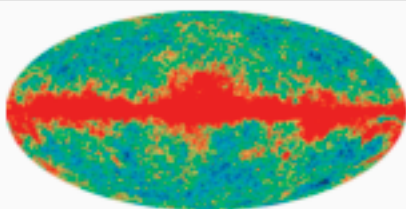
Now
13,700,000,000 YEARS
AFTER BIG BANG



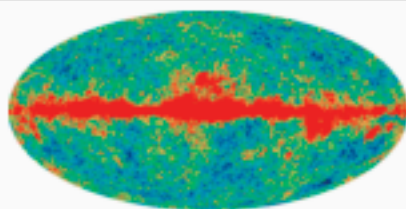
WMAP Cosmic Microwave Background



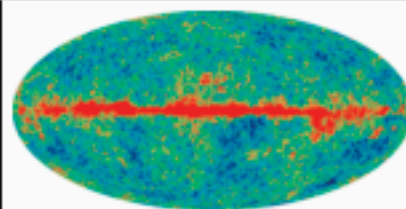
23 GHz



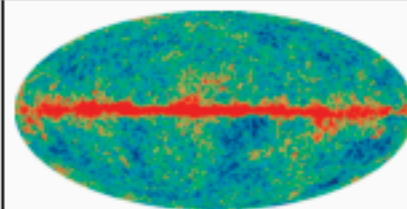
33 GHz



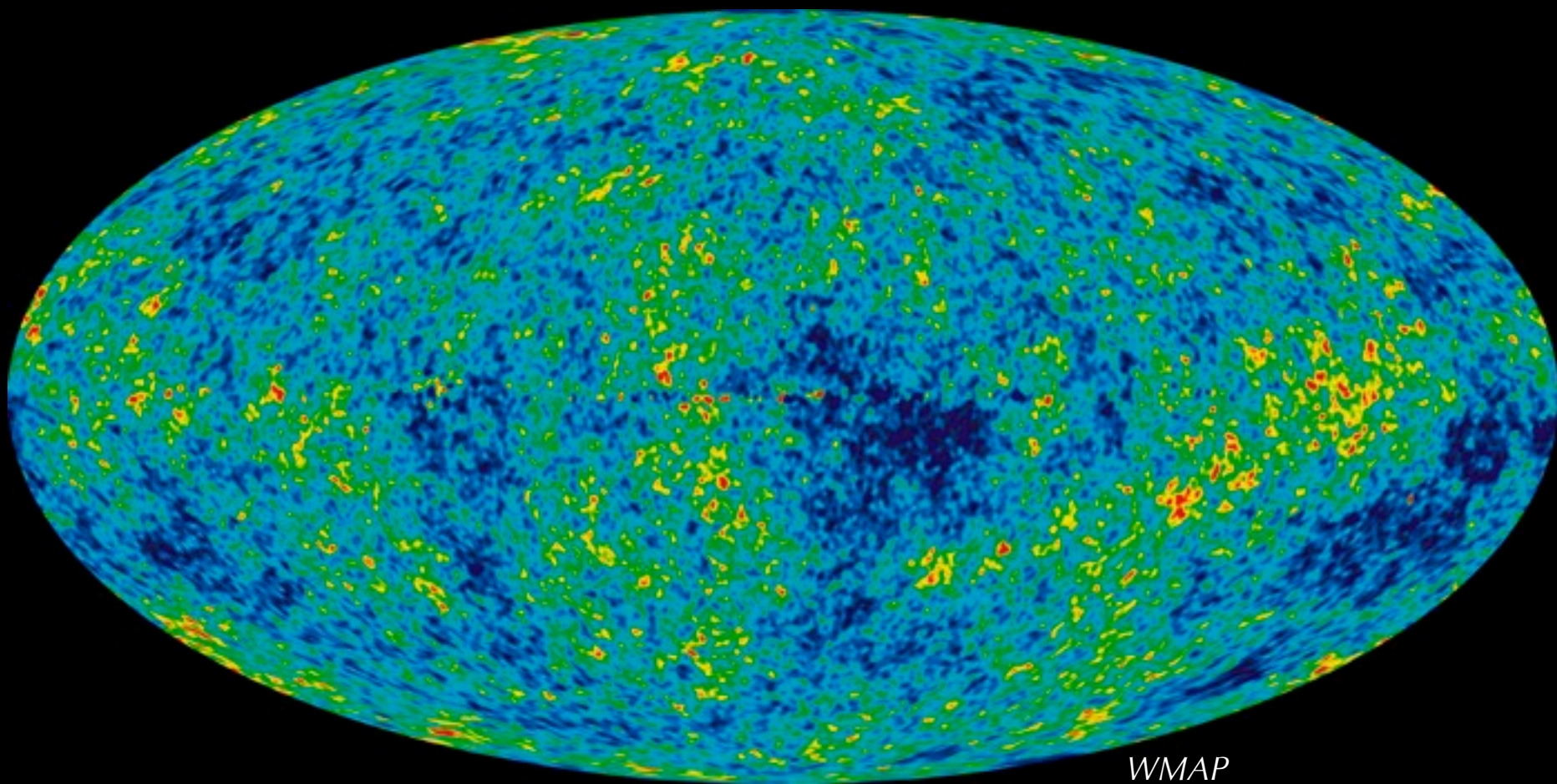
41 GHz



61 GHz



94 GHz

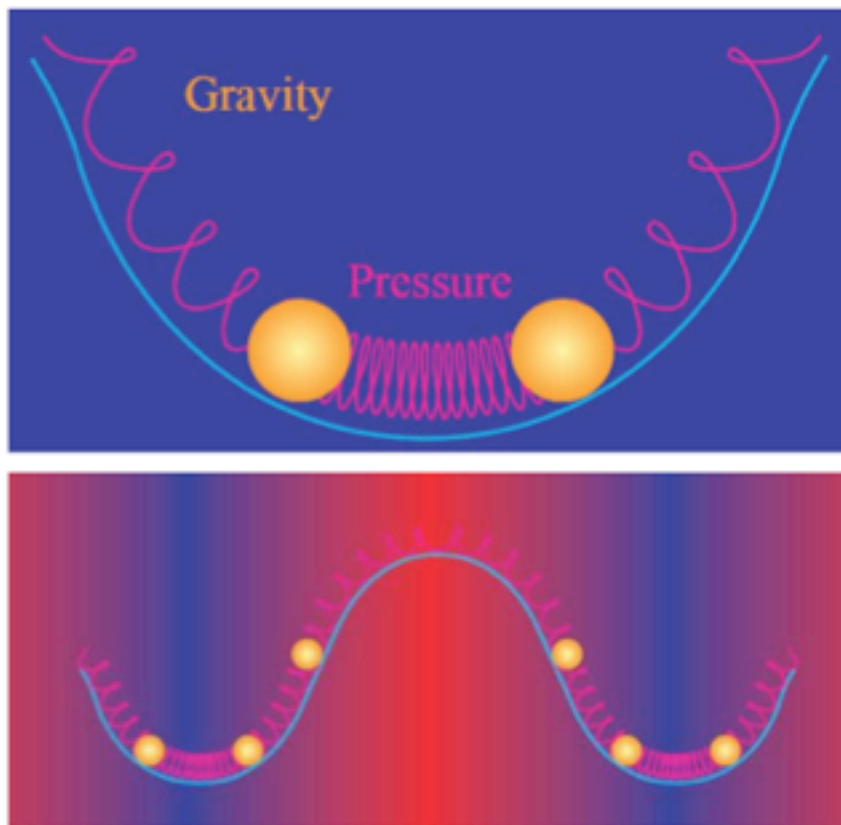


WMAP

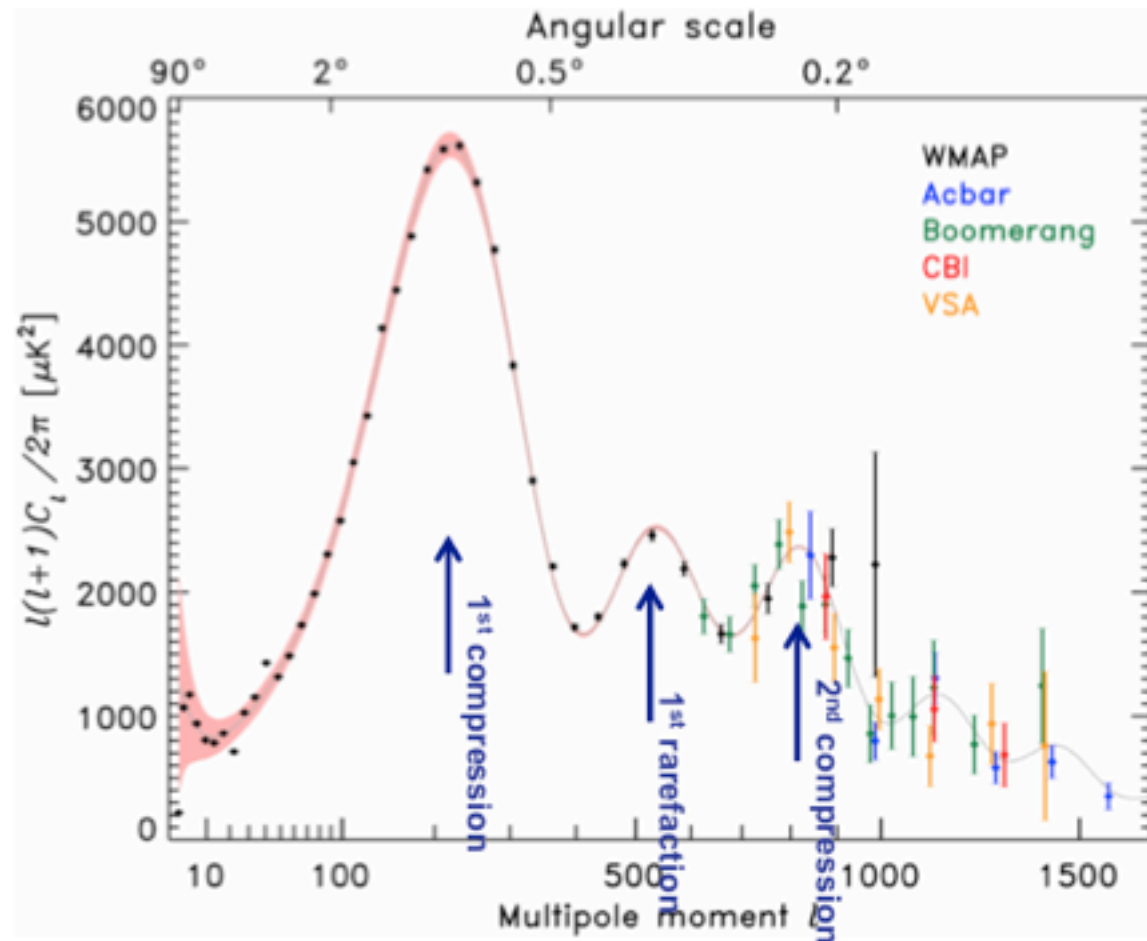
CMB Acoustic Spectrum

CMB “Baryometer”

From W. Wu (Chicago)

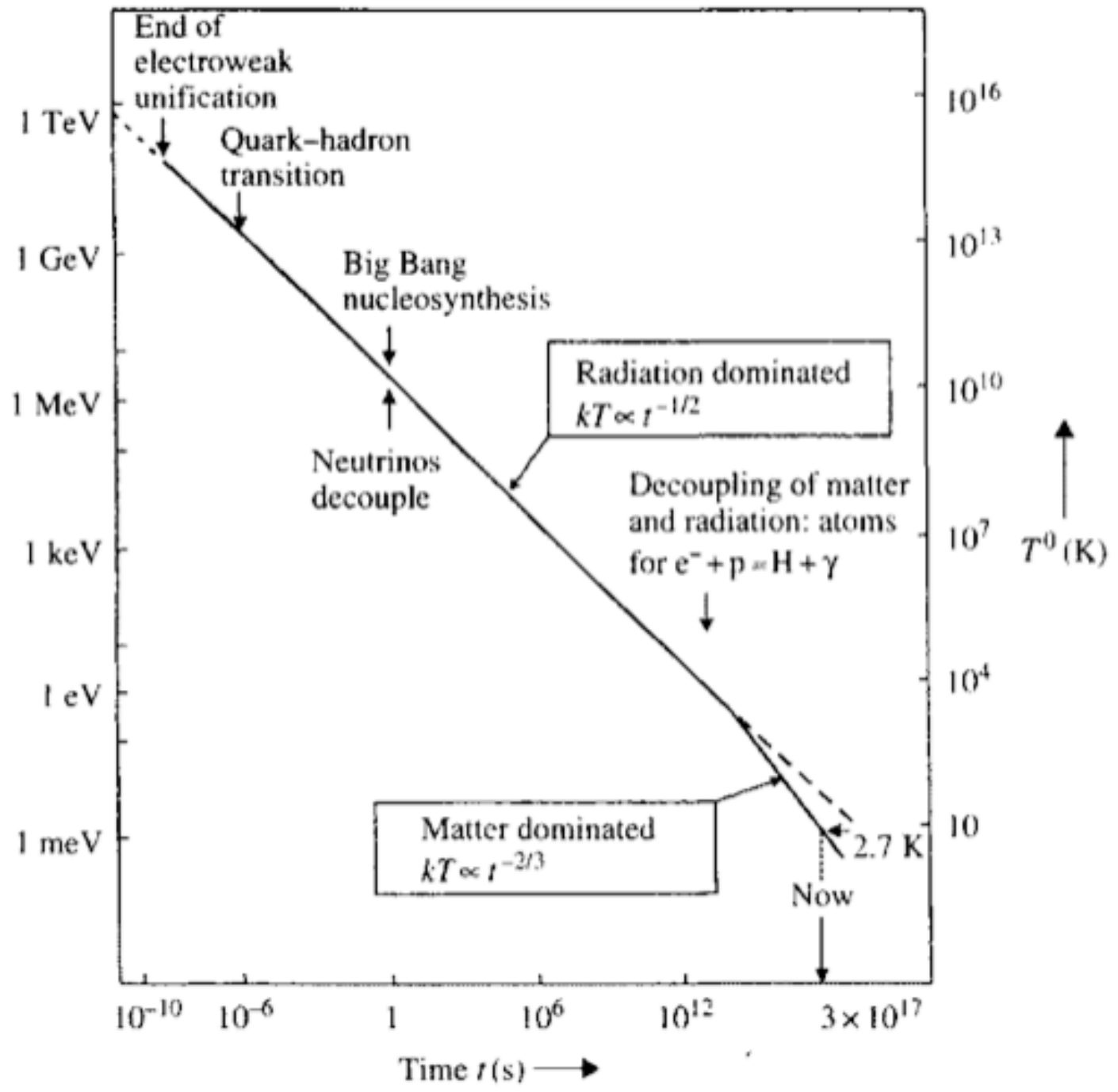


Baryon Acoustic Oscillations



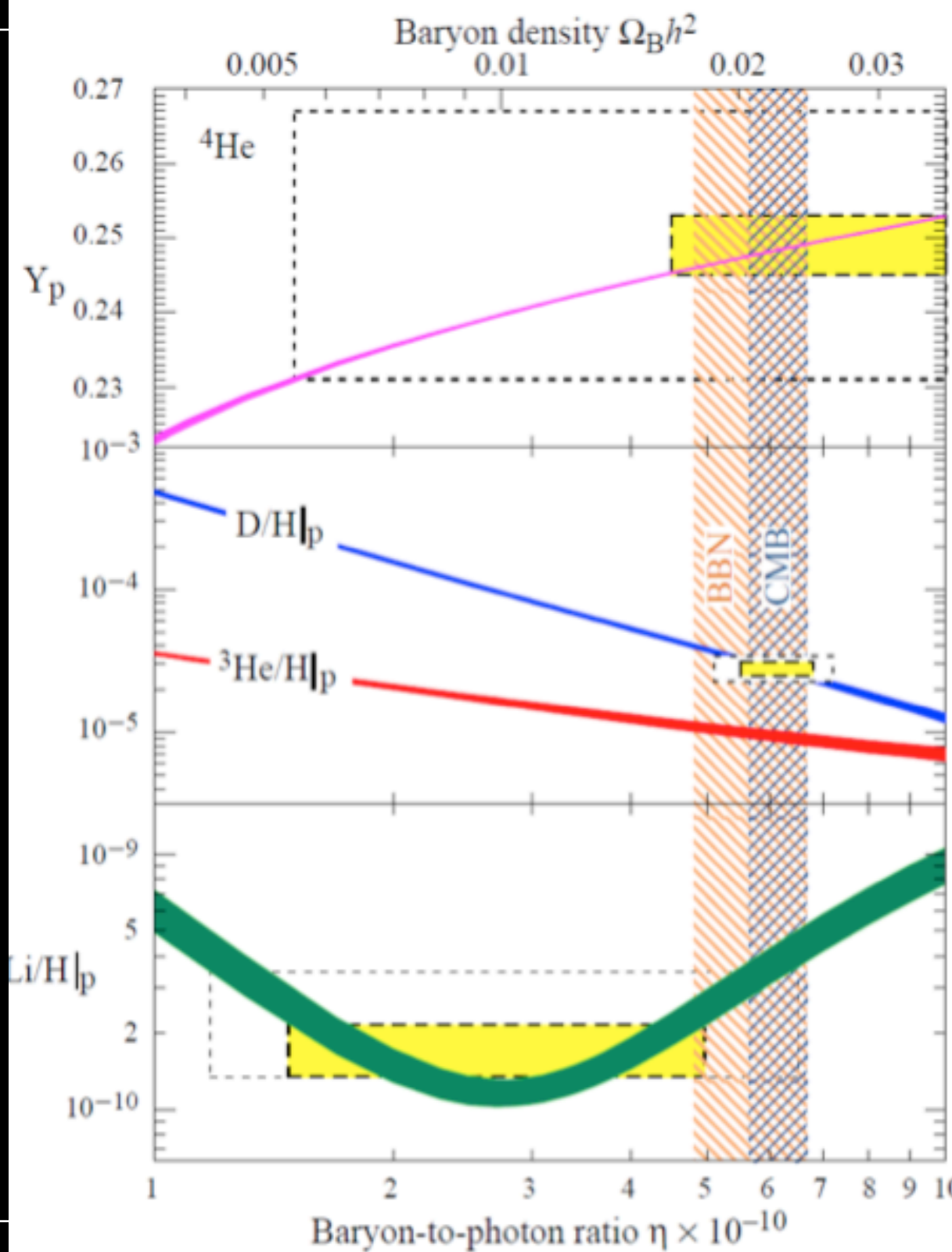
Baryons compress photon-baryon plasma at recombination, photons exert pressure, competition gives rise to pressure wave

BBN



D. H. Perkins, Particle Astrophysics (2004)

BBN

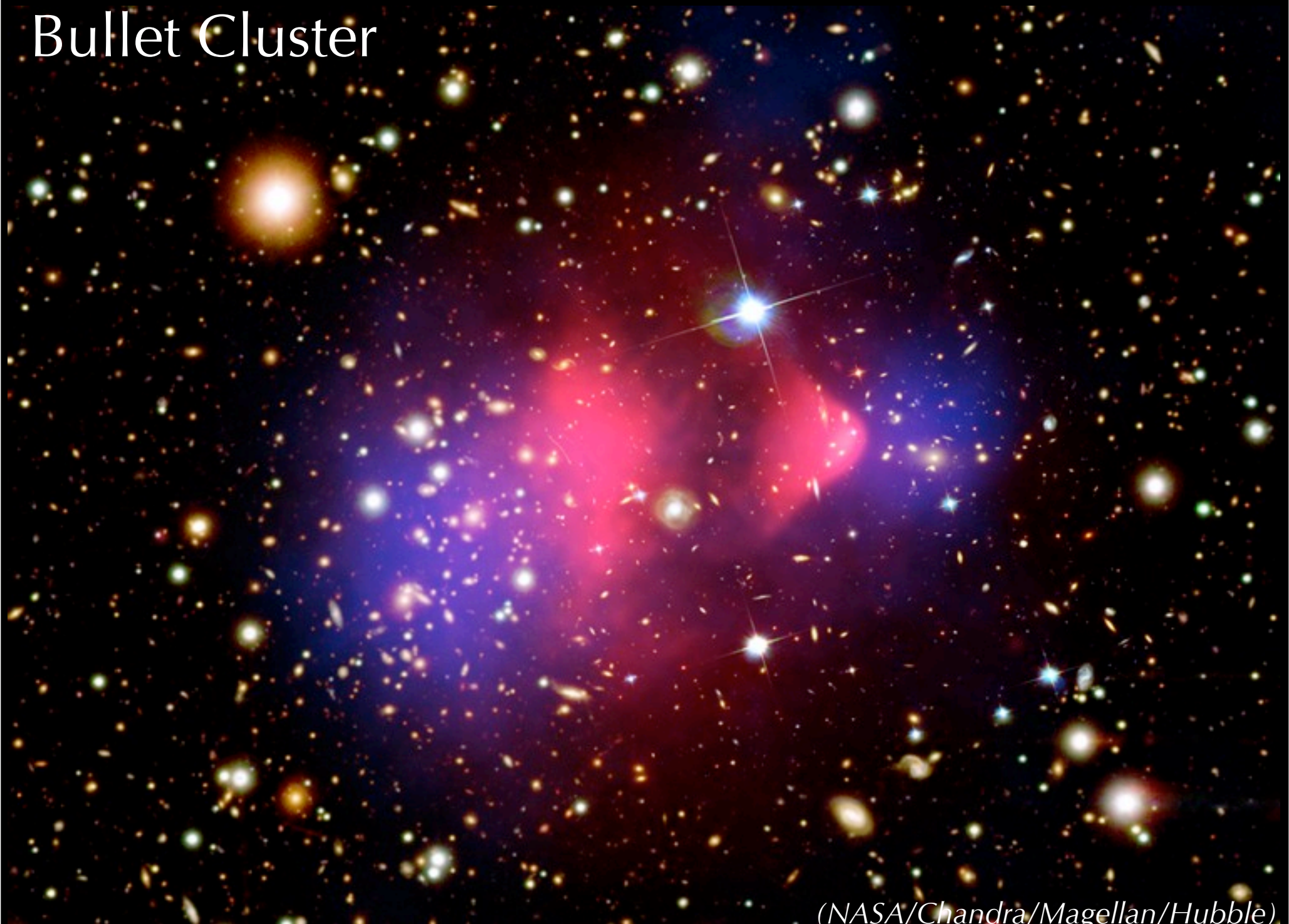


lines =
predicted
light element
abundances
vs. baryon
density

boxes =
observed
abundances
(1, 2 sigma)

vertical band =
CMB measure
of baryon
density
(PDG)

Bullet Cluster



(NASA/Chandra/Magellan/Hubble)

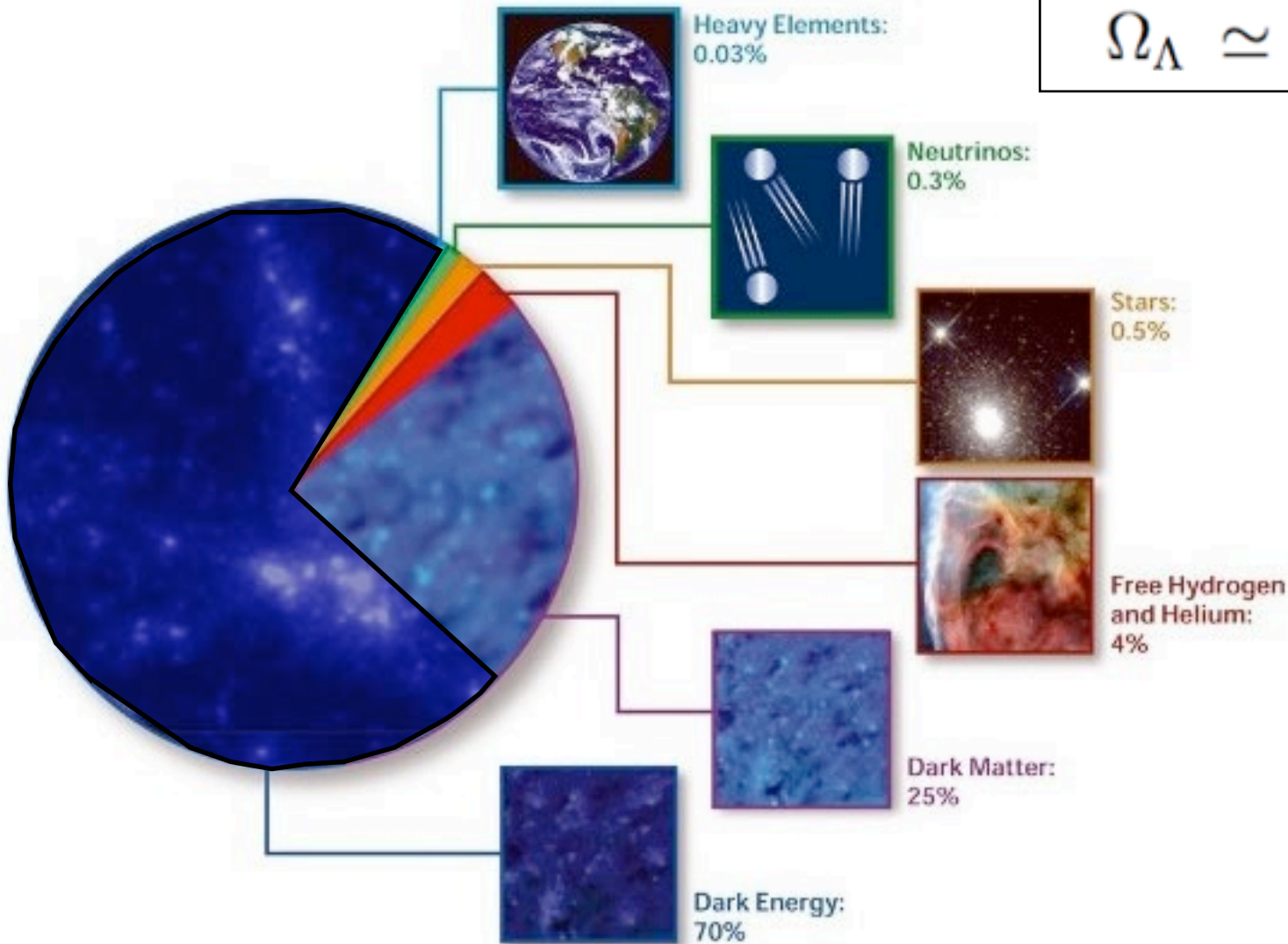
The Standard Model of Cosmology

$$\Omega_B \simeq 0.0456 \pm 0.0016$$

$$\Omega_{DM} \simeq 0.227 \pm 0.014$$

$$\Omega_\Lambda \simeq 0.728 \pm 0.015 .$$

E. Komatsu et al., Astrophys. J. Suppl 192 (2011) 18



Dark Matter is ~23% of the universe.

