

DIRAC EQUATION

$$E\psi = (\alpha_i p_i + \beta m)\psi$$

$$E^2\psi = (\alpha_i p_i + \beta m)(\alpha_i p_i + \beta m)\psi$$

$$(\alpha_i \alpha_j p_i p_j + (\alpha_i \beta + \beta \alpha_j) p_i m + \beta^2 m^2)\psi$$

$$\text{BUT } E^2 = p^2 + m^2$$

$$\rightarrow \alpha_i \alpha_j + \alpha_j \alpha_i = 2\delta_{ij}$$

$$\alpha_i \beta + \beta \alpha_j = 0$$

$$\beta^2 = \underline{1}$$

ALSO WE WANT

α AND β ARE HERMITIAN

$$\alpha = \alpha^\dagger \quad \beta = \beta^\dagger$$

$$\alpha^2 = \underline{1}$$

FROM $\alpha_i \beta + \beta \alpha_j = 0$, THEY MUST BE TRACELESS :

$$\alpha_i \beta = \beta \alpha_j$$

MULTIPLY BY β

$$\alpha_i \beta^2 = -\beta \alpha_j \beta$$

$$\alpha_i = -\beta \alpha_j \beta$$

$$\text{Tr}(\alpha_i) = -\text{Tr}(\beta \alpha_j \beta)$$

$$\text{BUT } \text{Tr}(\alpha \beta \alpha \beta \dots) = \text{Tr}(\alpha \alpha \beta \beta \dots)$$

$$\rightarrow \text{Tr}(\alpha_i) = -\text{Tr}(\beta^2 \alpha_j)$$

$$\text{FOR } i=j \quad = +\text{Tr}(\alpha_j) = -\text{Tr}(\alpha_i)$$

$$\rightarrow \text{Tr}(\alpha_i) = 0$$

DEFINE

$$\alpha_i = \begin{pmatrix} 0 & \sigma_i \\ \sigma_i & 0 \end{pmatrix}$$

$$\beta = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

COVARIANT FORM

$$E\psi = (\alpha_i \cdot \underline{p} + \beta m)\psi$$

REPLACE ENERGY AND
MOMENTUM BY OPERATORS

$$p \rightarrow -i \underline{\nabla} \quad E \rightarrow i \frac{\partial}{\partial t}$$

$$\Rightarrow i \frac{\partial \psi}{\partial t} = -i \alpha \cdot \nabla \psi + \beta m \psi$$

MULTIPLY BY β

$$i \beta \frac{\partial \psi}{\partial t} = -i \beta \alpha \cdot \nabla \psi + m \psi$$

DEFINE $\gamma^0 = \beta$ $\gamma^k = \beta \alpha^k$

$$i \gamma^0 \frac{\partial \psi}{\partial t} + i \gamma^k \nabla \psi - m \psi = 0$$

$$i \gamma^0 \frac{\partial \psi}{\partial x_0} + i \gamma^k \frac{\partial \psi}{\partial x_k} - m \psi = 0$$

$$\rightarrow (i\gamma^{\mu} \partial_{\mu} - m) \psi = 0$$

$$\gamma^{\mu} = (\gamma^0, \gamma^k)$$

$$\partial_{\mu} = \left(\frac{\partial}{\partial x^0}, \nabla \right)$$

β IS HERMITIAN $\rightarrow \gamma^0$ IS

BUT $(\gamma^k)^{\dagger} = -\gamma^k$

BECAUSE:

$$\gamma^k = \beta \alpha^k$$

$$\gamma^{\dagger} = (\alpha \alpha^k)^{\dagger} = \alpha^{k\dagger} \beta^{\dagger}$$

$$= \alpha^k \beta = -\beta \alpha^k = -\gamma^k$$

$$- (\gamma^0)^2 = \mathbb{1} \quad \text{SINCE } \beta^2 = \mathbb{1}$$

$$\begin{aligned} (\gamma^k)^2 &= \gamma^k \gamma^k = \beta \alpha^k \beta \alpha^k = -\beta \alpha^k \alpha^k \beta = \\ &= -\beta^2 = -\mathbb{1} \end{aligned}$$

TAKE NOW HERMITIAN
CONJUGATE OF DIRAC EQ.

$$-i \frac{\partial \psi^\dagger}{\partial t} \gamma^0 + i \frac{\partial}{\partial x^k} (-\gamma^k) \psi^\dagger - m \psi^\dagger = 0$$

MULTIPLY BY γ^0

$$-i \frac{\partial \psi^\dagger}{\partial t} - i \frac{\partial}{\partial x^k} \gamma^k \gamma^0 \psi^\dagger - m \psi^\dagger \gamma^0 = 0$$

$\gamma^0 \gamma^k = -\gamma^k \gamma^0$

$$+ i \frac{\partial}{\partial x^k} \gamma^0 \gamma^k \psi^\dagger$$

DEFINE ADJOINT $\bar{\psi} = \psi^\dagger \gamma^0$

$$-i \frac{\partial \bar{\psi}}{\partial t} \gamma^0 - i \frac{\partial \bar{\psi}}{\partial x^k} \gamma^k - m \bar{\psi} = 0$$

$$i \partial_\mu \bar{\psi} \gamma^\mu + m \bar{\psi} = 0$$

ADJOINT FORM OF DIRAC EQ

MULTIPLY "STANDARD"
EQUATION BY $\bar{\psi}$ AND ADJOINT
ONE BY ψ

$$i \bar{\psi} \gamma^{\mu} \partial_{\mu} \psi - \bar{\psi} m \psi = 0$$