Dark Matter Candidates

strong

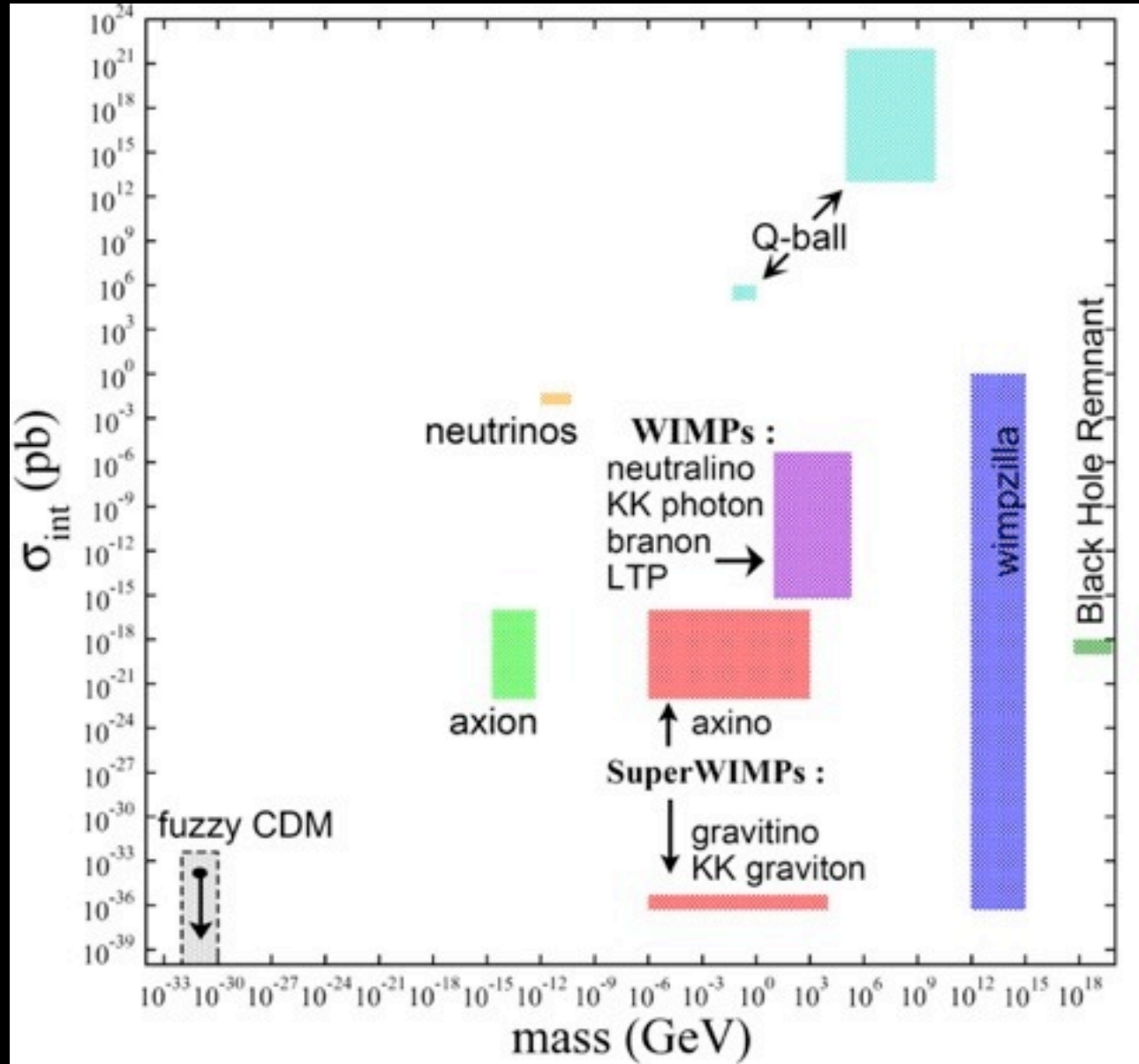
e.m.

weak

interaction
strengths



gravity



masses



neutrino?

electron

t-quark

Axions

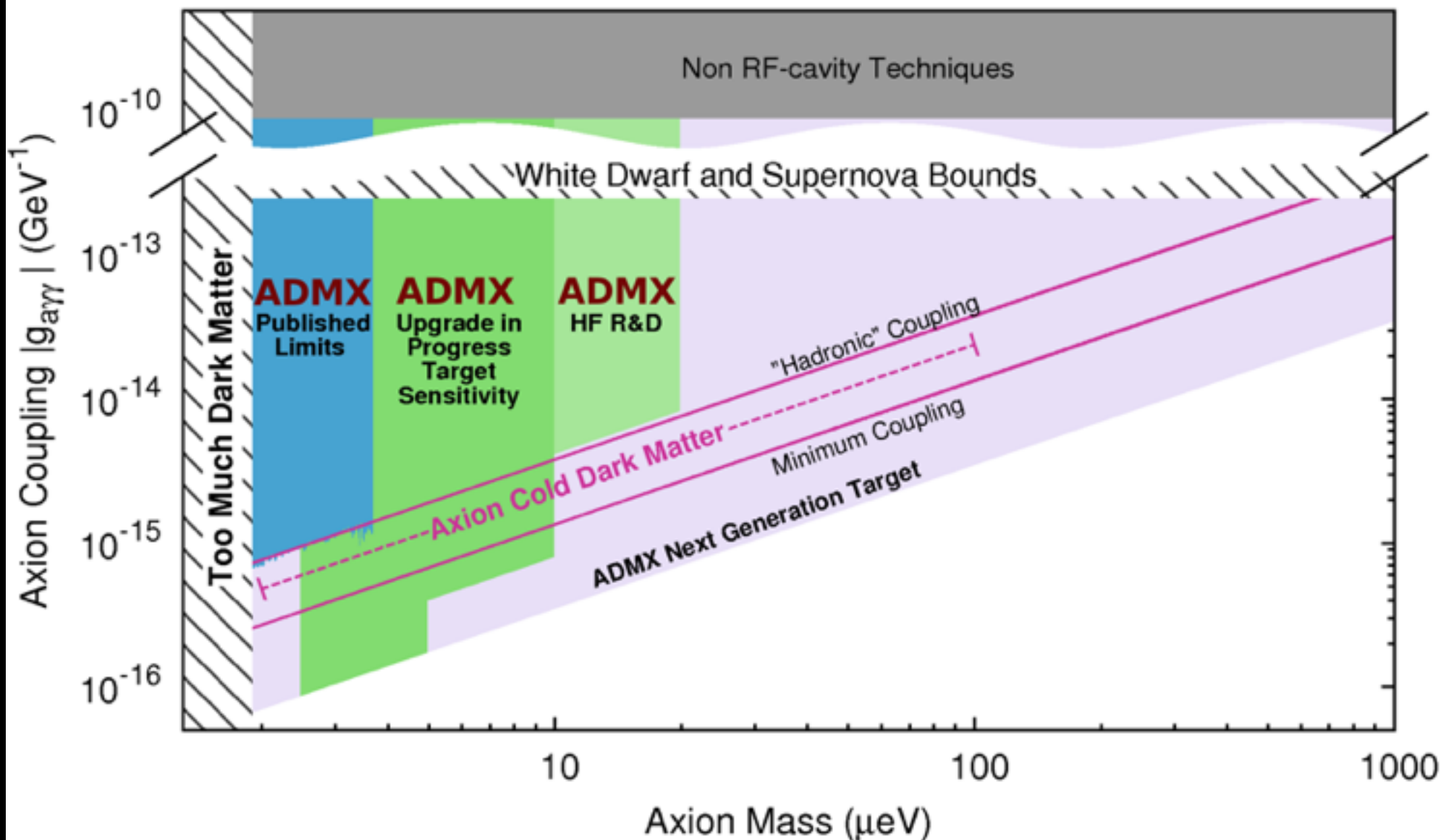
ADMX Achieved and Projected Sensitivity

Cavity Frequency (GHz)

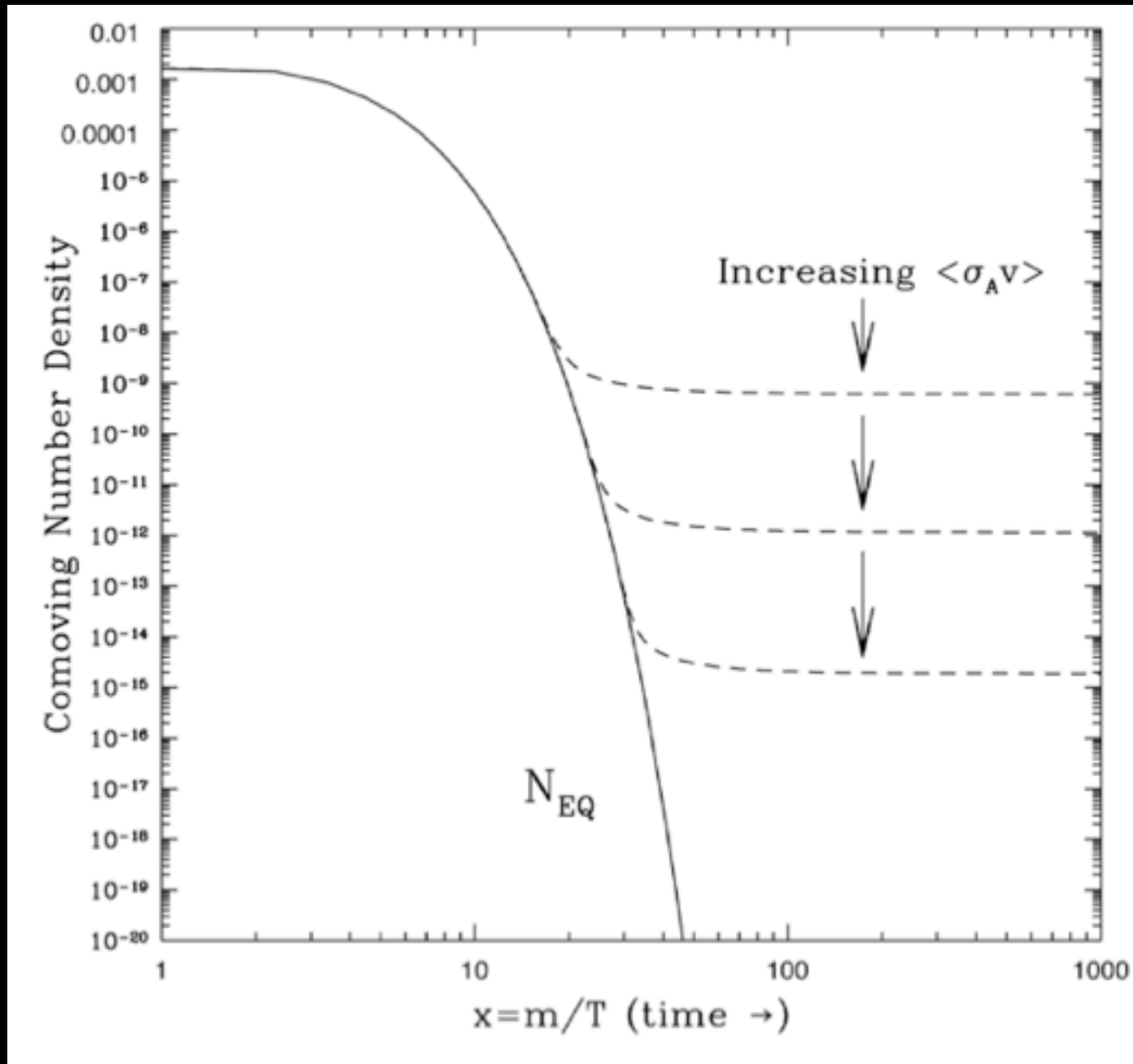
1

10

100



WIMP Number Density



D. H. Perkins, Particle Astrophysics (2004)

WIMP Mass Range

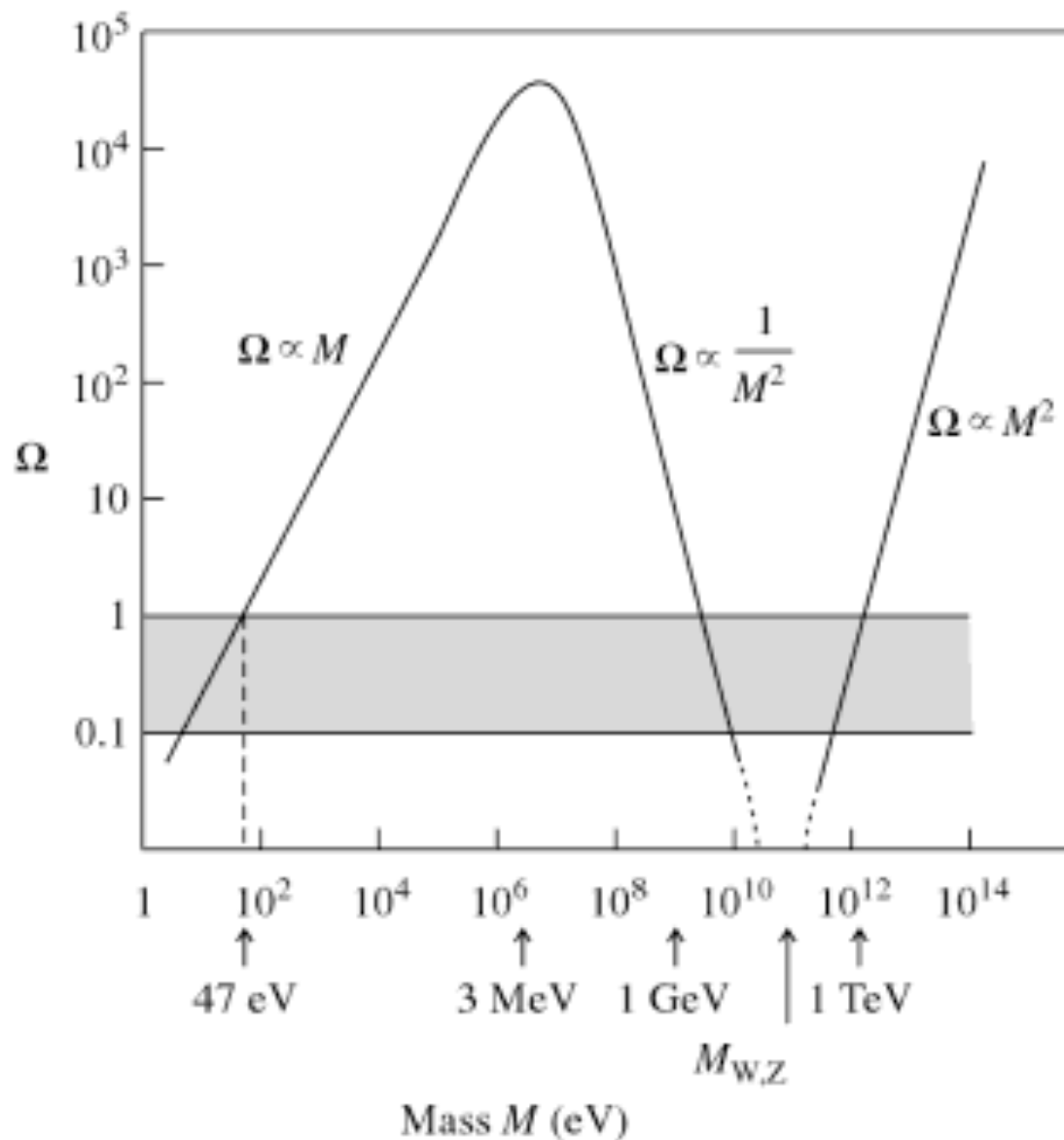
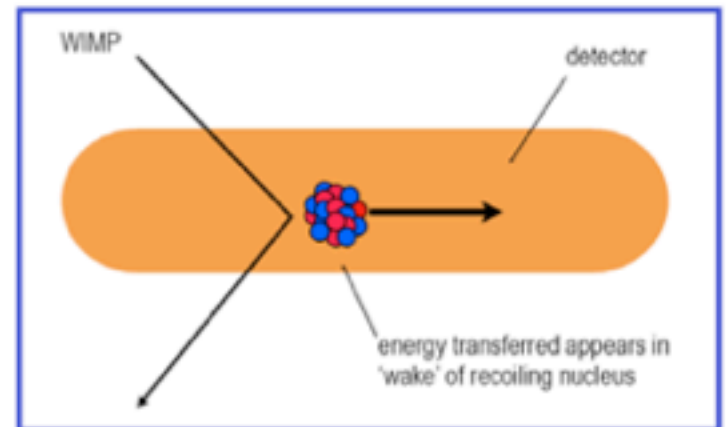
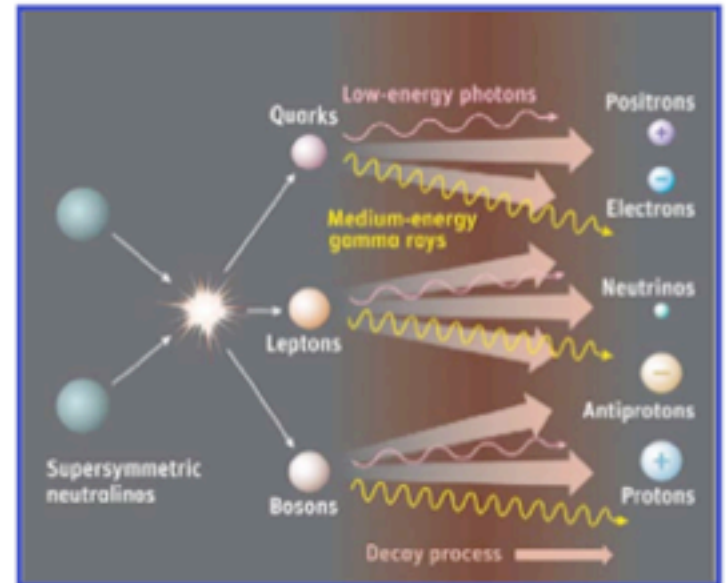
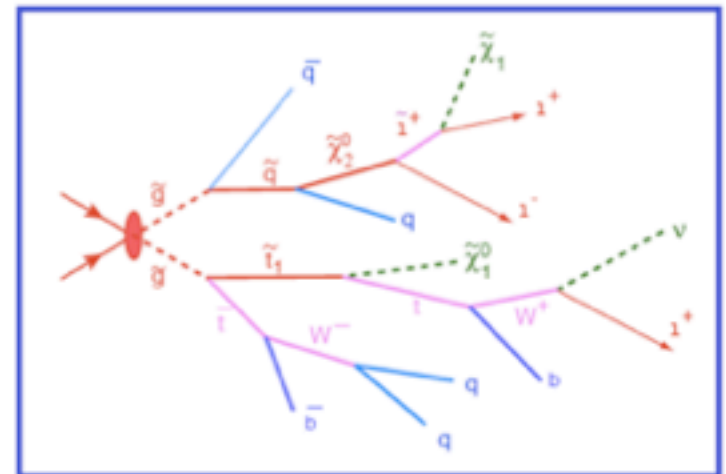
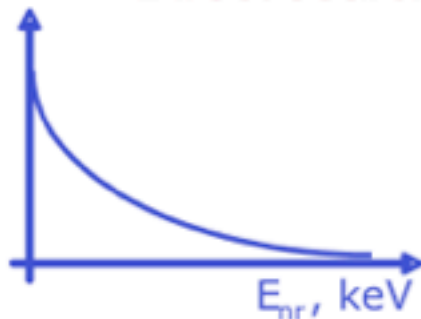


Fig. 7.11 Variation of the closure parameter with WIMP mass, assuming conventional weak coupling. The shaded region, corresponding to $\Omega = 0.1 - 1$, is that in which the contribution to the closure parameter from massive neutrinos or WIMPs must lie, thus excluding the range of masses 100 eV–3 GeV. Accelerator experiments suggest that WIMPs must have masses exceeding $M_Z/2 = 45$ GeV, otherwise Z bosons could decay into WIMP–antiWIMP pairs. However, for masses which are large compared with the Z boson mass, the weak cross-section falls rapidly because of propagator effects, so that WIMPs in the TeV mass range are possible dark matter candidates, depending on the precise values of the WIMP coupling.

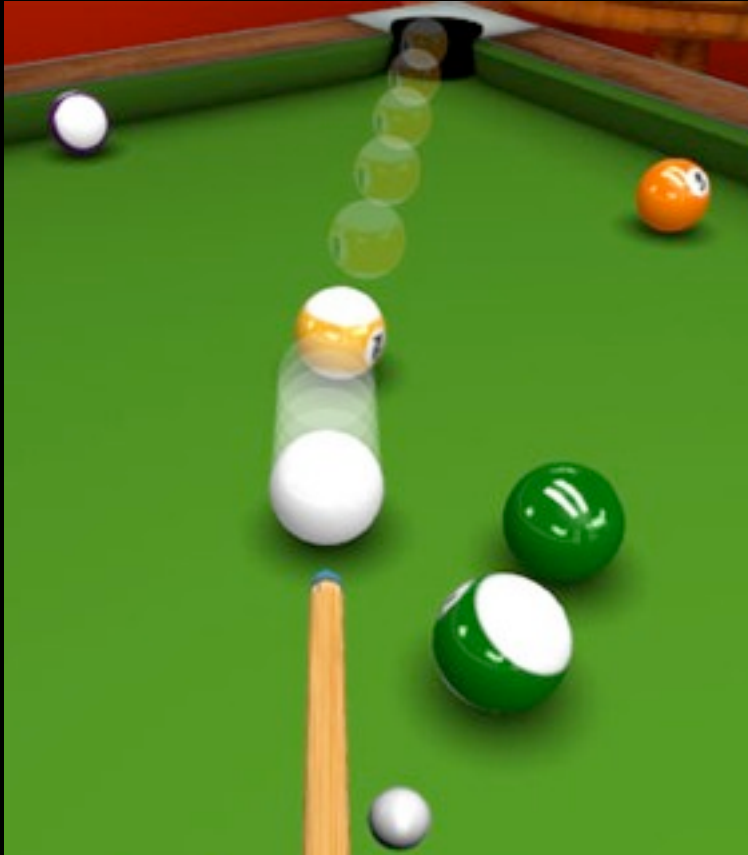
D. H. Perkins, Particle Astrophysics (2004)

Neutralino interactions

- Production
 - Accelerators such as LHC
 - Searches for missing E_T
- Annihilation
 - into fermion pairs, photons, ...
 - Indirect searches (FERMI, AMS,...)
- Elastic interaction
 - Spin-independent, $\sigma \sim A^2$
 - Spin-dependent, $\sigma \sim J/(J+1)$
 - Direct searches with terrestrial detectors



WIMP-Nucleon Scattering Event Rate



kinematics: $v/c \sim 1E-3!$

$$M_\chi = 100 \text{ GeV}$$

$$\sigma_{SI} = 1E-44 \text{ cm}^2$$

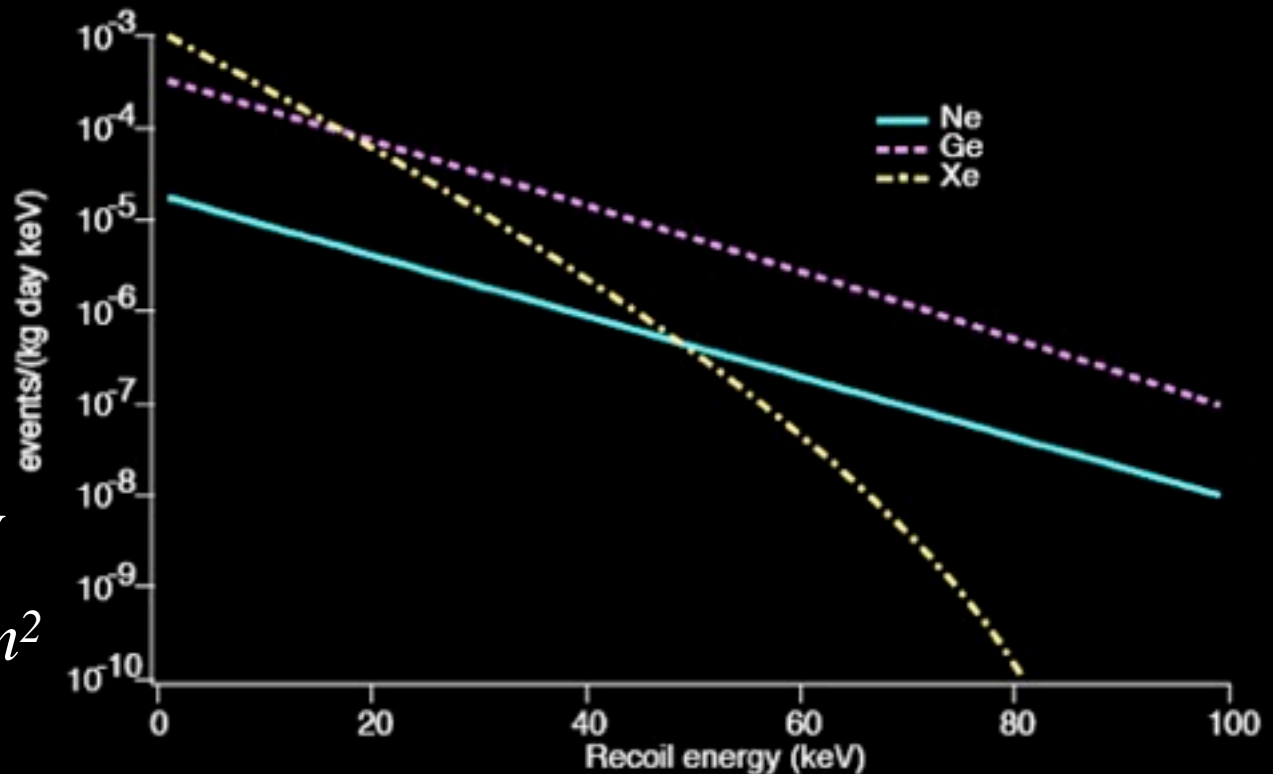
Spin Independent:

χ scatters coherently off of the entire nucleus A : $\sigma \sim A^2$

D. Z. Freedman, PRD 9, 1389 (1974)

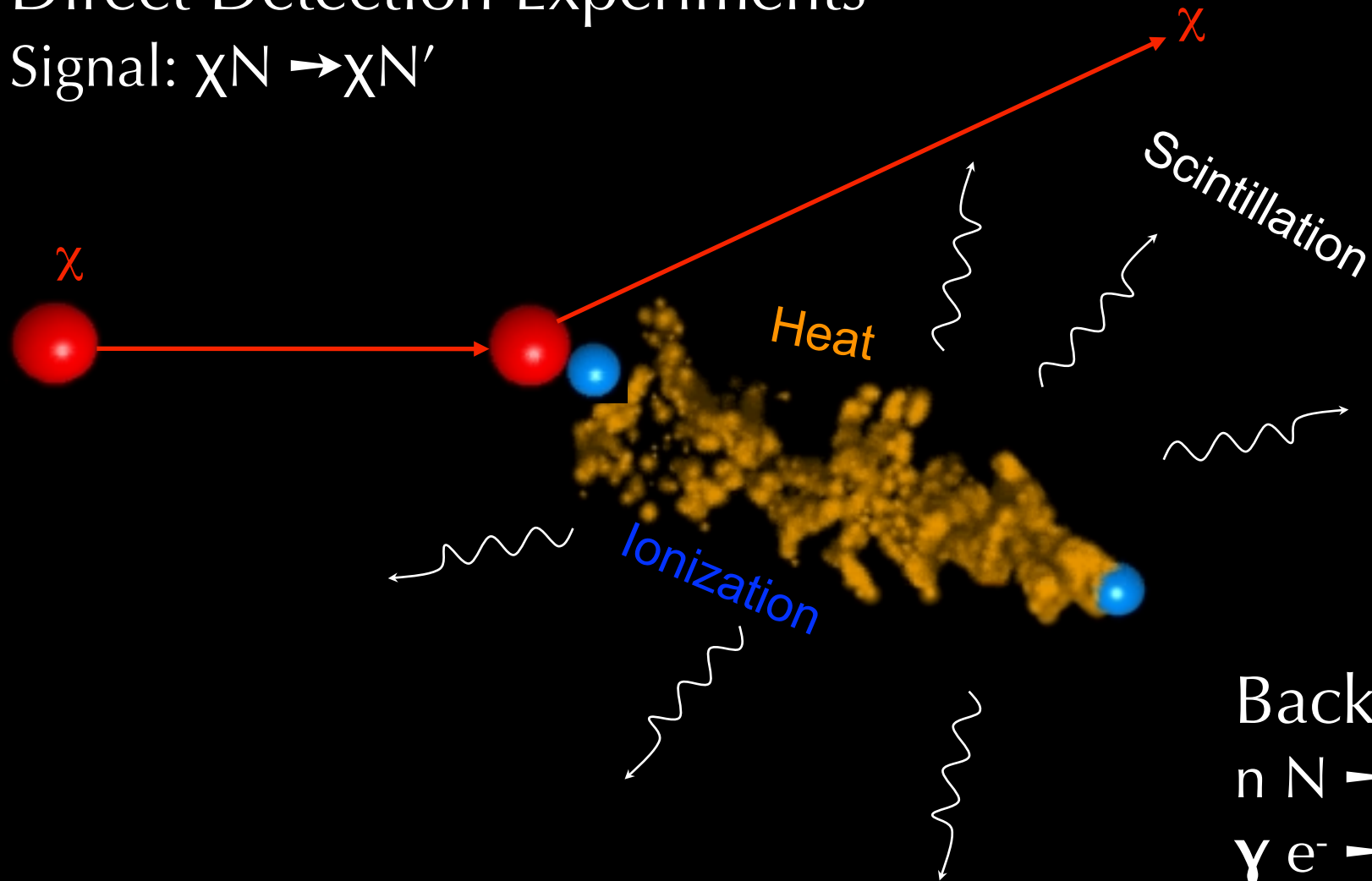
Spin Dependent:

only unpaired nucleons contribute to scattering amplitude: $\sigma \sim J(J+1)$



Direct Detection Experiments

Signal: $\chi N \rightarrow \chi N'$



Backgrounds:

$$n N \rightarrow n N'$$

$$\gamma e^- \rightarrow \gamma e^-$$

$$N \rightarrow N' + \alpha, e^-$$

$$\nu N \rightarrow \nu N'$$

DRIFT

IGEX

Picasso

Newage

COUPP

WARP

ArDM

ZepInIII

CDMS

Edelweiss

Xenon100

LUX

CoGeNT

DMTPC

ANaIS

DEAP/CLEAN

XMASS

KIMS

Dama/LIBRA

CRESSTII

ROSEBUD

CRESST

Ionization!

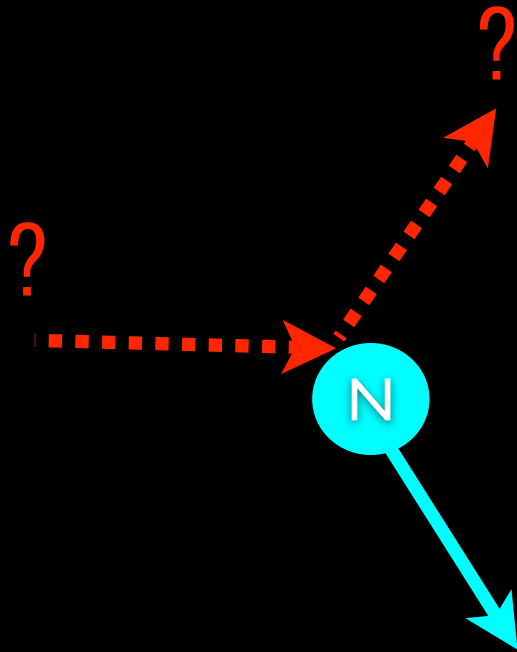
Heat!

Scintillation!

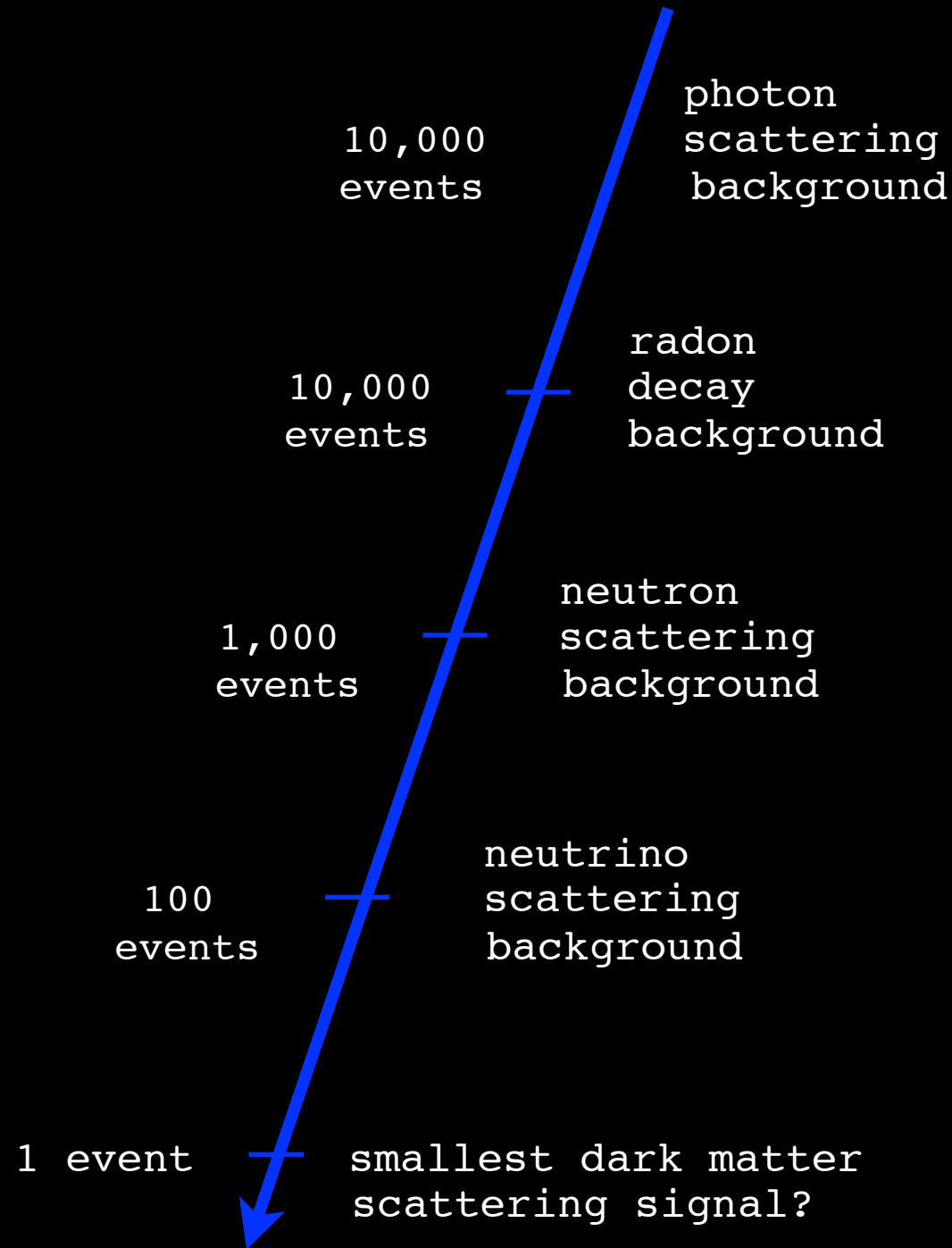


Backgrounds

In dark matter experiments...



*Anything else that does this
can fake a dark matter signal!*

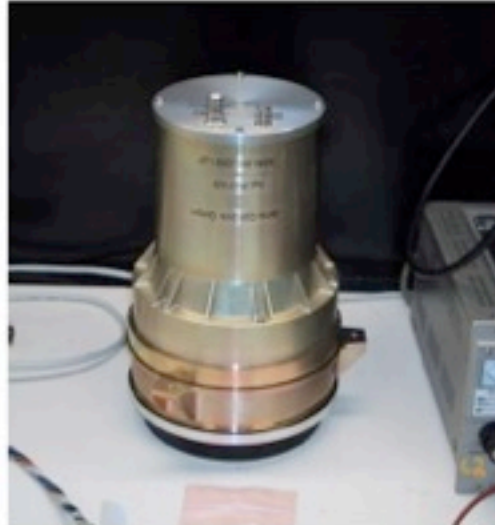


EM Backgrounds

(D. McKinsey)



Geiger counter



Sodium iodide crystal



Germanium

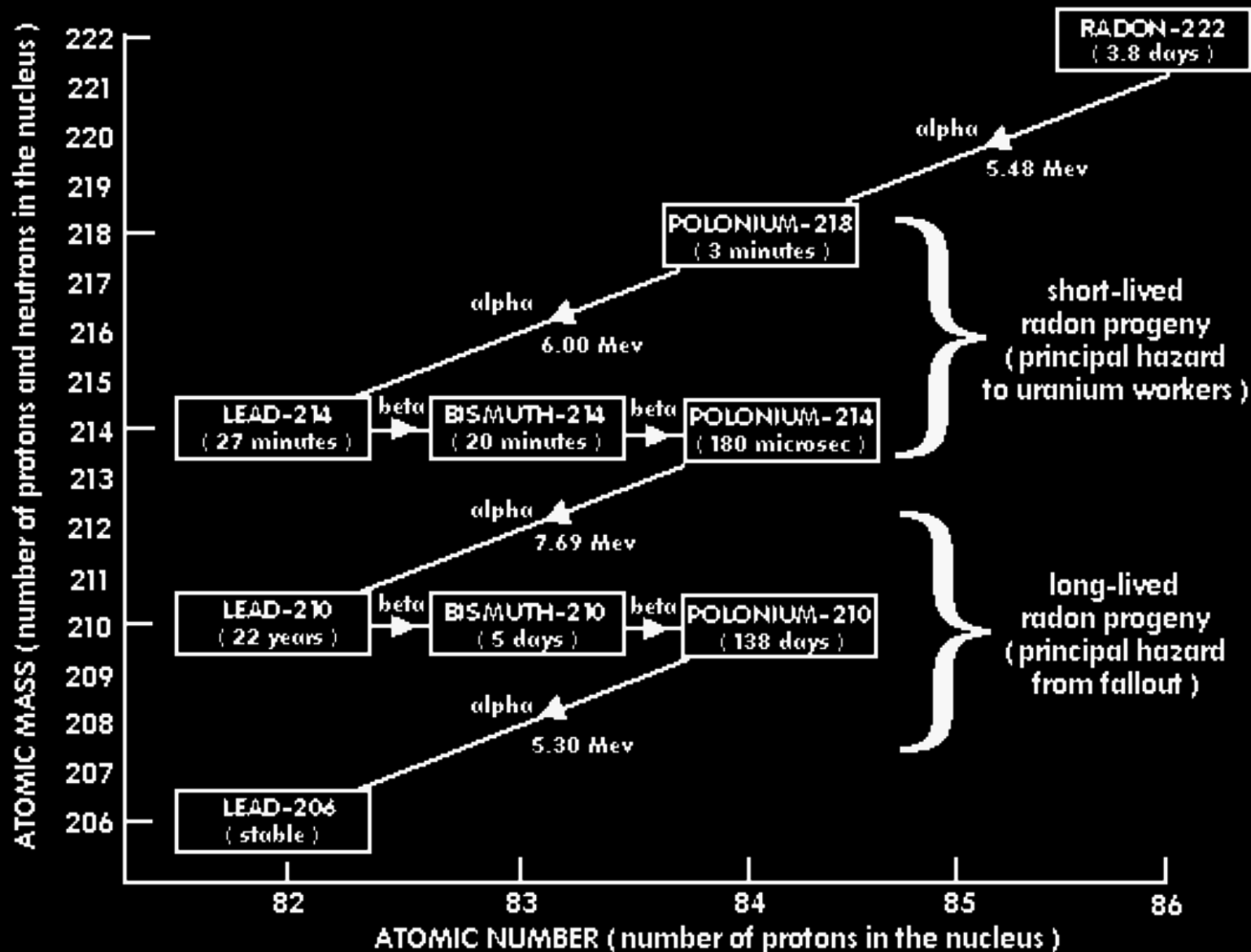
Gamma ray interaction rate is proportional to
(# of electrons in detector) \times (gamma ray flux)

Typical count rate = 100 events/s/kg = 10,000,000 events/day/kg
in a good lead shield, rate drops to 100 events/day/kg

Best dark matter detectors: sensitive to **0.01 events/day/kg**
($\sigma \sim 1\text{E-}44 \text{ cm}^2$)

U and Th Decay Backgrounds

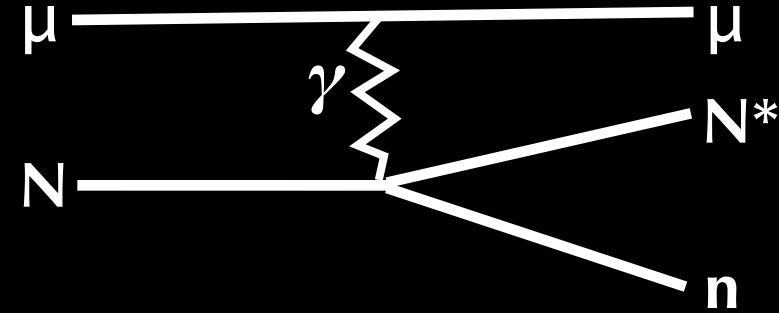
can't shield a detector from U and Th inside, recoiling progeny and associated betas can fake nuclear recoils



Neutron Backgrounds

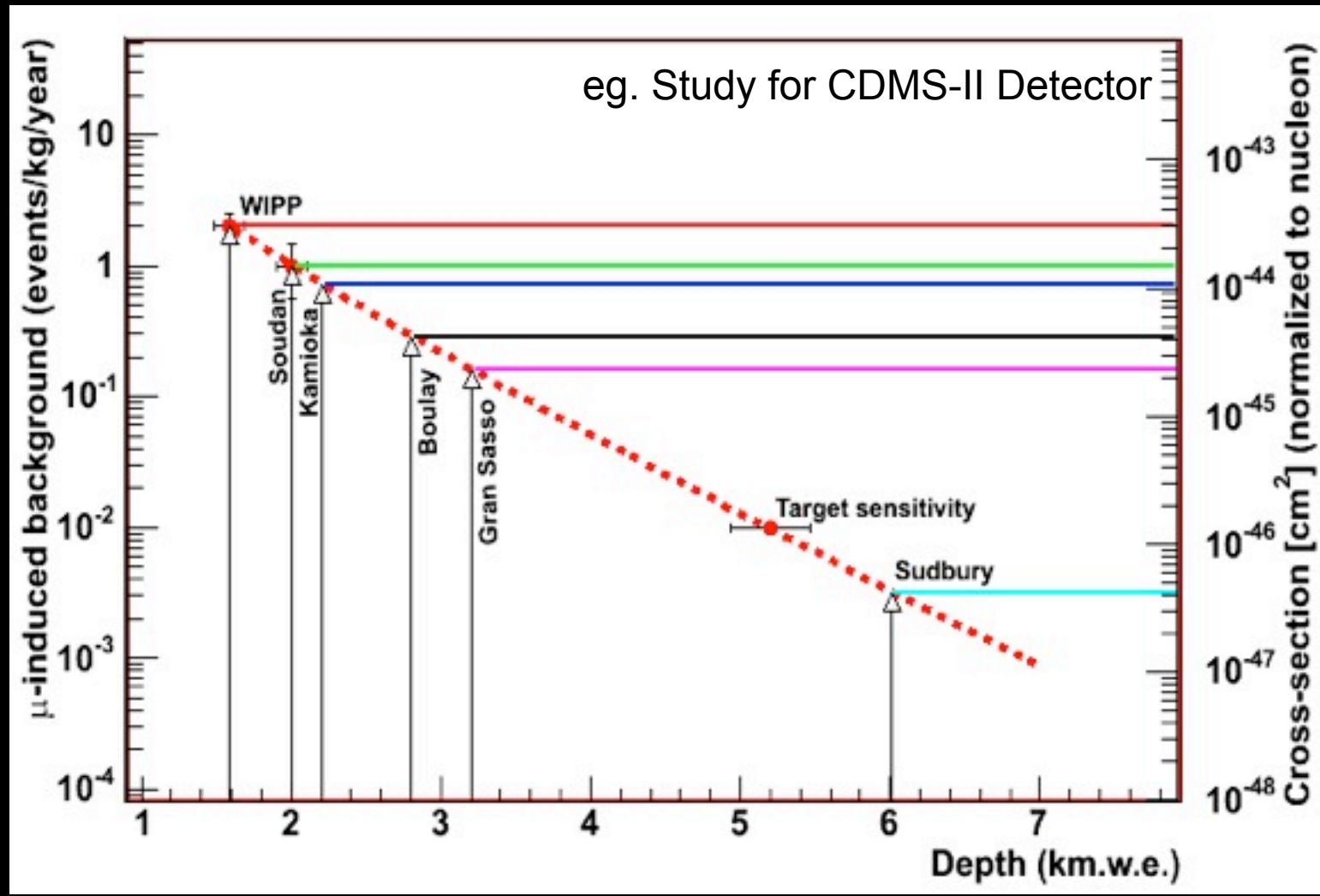


D.-M. Mei, A. Hime, PRD73:053004 (2006)



Cosmic muons
spall neutrons:
 $\sim 10^{-4}$ neutrons/
(100 GeV μ)/
gm/cm²

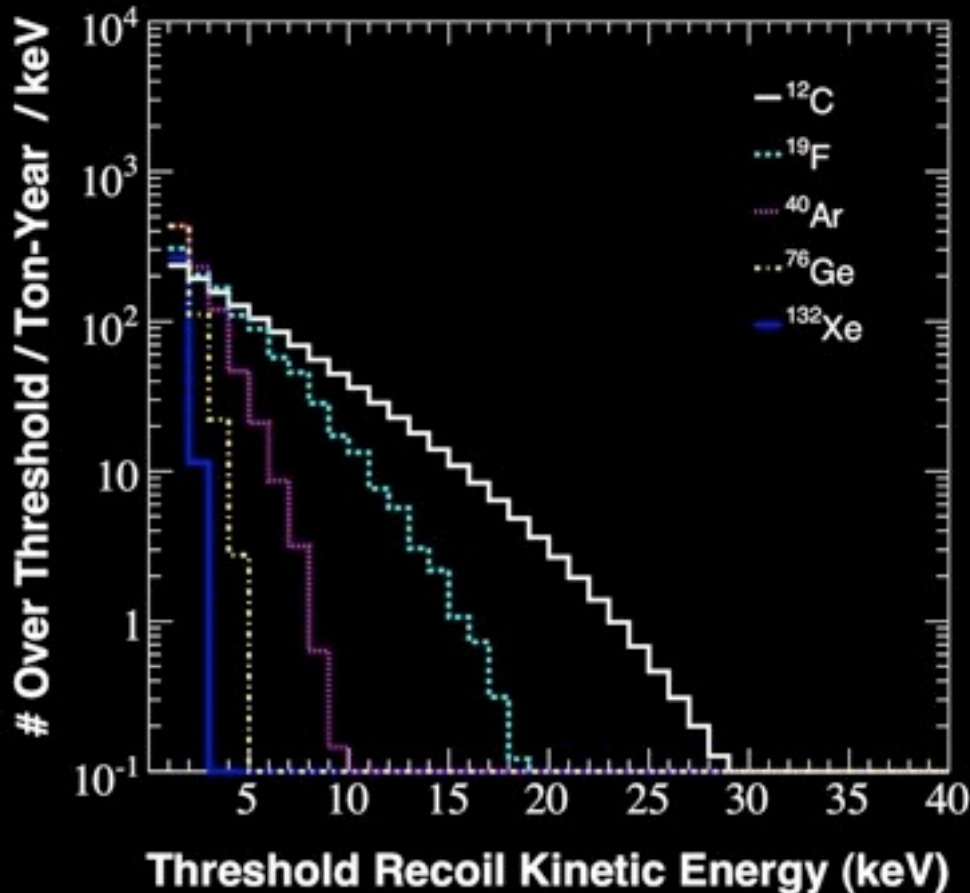
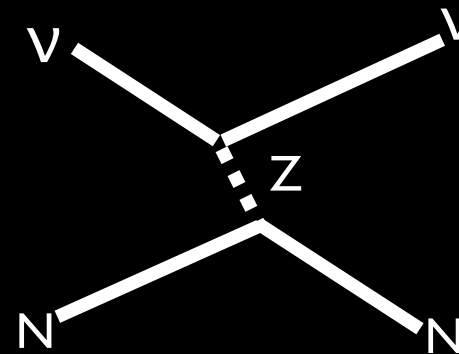
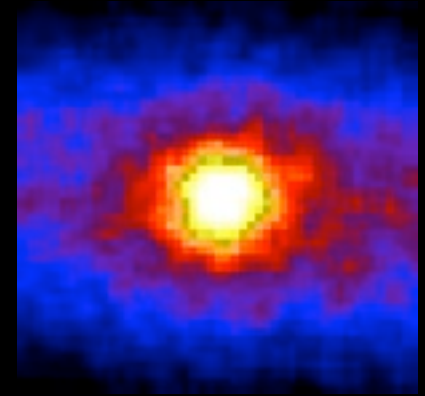
neutron flux:
 $10^{-8} - 10^{-10}$ /cm²/s
(range for depth)



ν Backgrounds

can't shield a detector from
coherent elastic scattering of
solar neutrinos

$$\Phi(B^8) = 5.86 \times 10^6 \text{ cm}^{-2} \text{ s}^{-1}$$

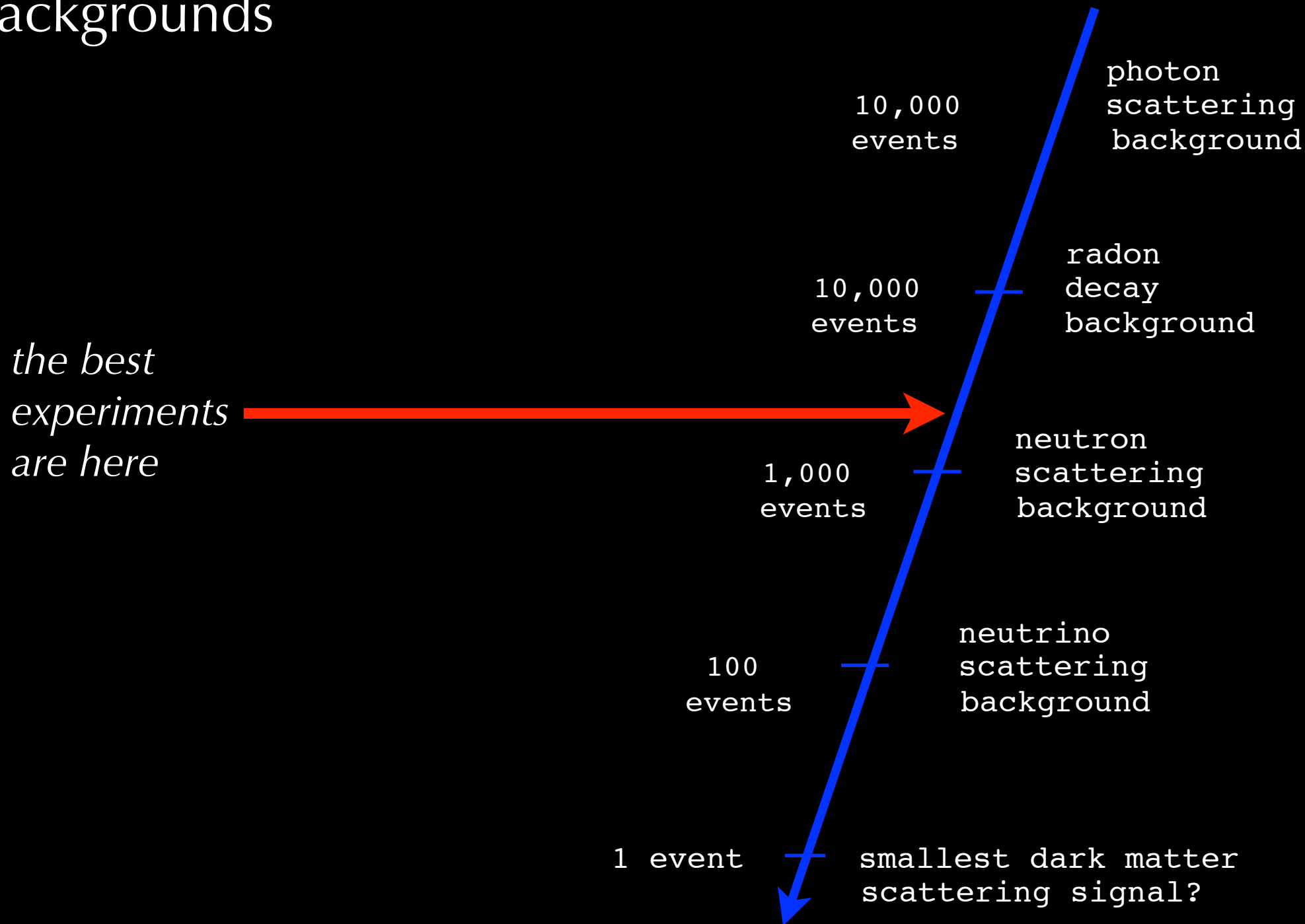


100 events/ton-year =
 $\sim 10^{-46} \text{ cm}^2$ limit

unless you measure
the direction!

JM, P. Fisher, PRD76:033007 (2007)

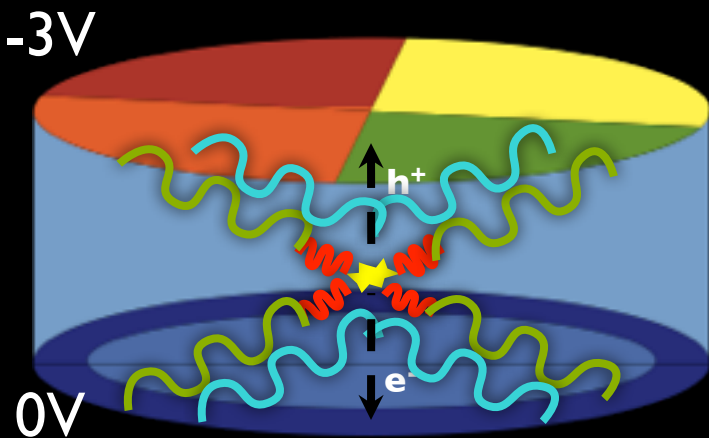
Backgrounds



CDMS

(material thanks to E. Figueroa)

Phonon side: 4 quadrants of phonon sensors for energy & position (timing)

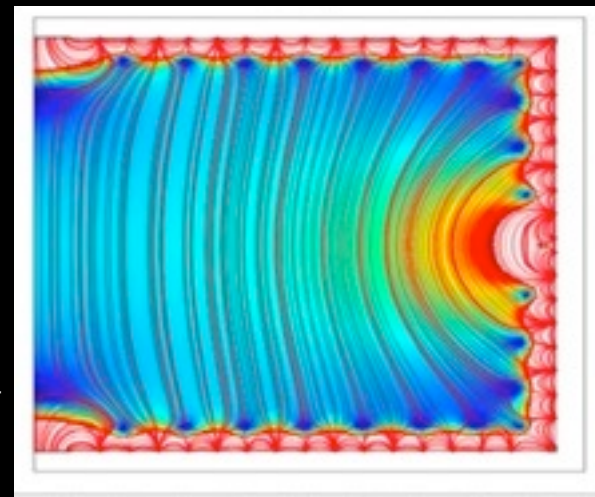


Transition Edge Sensors, operated at ~ 40 mK on Ge and Si crystals, in Soudan

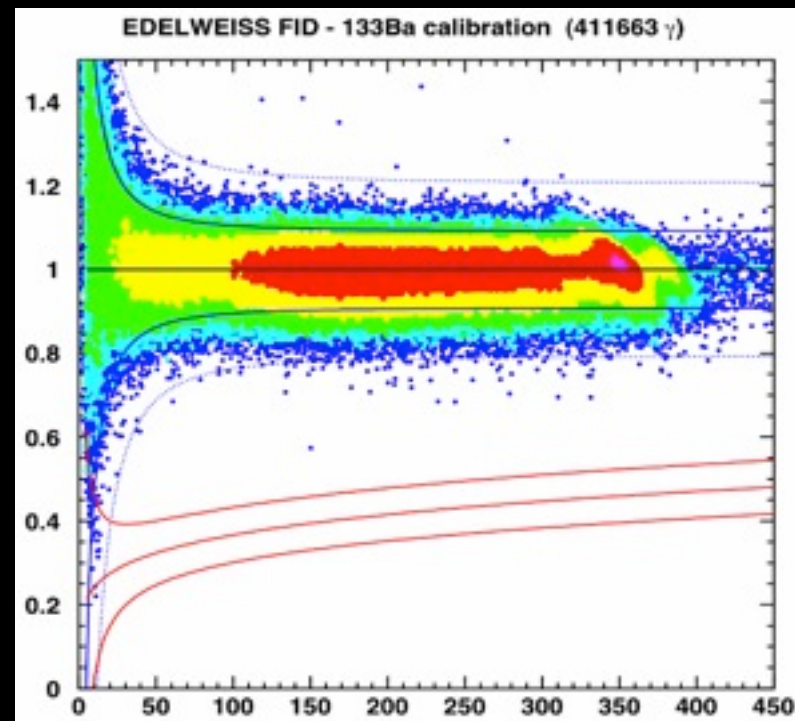
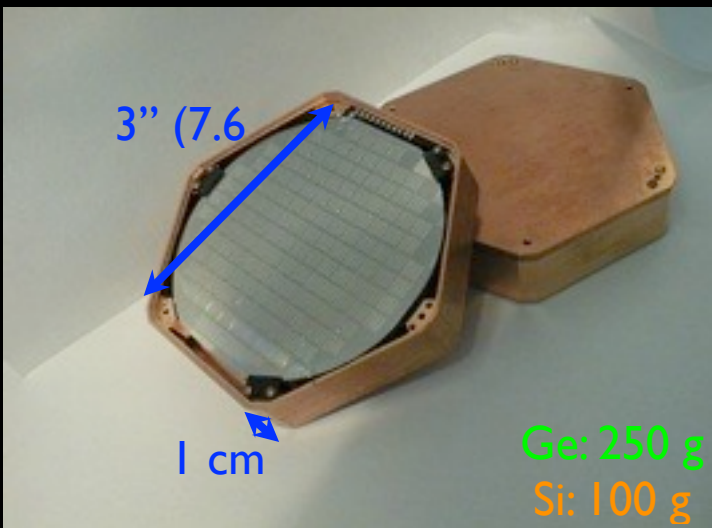
arXiv:0907.1438v1

arXiv:0912.3592v1

Detector re-design a la EDELWEISS to reduce surface backgrounds $\times 10^4$



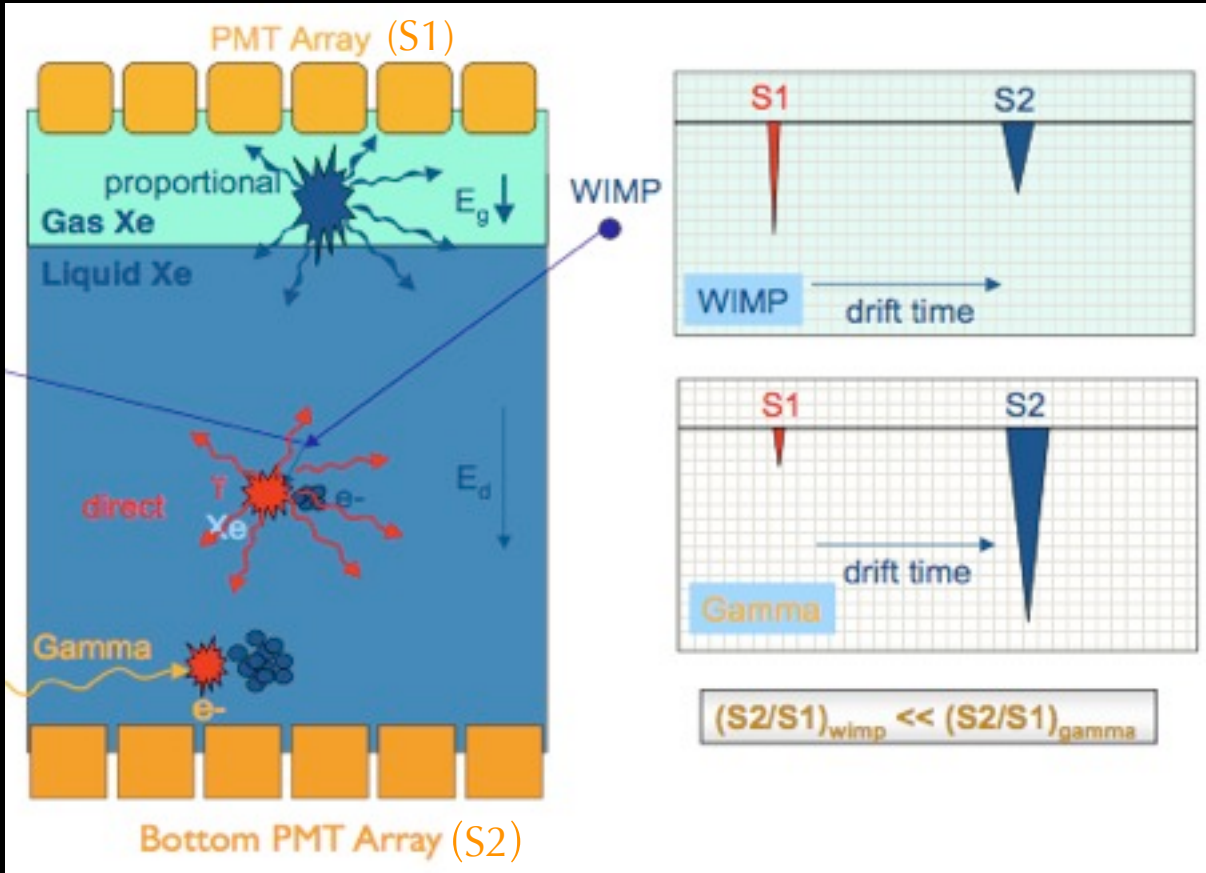
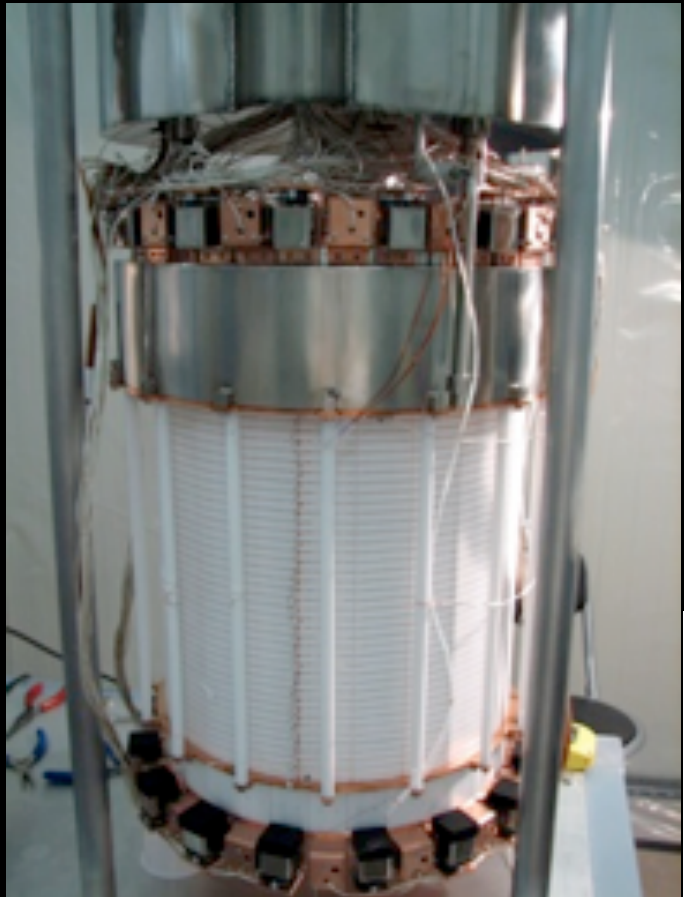
Charge side: 2 concentric electrodes (inner & outer) energy (& veto)



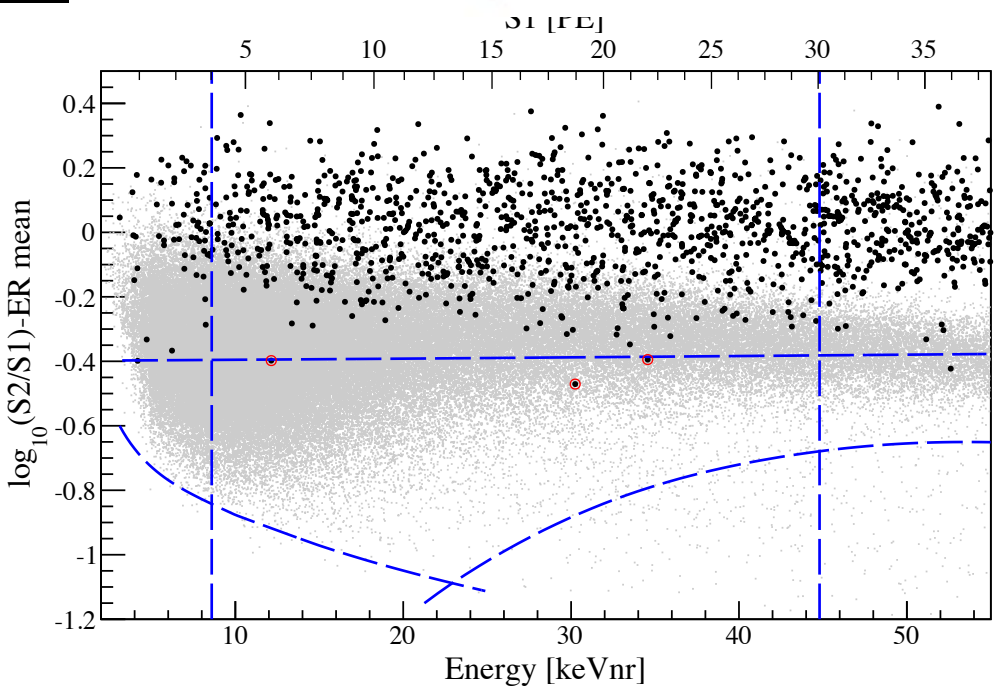
Ionization/Phonon yield vs. E_{recoil} (keV)

Xenon100

Two-Phase liquid Xenon TPC,
operated in LNGS



“S2”: primary scintillation
“S1”: amplified, drifted
ionization signal
both read out with PMTs

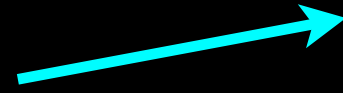


(E. Aprile)

arXiv:
1104.2549v2

Experimental Procedure

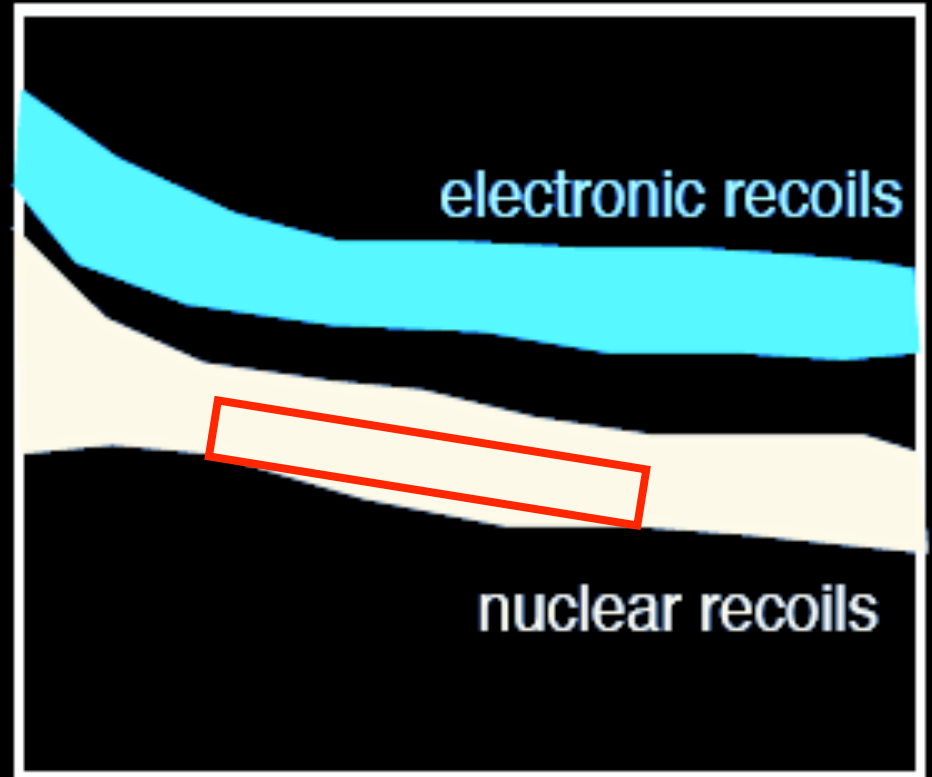
1. using calibration data,
define background
and signal regions



2. using calibration data,
define a region
with zero expected
background events



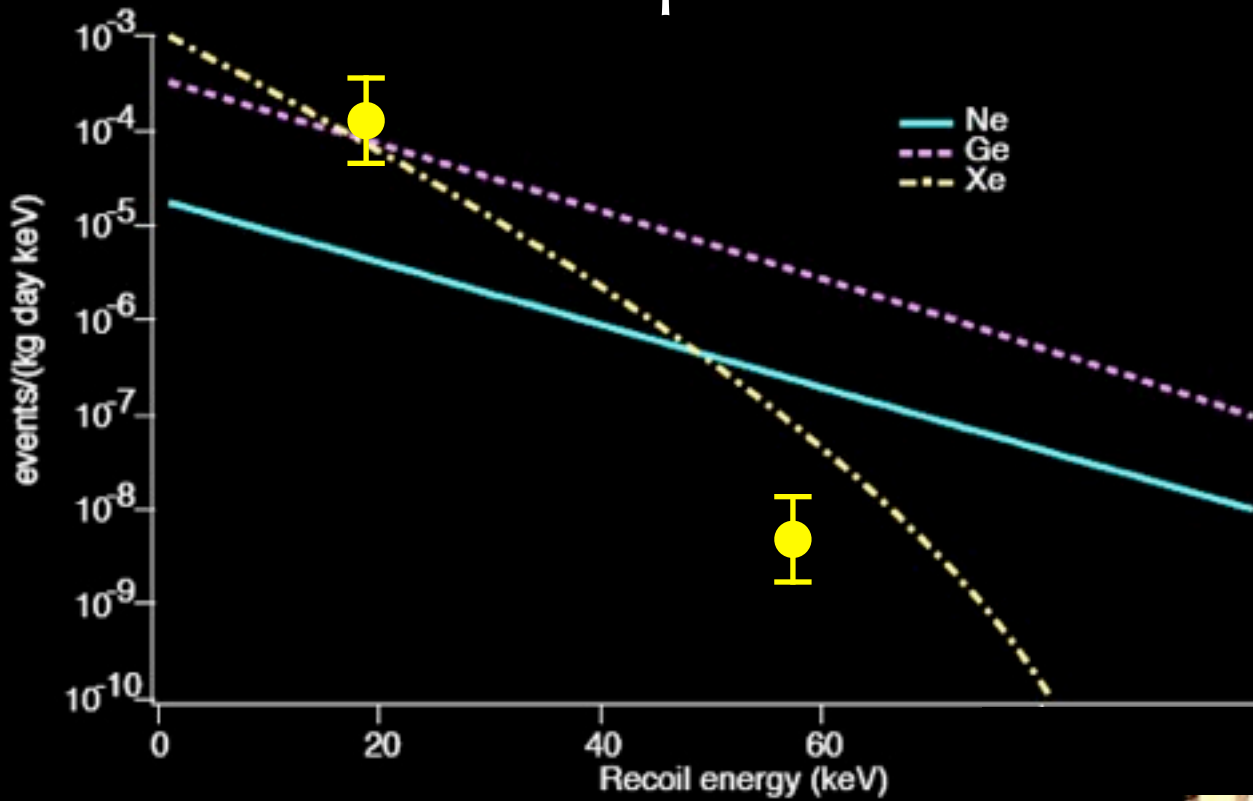
log (ionization/scintillation)



nuclear recoil energy

(any events in the blind region are signal candidates)

Experimental Procedure



theorist:

what model explains this rate?

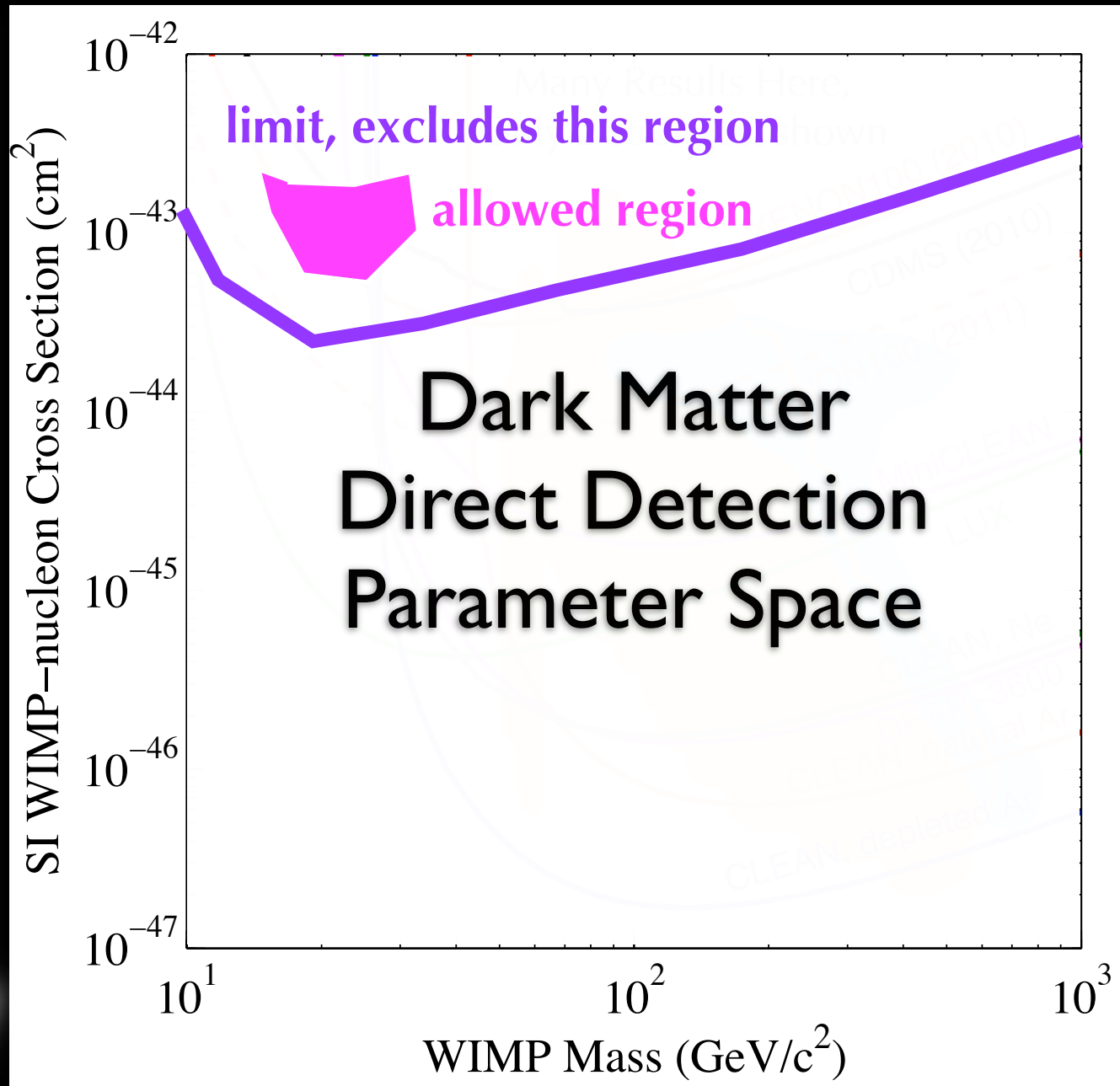
see a signal or set a limit?

experimentalist:



The Low Background Frontier

Scalability of Detector Technology



← 1 event/
kg/day

← 1 event/
100 kg/day

← 1 event/
100 kg/
100 days

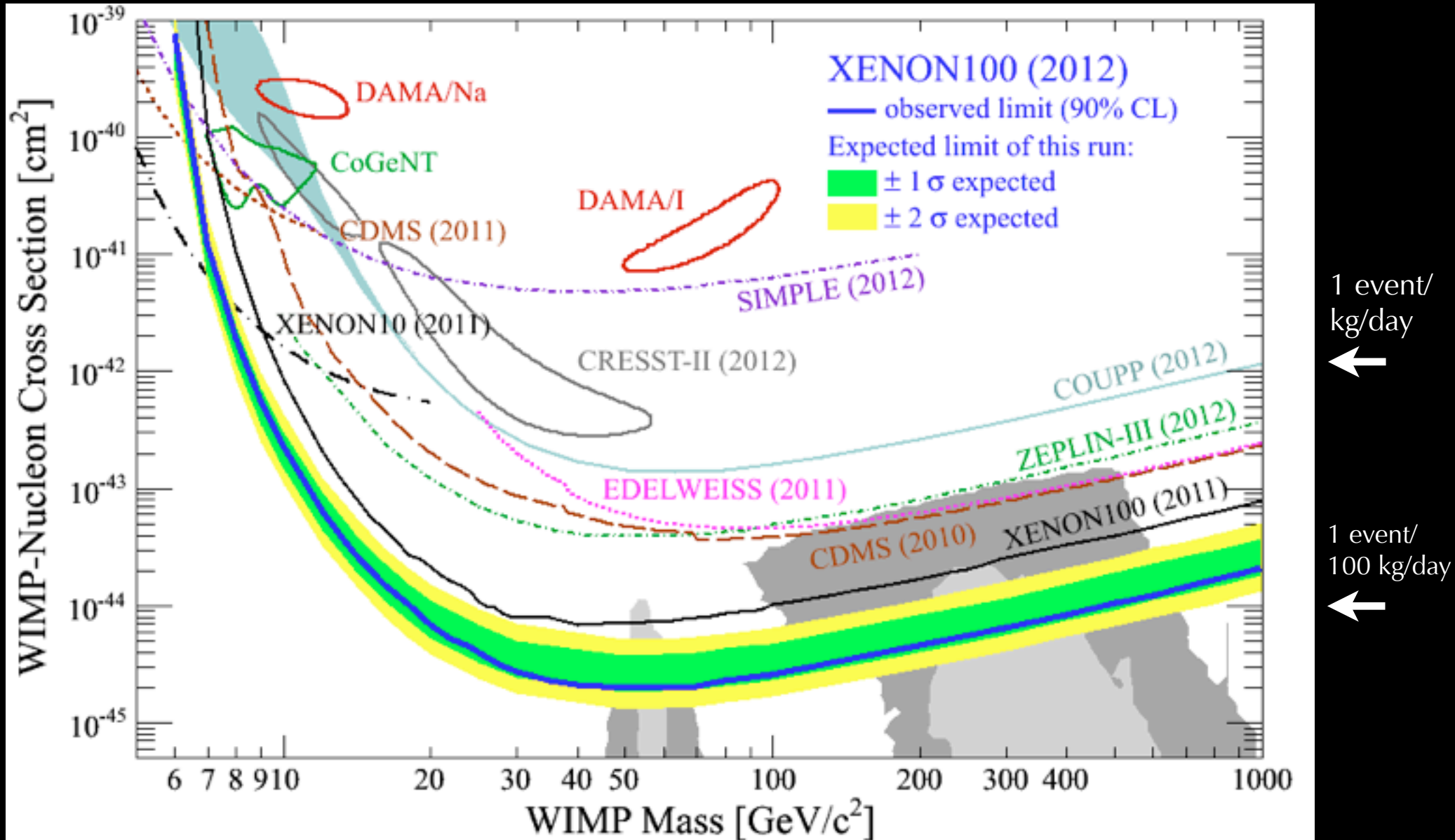
New Techniques for Backgrounds

Complementary with High-Energy Frontier

*need 100-1000
dark matter events
to measure mass,
cross section*

Spin-Independent Cross Section: Latest Experiment Results

E. Aprile et al., arXiv:1207.5988



tonne-scale detectors likely required for discovery

Spin-Independent Cross Section: Latest Theory Results

O. Buchmuller et al., arXiv:1207.7315

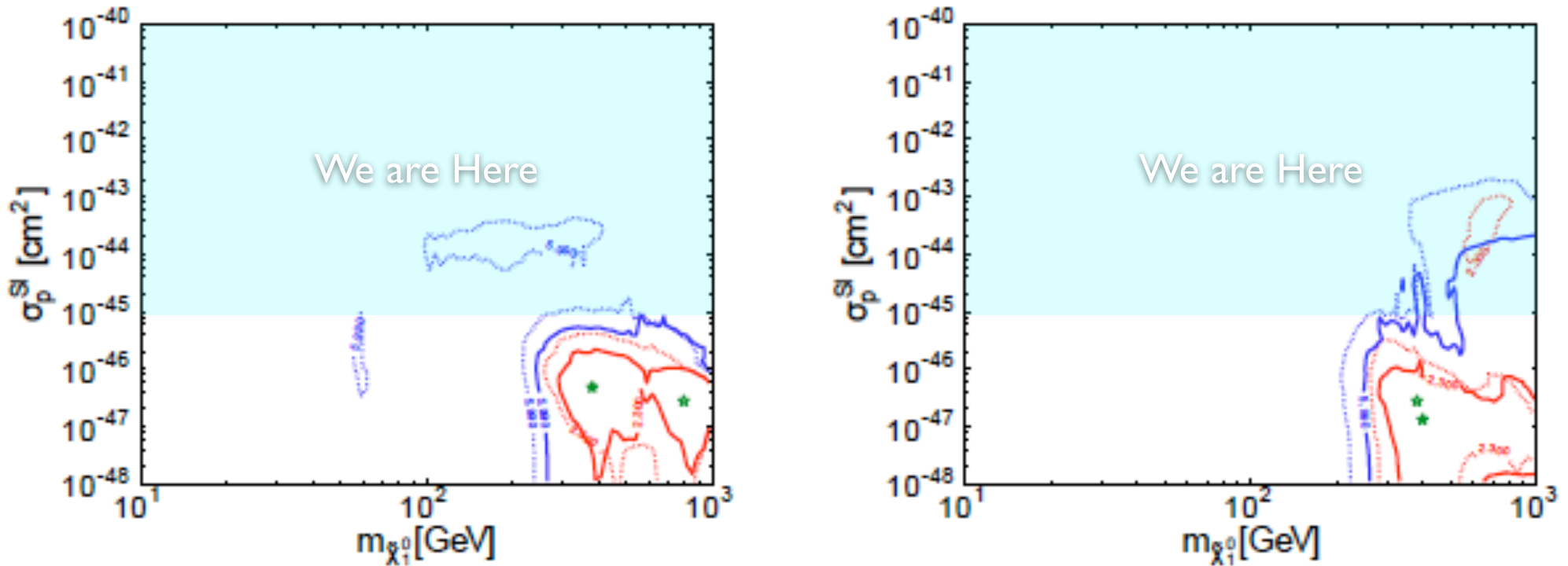
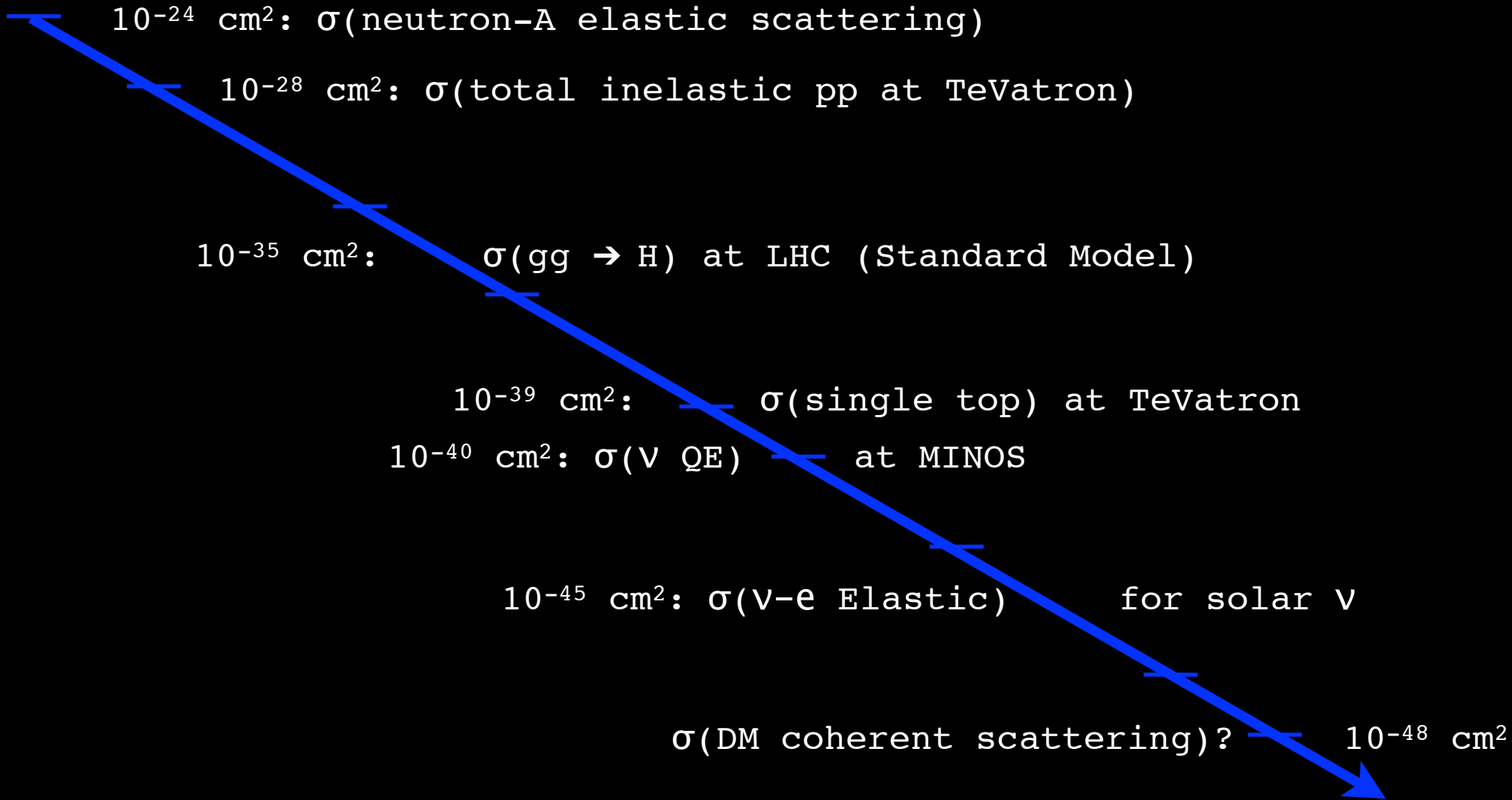


Figure 17. The $(m_{\tilde{\chi}_1^0}, \sigma_p^{\text{SI}})$ planes in the CMSSM (left panel) and the NUHM1 (right panel). The $\Delta\chi^2 = 2.30(5.99)$ contours, corresponding to the 68(95)% CL are coloured red (blue). The solid (dashed) lines are for global fits to the LHC_{5/fb}, BR($B_s \rightarrow \mu^+\mu^-$) and new XENON100 (LHC_{1/fb}) data, and the corresponding best-fit points are indicated by solid (open) green stars.

CAVEAT: many more exotic models (Asymmetric DM, Dark Forces Models, Magnetic Inelastic DM, Sterile Nus + Freeze-In, Isospin-Violating DM, Emergent DM and L# models

10^3 is a lot of σ

Not to Scale

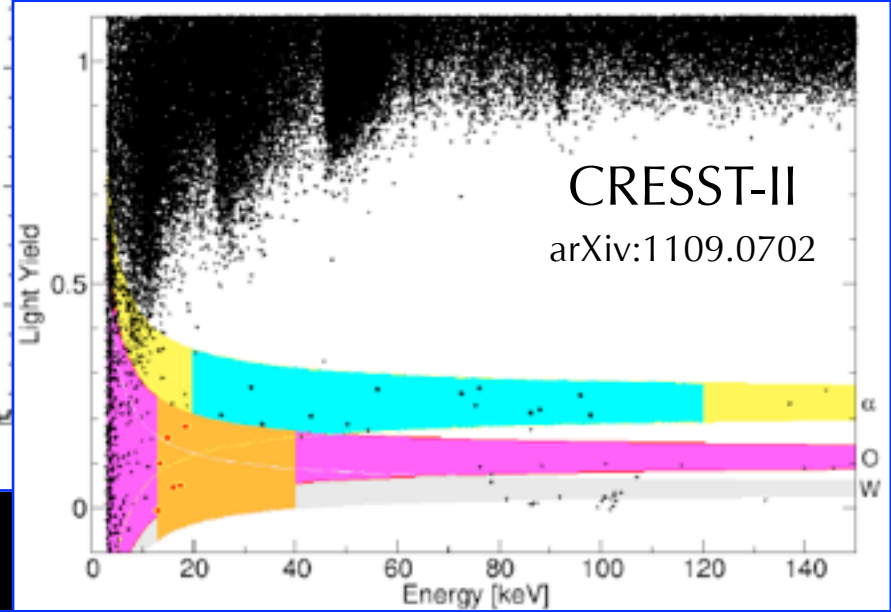
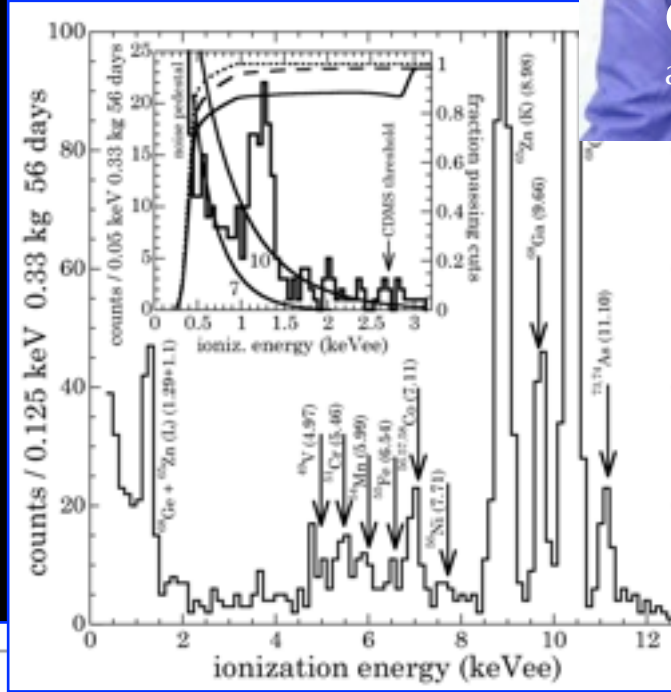


Direct Dark Matter Signals?

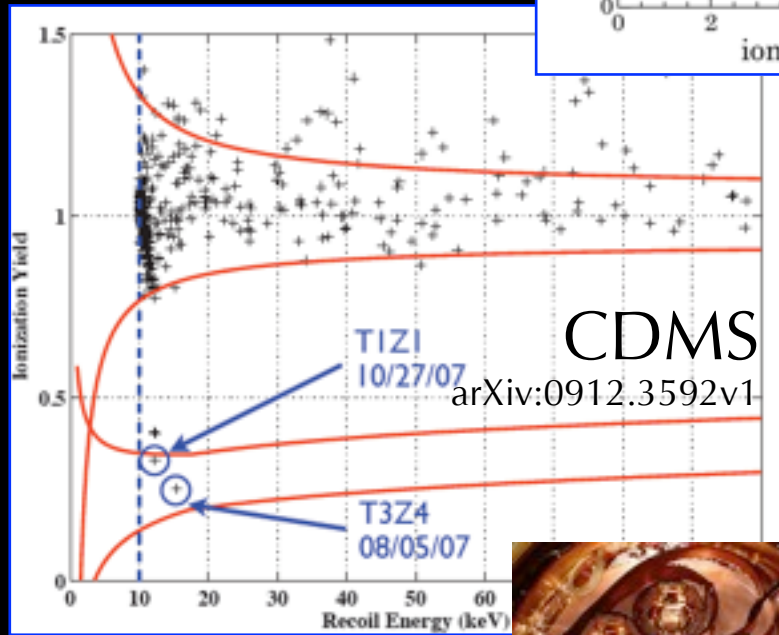
DAMA/Libra



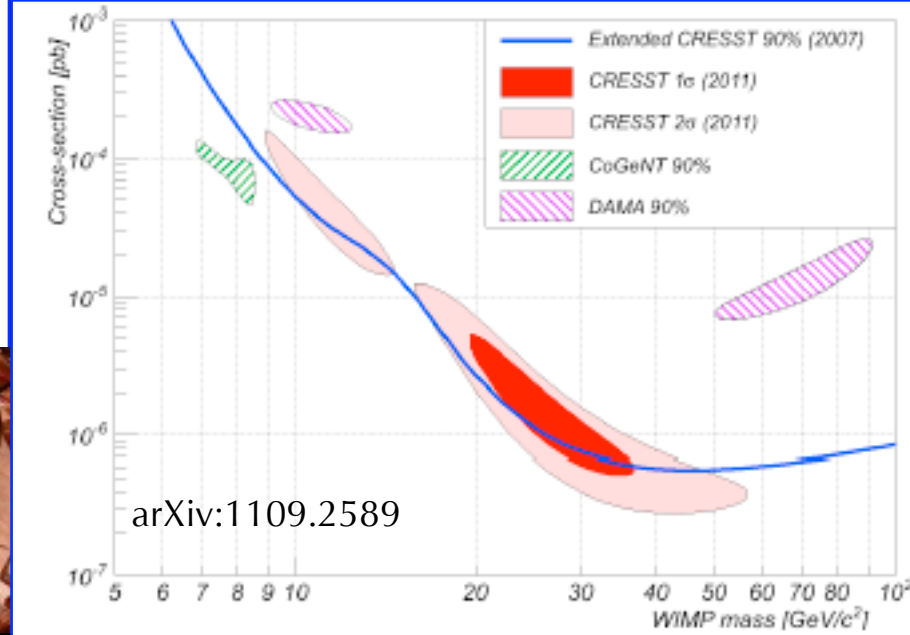
COGENT
arXiv:1002.4703



CRESST-II
arXiv:1109.0702



CDMS
arXiv:0912.3592v1

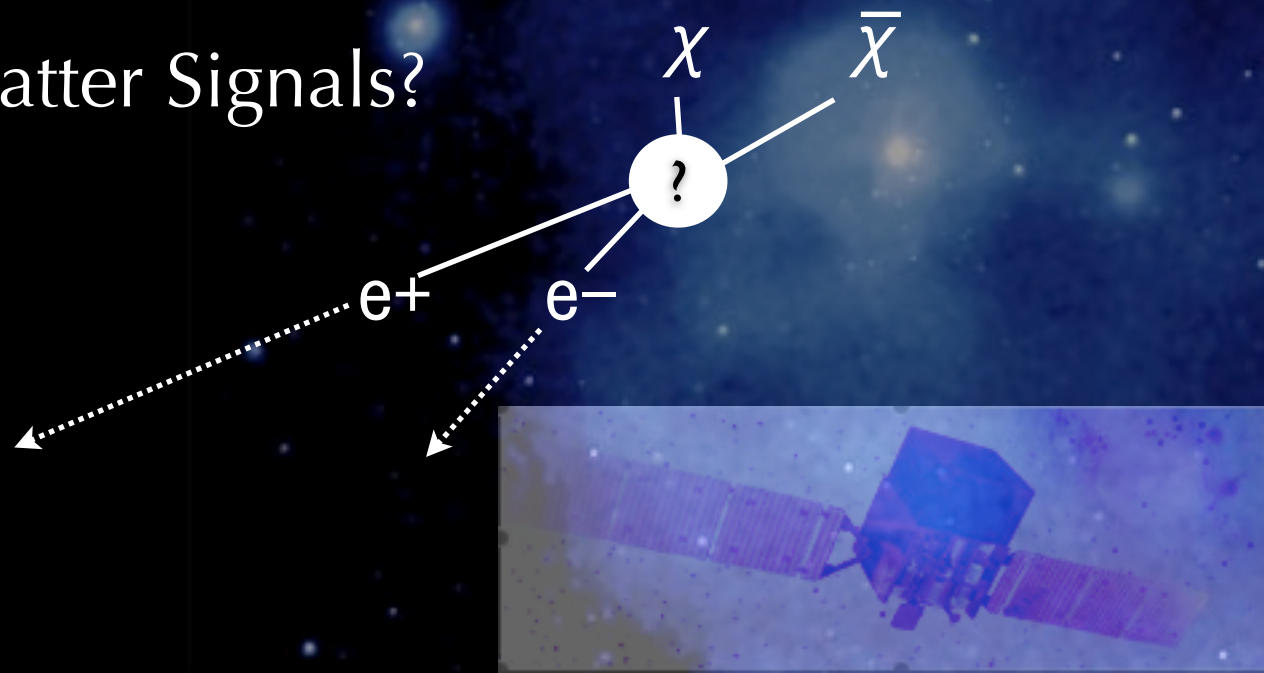
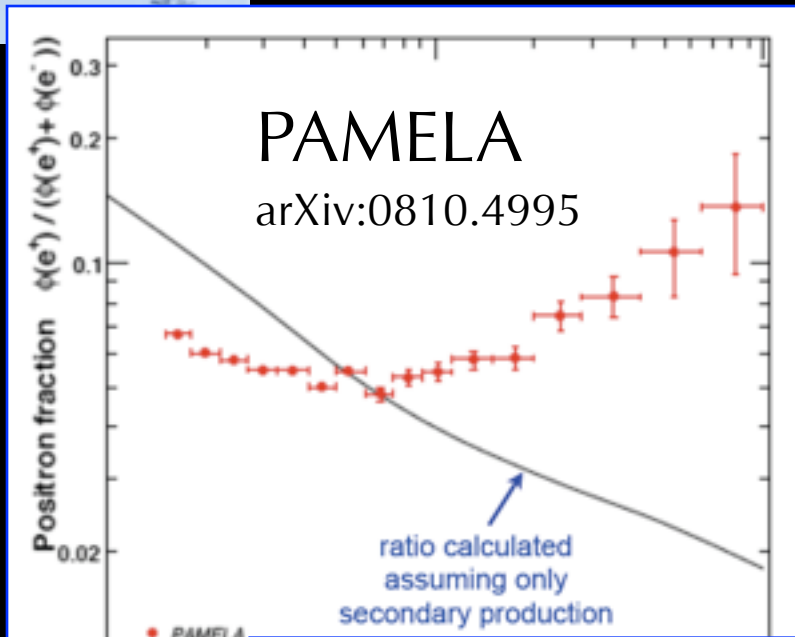


arXiv:1109.2589

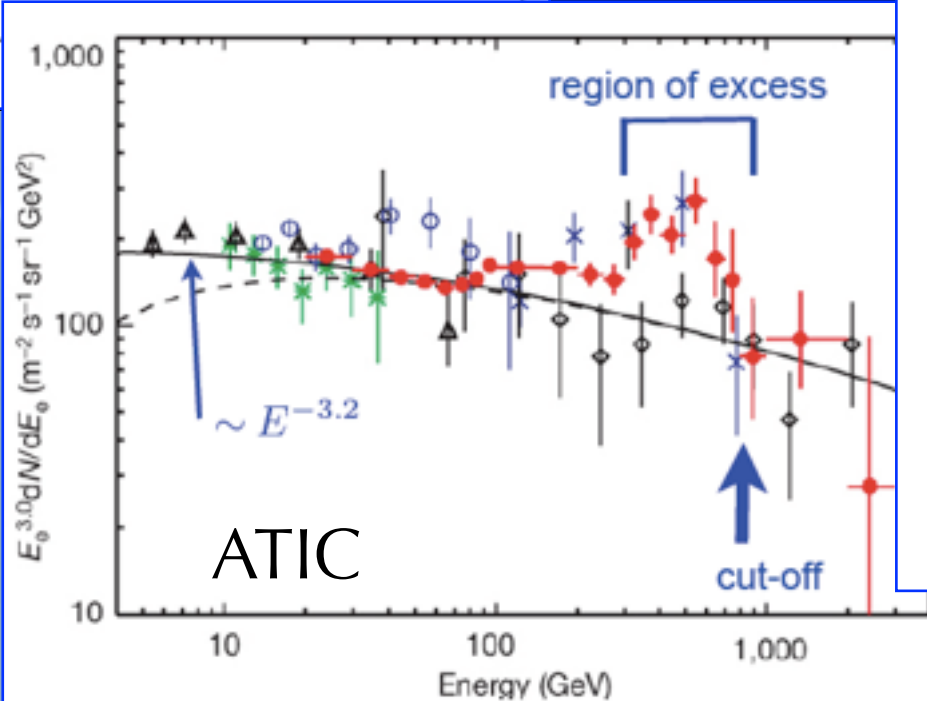
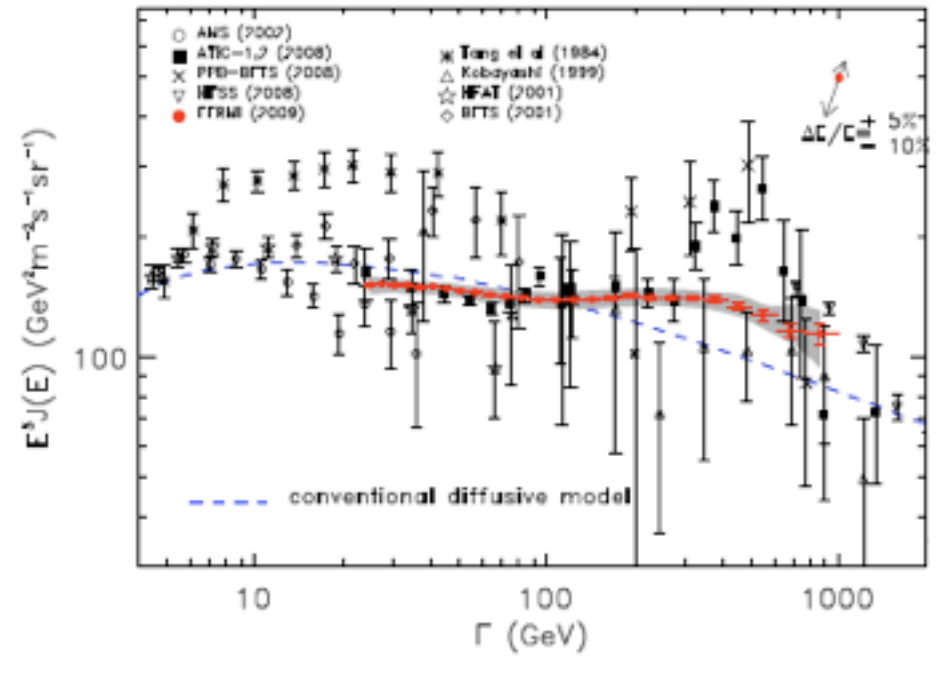


dark matter?
backgrounds?

Indirect Dark Matter Signals?



Fermi LAT arXiv:0905.0025



J. Chang *et al.* Nature **456** 362-365 (2008)

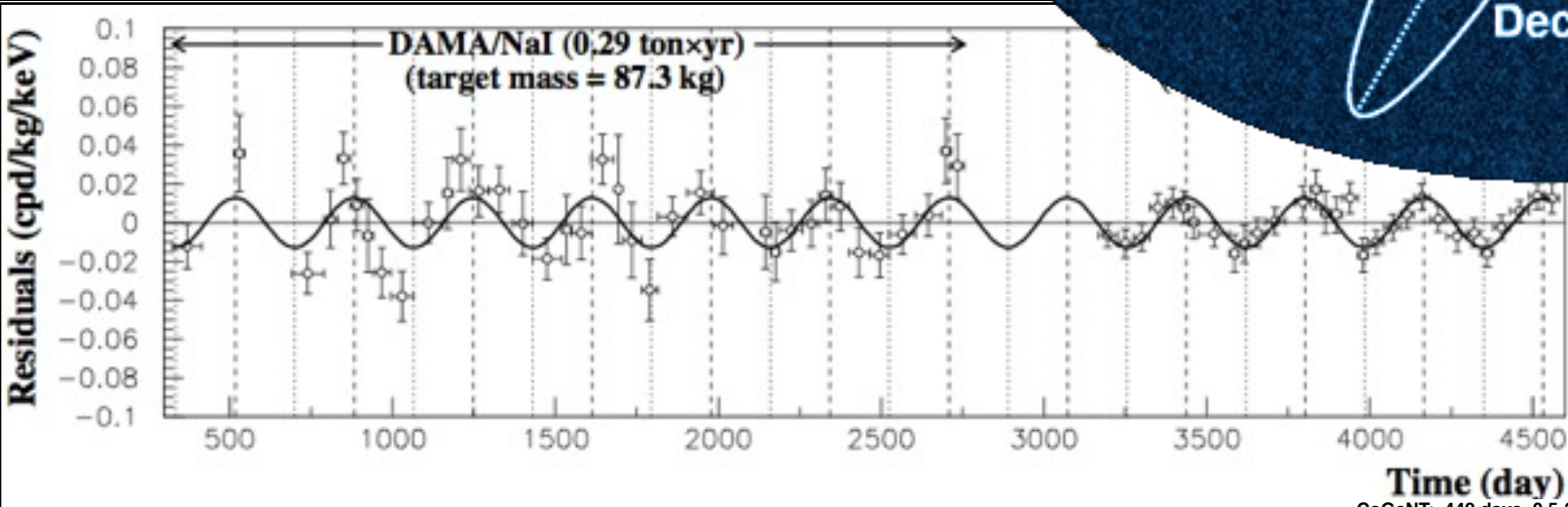
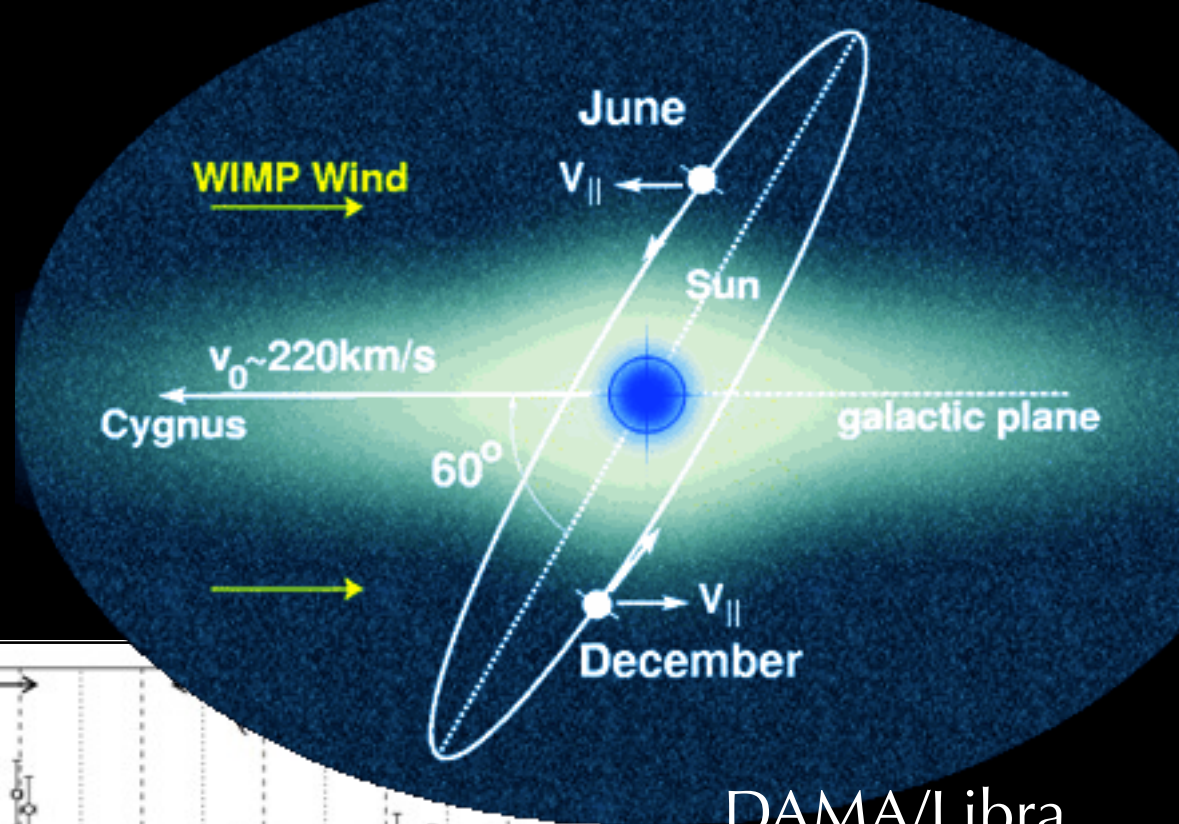
dark matter? local astrophysics?

Annual Modulation

June-December event rate asymmetry $\sim 2-10\%$

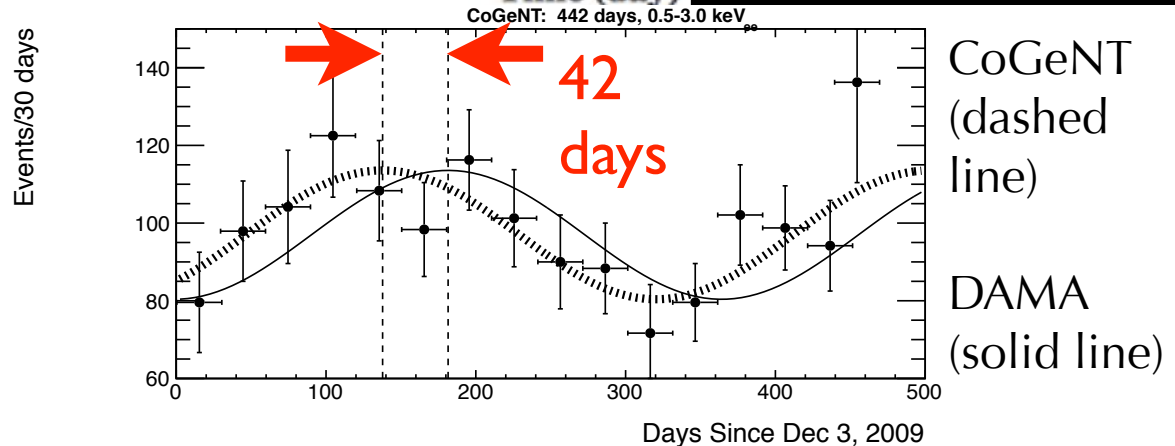
Drukier, Freese, Spergel,
*Phys. Rev. D*33:3495 (1986)

*Eur. Phys. J. C*56:333-355 (2008)



DAMA/Libra
 positive result,
 $>8\sigma$, inconsistent
 with many expts

CoGeNT modulation
 result, 2.8σ , \sim consistent
 with DAMA/Libra
 J. Collar, STSI (2011),
 arXiv:1106.0650v1

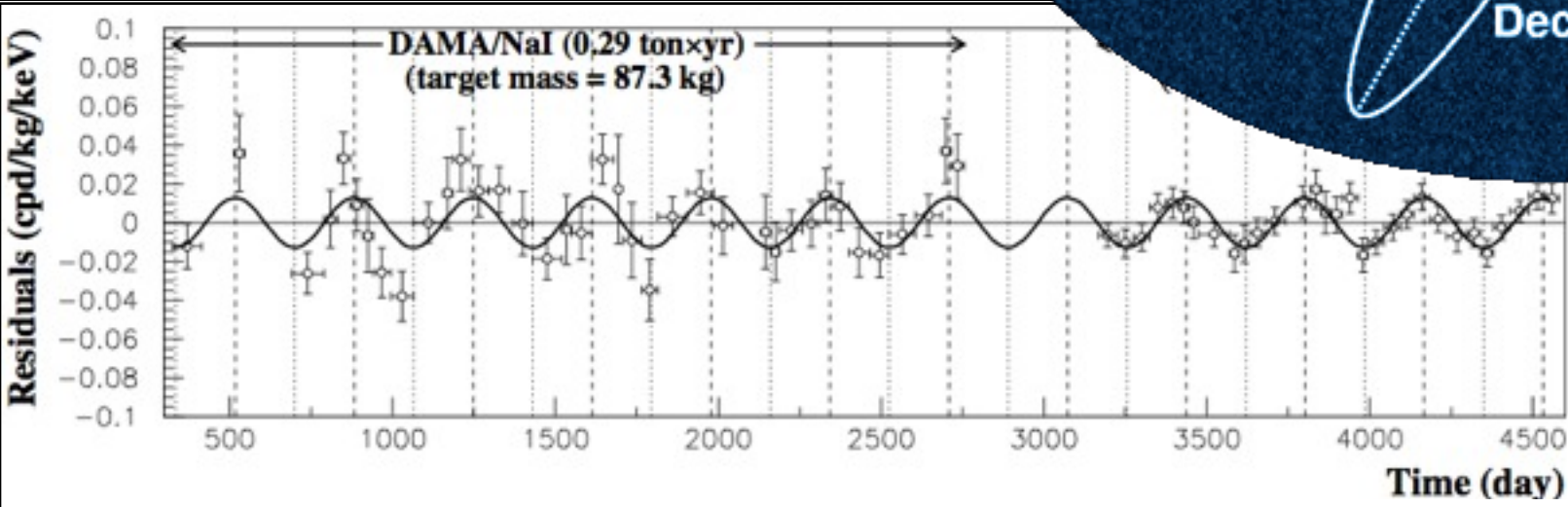
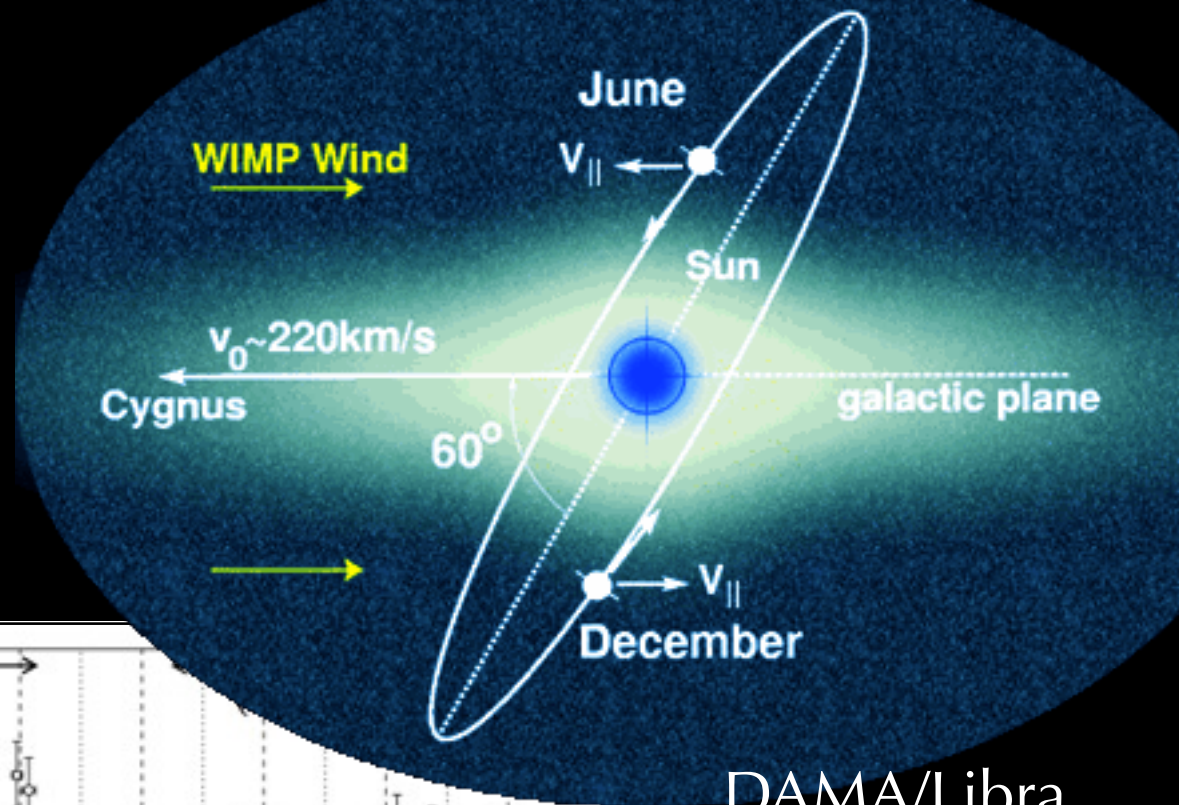


Annual Modulation

June-December event rate asymmetry $\sim 2-10\%$

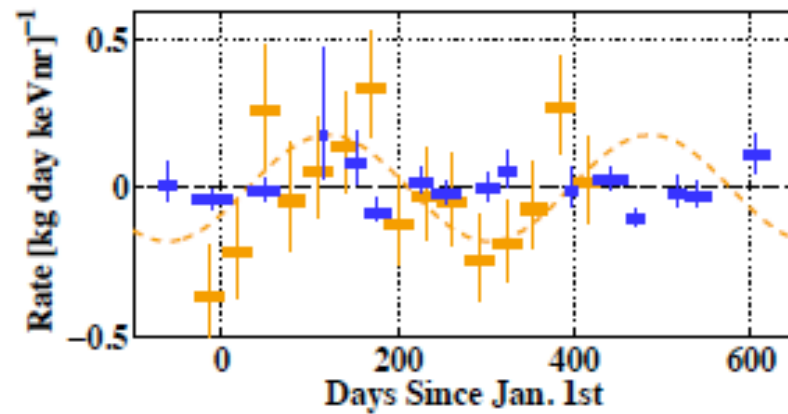
Drukier, Freese, Spergel,
Phys. Rev. D33:3495 (1986)

Eur. Phys. J. C56:333-355 (2008)



DAMA/Libra
positive result,
 $>8\sigma$, inconsistent
with many expts

BUT...
CDMS modulation search
not consistent with CoGeNT
or DAMA/Libra (98.3% CL)
arXiv:1203.1309

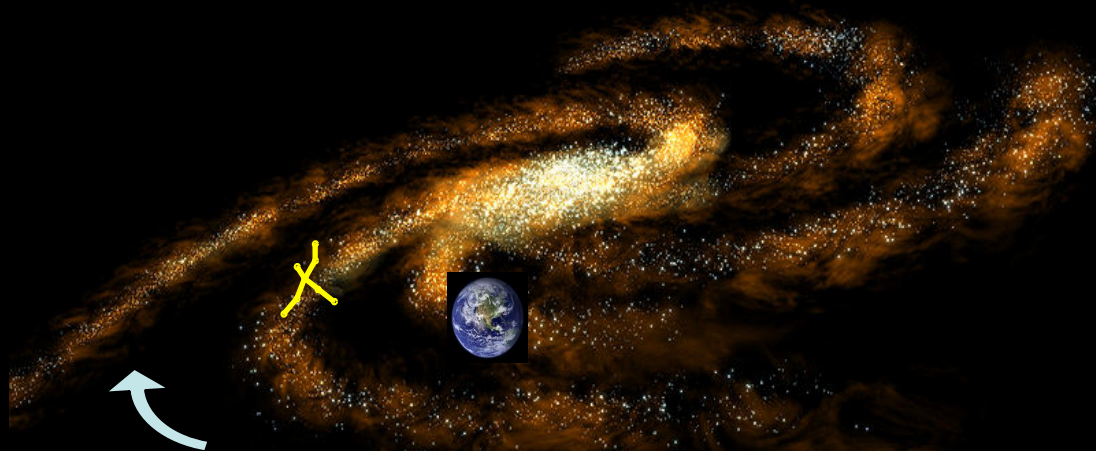


CoGeNT
(dashed line)

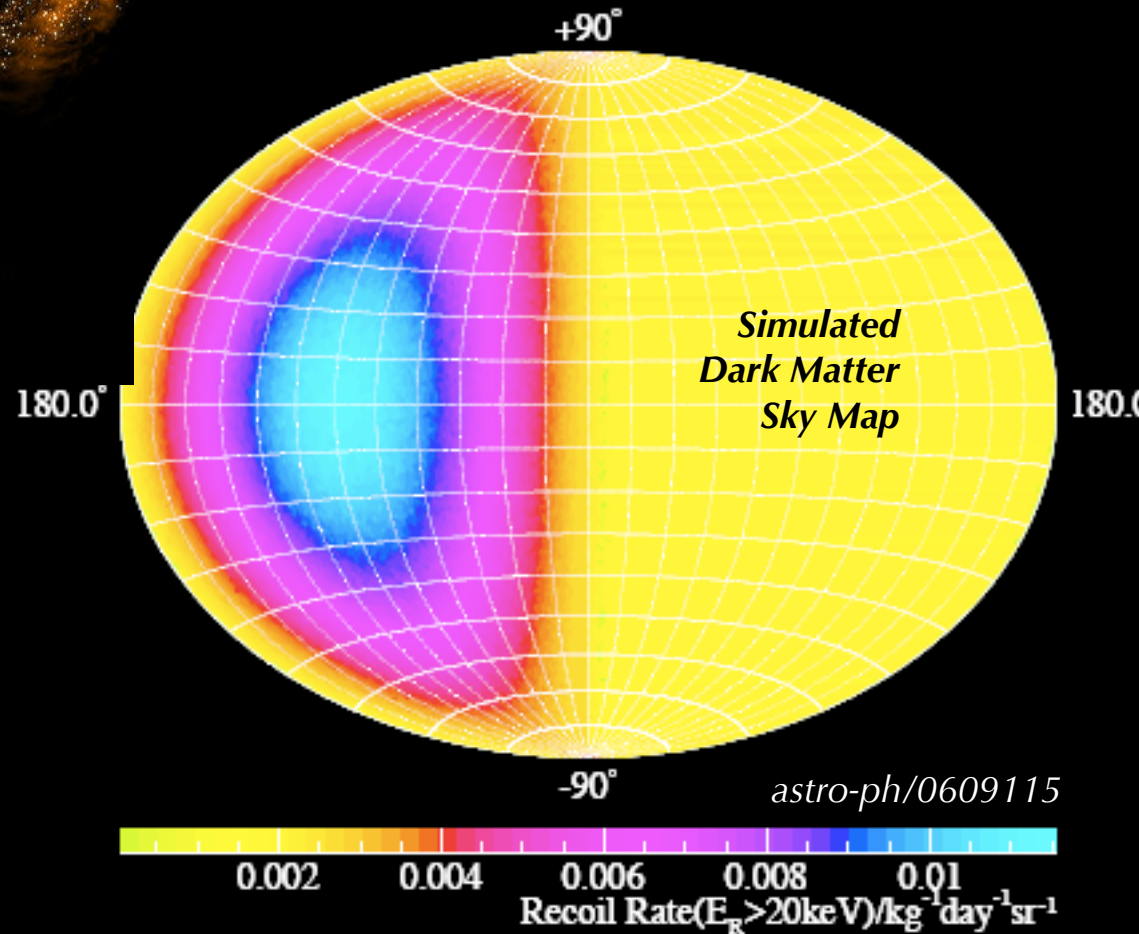
CDMS
(solid points)

Directional Detection

The Dark Matter Wind “blows” from Cygnus



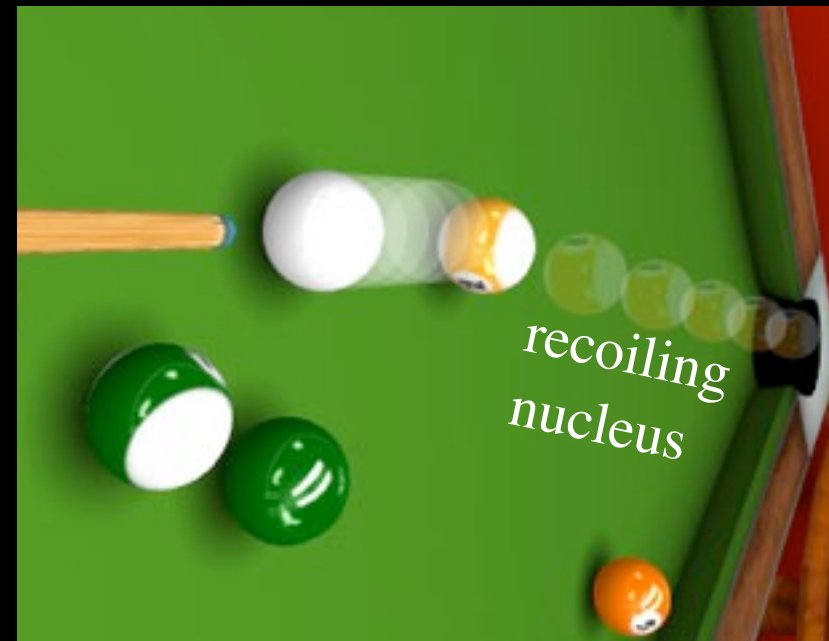
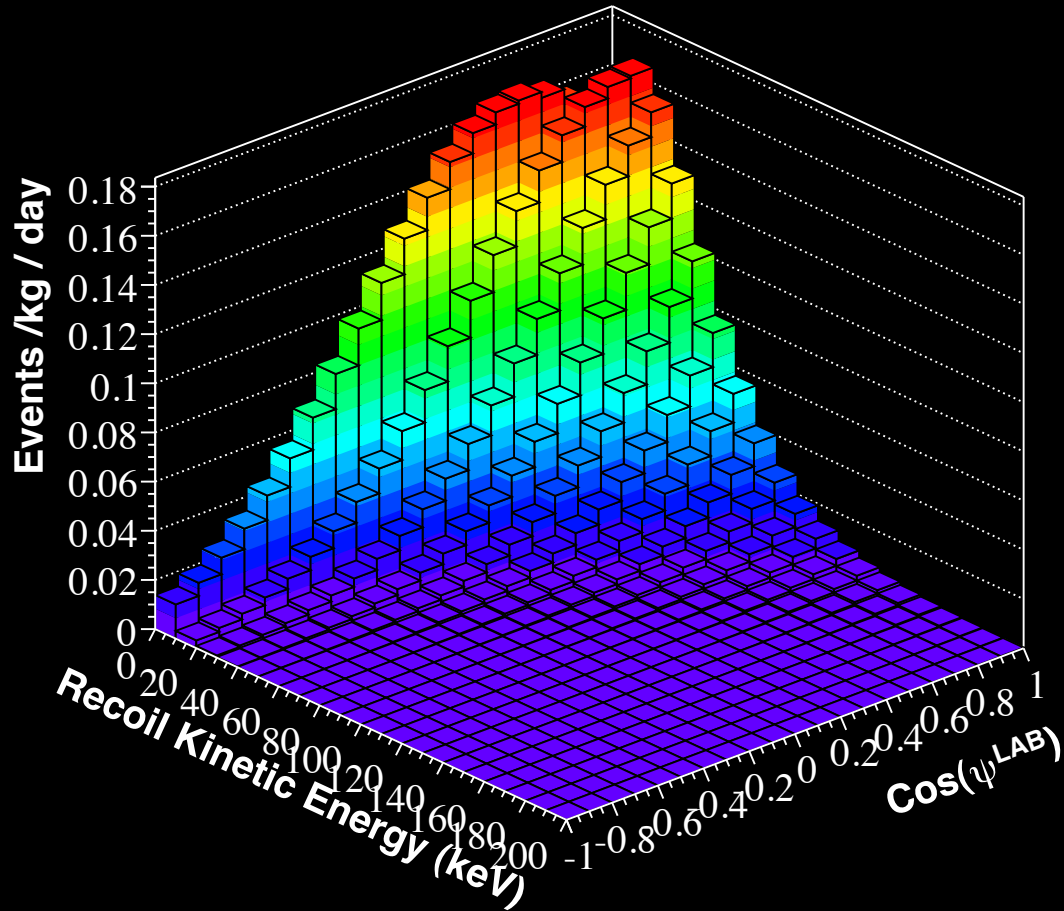
search for a dark matter source



simulated reconstructed dark matter sky map: search for anisotropy

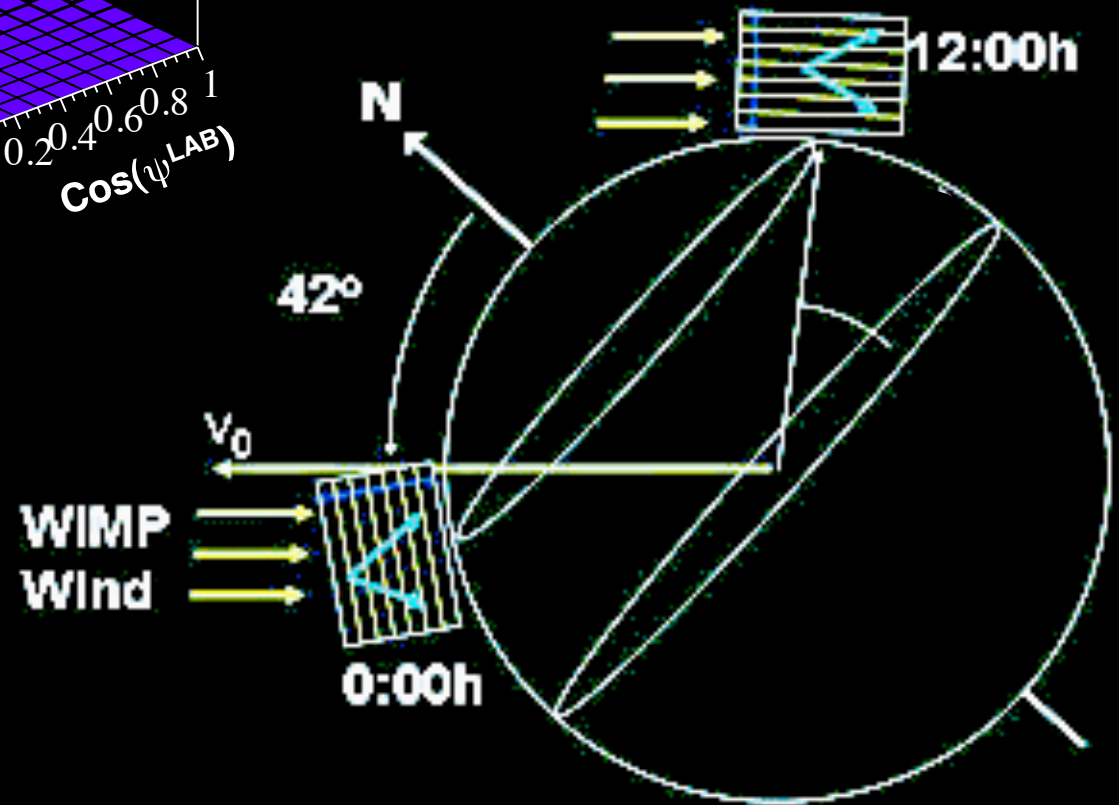
Unambiguous proof:
Correlation of WIMP-induced nuclear recoil signal with galactic motion

Directional Detection



Daily direction modulation:
asymmetry $\sim 20\text{-}100\%$
in forward-backward
event rate.

Spergel, Phys. Rev. D36:1353 (1988)



Directionality Around the World

DRIFT: in Boulby (UK)

S. Burgos et al., Astropart. Phys. 28, 409 (2007)



NEWAGE: in Kamioka (JP),
first directional dark matter limit!

K. Miuchi, et al., Phys.Lett.B654:58-64 (2007)



MiMAC-He3: ILL/Modane (FR)

D. Santos, et al., J. Phys. Conf. Ser. 65, 021012 (2007)

DMTPC: in WIPP (US)
CCD readout

D. Dujmic, et al., NIM A 584:337 (2008)

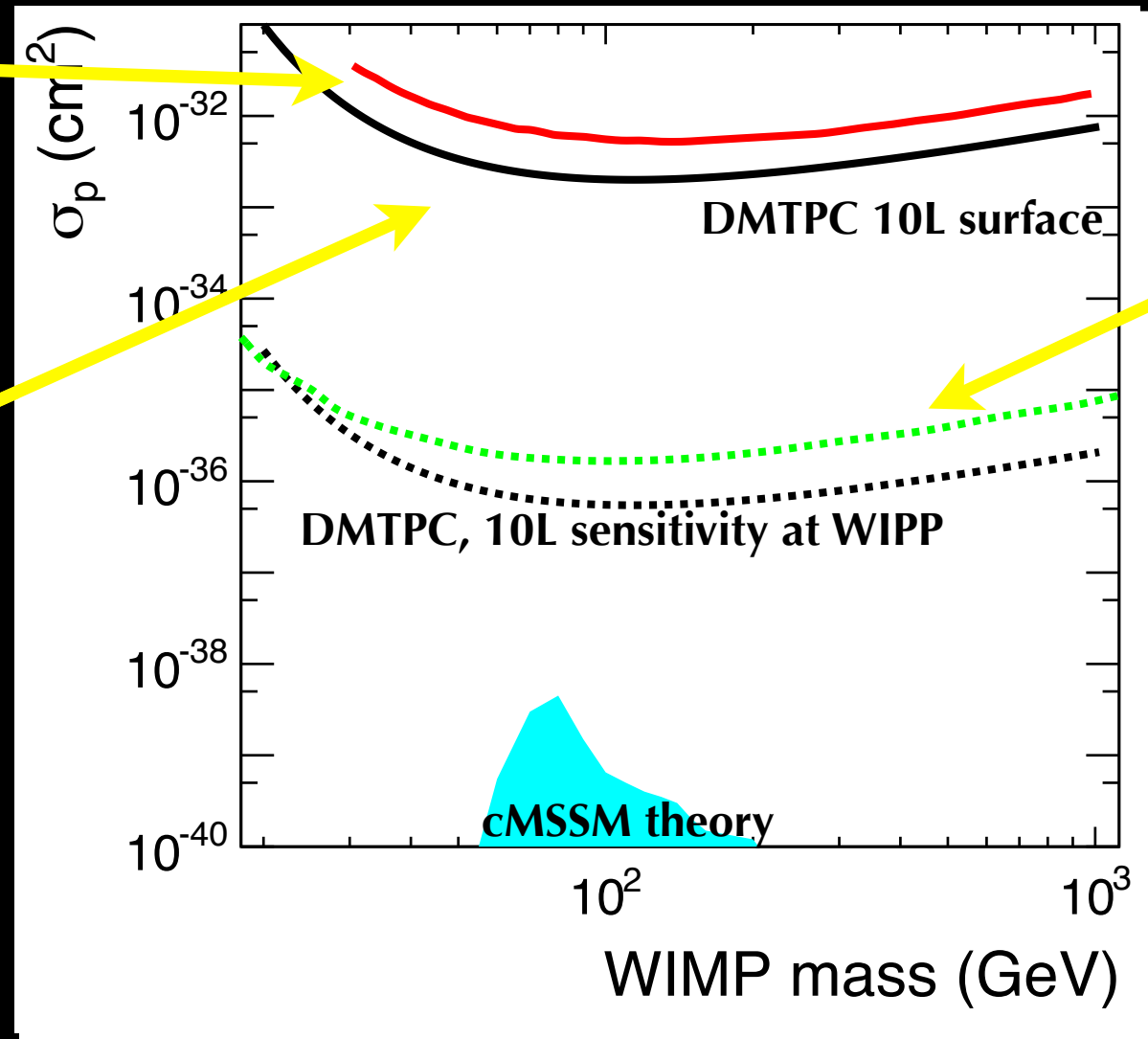
Spin-Dependent Cross Section: Latest Experiment Results

NEWAGE limit
(Kamioka)

*K. Miuchi et al.,
Phys.Lett.B686:11-17
(2010)*

DMTPC limit
(surface, 38 gm-day)

*S. Ahlen et al.,
Phys. Lett. B 695 (2011)*



directional
searches

DRIFT
sensitivity,
Astropart.Phys.
25 (2012) 397

Theory region:

Rozkowski et al JHEP 07 (2007) 075

Ellis et al PRD63 (2001) 065016

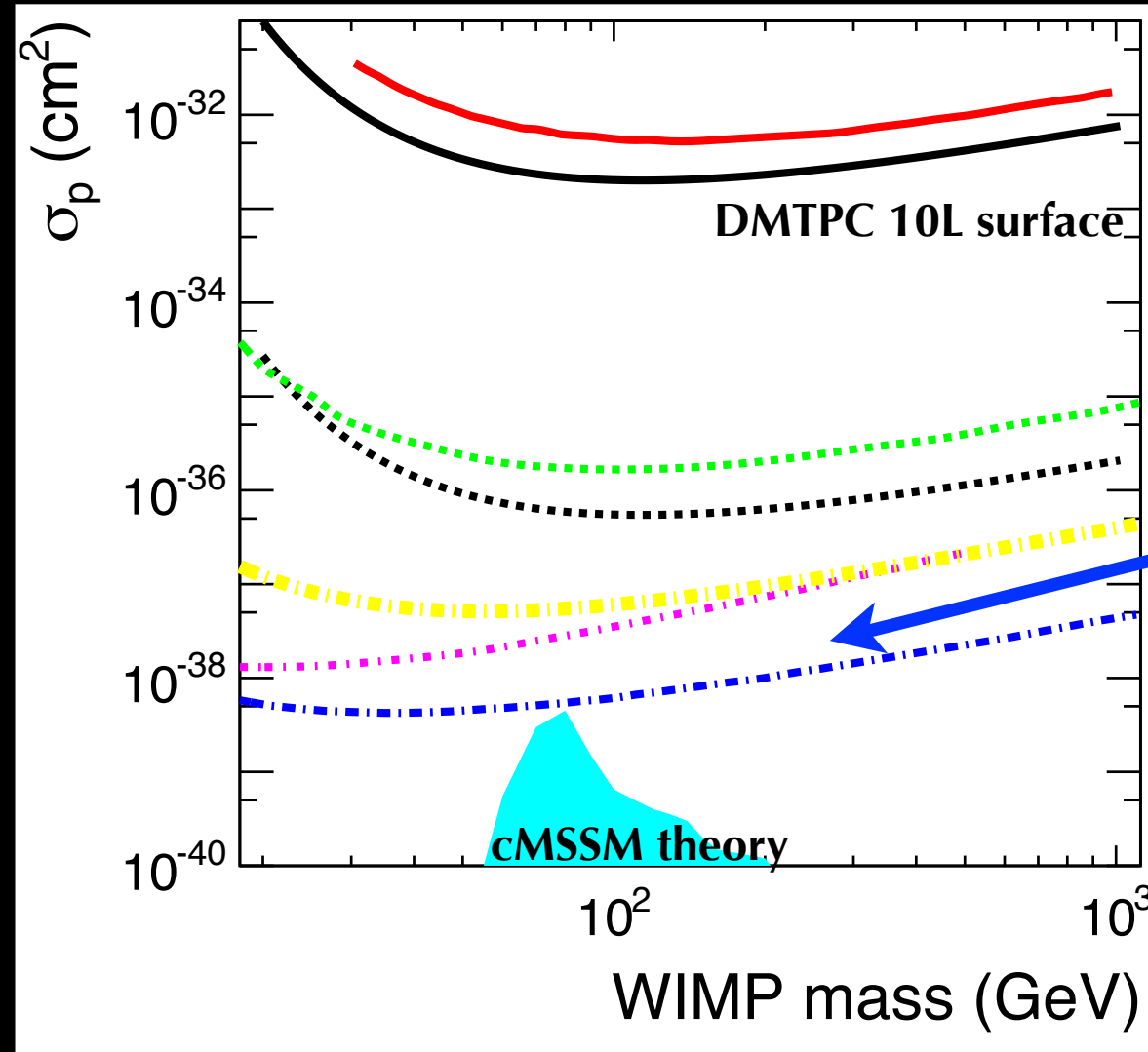
Spin-Dependent Cross Section: Latest Experiment Results

NEWAGE limit
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*directional
searches*

DRIFT
*sensitivity,
Astropart.Phys.
25 (2012) 397*

1D results

COUPP
*PRL.106, 021303
(2011)*

SIMPLE
arXiv:1106.3014

PICASSO
arXiv:1202.1240

Theory region:

Rozkowski et al JHEP 07 (2007) 075

Ellis et al PRD63 (2001) 065016

Spin-Dependent Cross Section: Latest Experiment Results

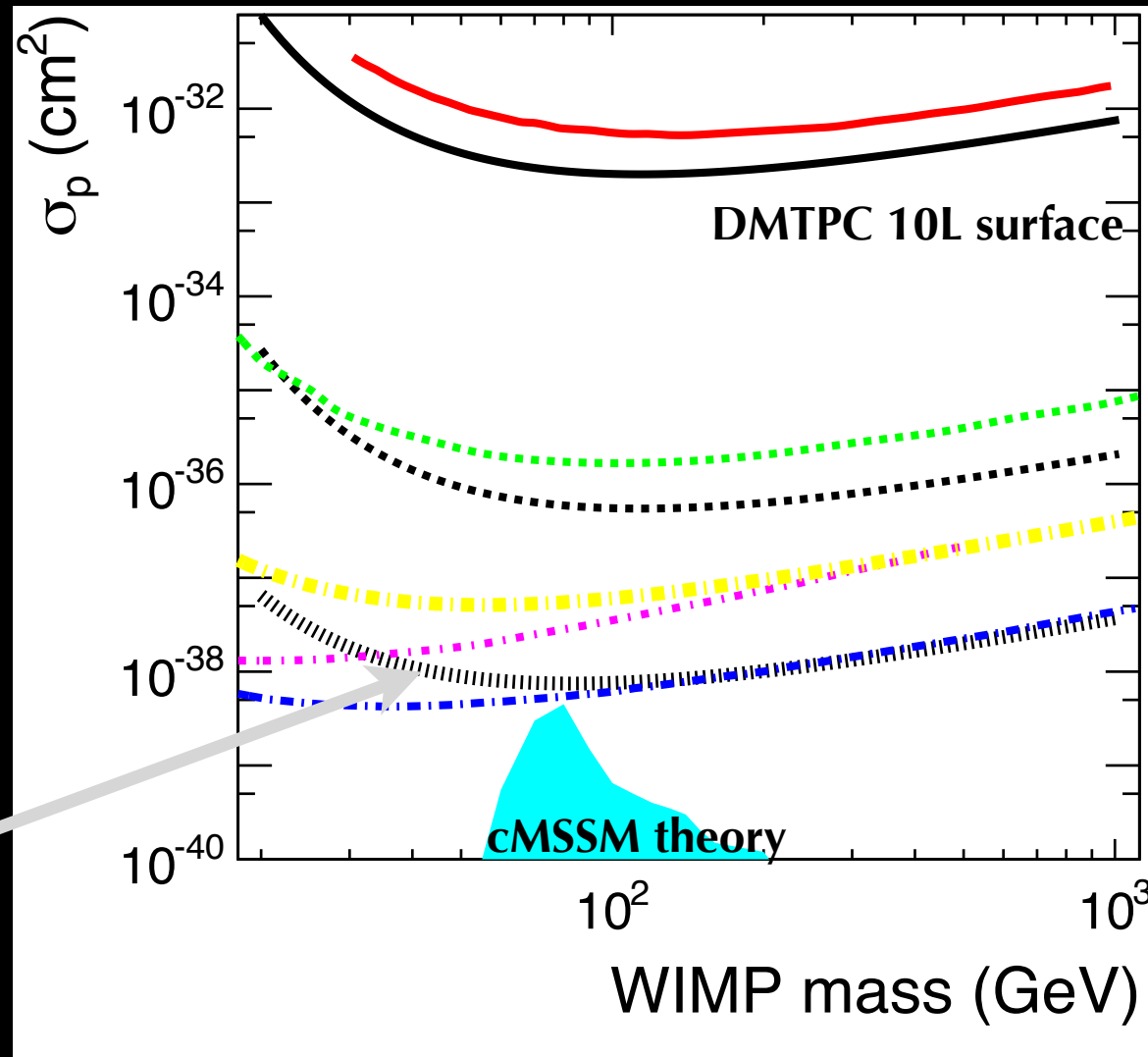
NEWAGE limit
(Kamioka)

*K. Miuchi et al.,
Phys.Lett.B686:11-17
(2010)*

DMTPC limit
(surface, 38 gm-day)

*S. Ahlen et al.,
Phys. Lett. B 695 (2011)*

1 m³ at WIPP
(DMTPCino)
projected
sensitivity



*directional
searches*

DRIFT
sensitivity,
Astropart.Phys.
25 (2012) 397

1D results

COUPP
*PRL.106, 021303
(2011)*

SIMPLE
arXiv:1106.3014

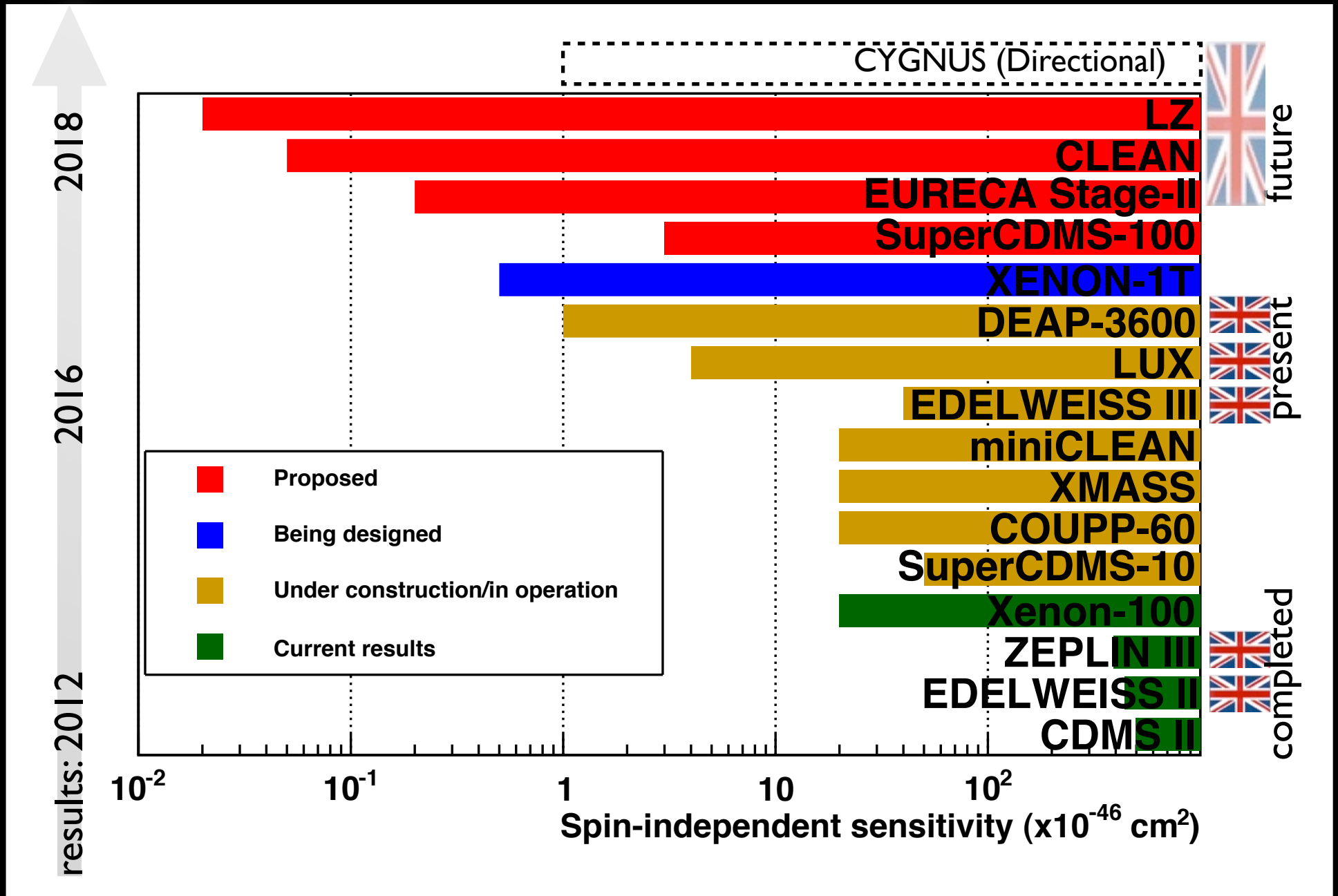
PICASSO
arXiv:1202.1240

Theory region:

Rozkowski et al JHEP 07 (2007) 075

Ellis et al PRD63 (2001) 065016

Global SI Dark Matter Programme, Sensitivity Reach



NB: projected sensitivities (all except green) assume zero background.

Backup Slides

Setting a Limit

1. The theoretical dark matter interaction rate is:

$$\frac{dR}{dE_R} = \left(\frac{c_1 R_0}{E_0 r} \right) \exp\left(\frac{-c_2 E_R}{E_0 r} \right) \quad \begin{array}{l} E_R = \text{nuclear recoil energy,} \\ E_0 = \text{dark matter particle energy} \end{array}$$

2. Experiments measure:

$$R_0 = \left[\left(\frac{2v_0}{\sqrt{\pi}} \right) \left(\frac{N_0(\rho_D/m_D)}{A} \right) \right] \sigma_0 \times \text{exposure}$$

$$\sigma_A = \sigma_0 F^2(E_R, A) I_C, \quad F^2(E_R, A) = \text{nuclear form factor}, \quad I_C = A^2$$

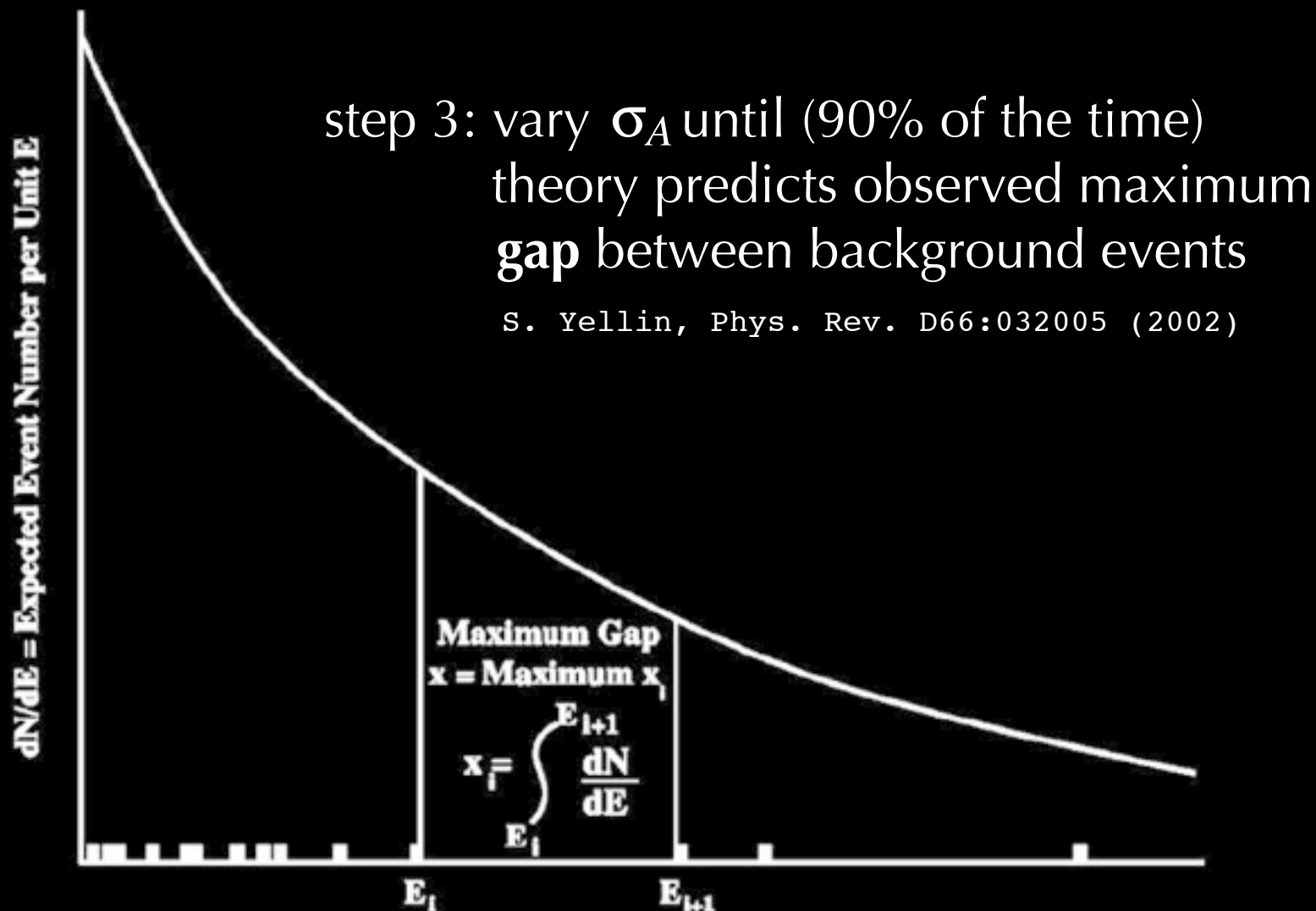
3. vary σ_A until (90% of the time) theory predicts observed rate

4. Normalize to σ_{W-N} to compare limits:

$$\sigma_{W-N} = \left(\frac{\mu_1}{\mu_A} \right)^2 \left(\frac{1}{A} \right)^2 \sigma_A$$

$$\mu = \frac{m_D m_{\text{target}}}{(m_D + m_{\text{target}})}$$

... in the Presence of Background



Yellin gap method: a way to make a “zero-background” measurement over a restricted range of an experiment’s acceptance (zero signal too)

$$\sigma = \sigma_{\text{SI}} + \sigma_{\text{SD}};$$

$$\sigma_{0,\text{SI}} = \frac{4\mu^2}{\pi} [Zf_p + (A - Z)f_n]^2.$$

$$\sigma_{0,\text{SI}} = \sigma_{\text{p,SI}} \left(\frac{\mu}{\mu_p} \right)^2 A^2$$

$$\sigma = \sigma_0 |F(E)|^2 \quad (\text{nuclear form factor})$$

$$\sigma_{\text{SD}}(q) = \frac{32\mu^2 G_F^2}{2J + 1} [a_p^2 S_{pp}(q) + a_p a_n S_{pn}(q) + a_n^2 S_{nn}(q)].$$

Many of the parameters that factor into the expected recoil rates for a scattering detector are unknown, including the WIMP mass, four WIMP-nucleon couplings (SI and SD couplings to each of protons and neutrons), the local WIMP density, and the WIMP velocity distribution in the halo. In this paper, we shall fix the halo model to the SHM and the local density to 0.3 GeV/cm^3 . In addition, we shall take $f_p = f_n$ (equal SI couplings) so that there are only three independent scattering couplings; the SI coupling will be given in terms of the SI scattering cross-section off the proton, $\sigma_{\text{p,SI}}$. The parameter space we examine will then consist of the four parameters m , $\sigma_{\text{p,SI}}$, a_p , and a_n .

Scale Parameters

Table 5.2 Energy density and scale parameters for different regimes

Dominant regime	Equation of state	Energy density	Scale parameter
Radiation	$P = \frac{\rho c^2}{3}$	$\rho \propto R^{-4} \propto t^{-2}$	$R \propto t^{1/2}$
Matter	$P = \left(\frac{2}{3}\right) \rho c^2 \times \left(\frac{v^2}{c^2}\right)$	$\rho \propto R^{-3} \propto t^{-2}$	$R \propto t^{2/3}$
Vacuum	$P = -\rho c^2$	$\rho = \text{constant}$	$R \propto \exp(\alpha t)$

SN1A Data

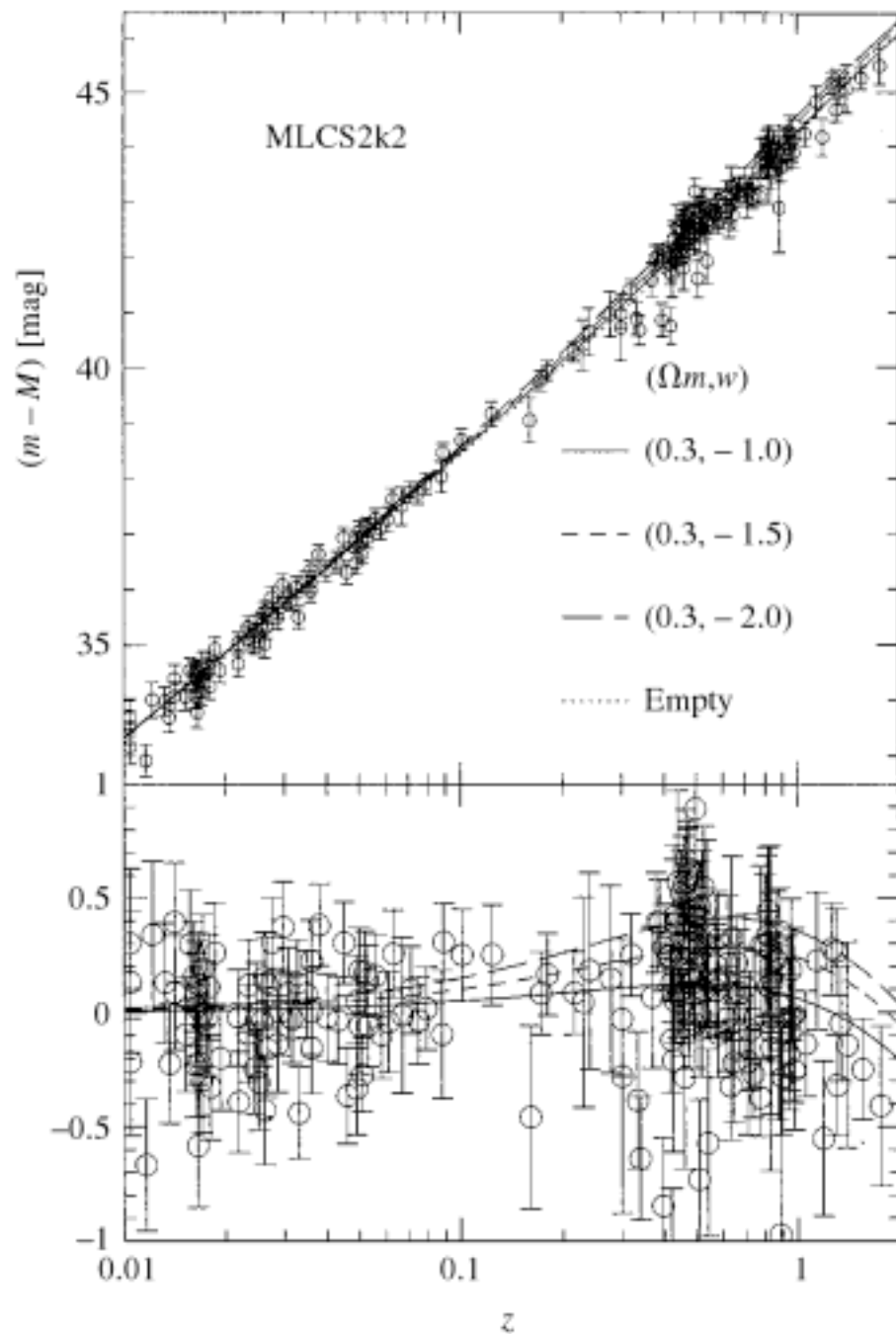


Fig. 7.13 Hubble plot from Type Ia supernovae at low and high redshifts, after Clacchiatti *et al.* (2006). The upper panel shows the measured values of the distance modulus (or logarithm of the luminosity distance) plotted against redshift. The lower panel shows the *difference* in magnitude as compared with the value expected for an empty universe. For averaged values, see Fig. 7.14.

Fig. 7.14 Differential Hubble plot from Type Ia supernovae, after Riess *et al.* (2004). The experimental points represent averages over several supernovae. An empty universe ($\Omega_k = 1, \Omega_m = \Omega_\Lambda = 0$) is represented by the horizontal dotted line. A flat, matter-dominated (so-called Einstein-de Sitter) universe ($\Omega_m = 1, \Omega_\Lambda = \Omega_k = 0$) is shown by the solid curve; while the dashed curve represents their best fit, with $\Omega_m = 0.27$,

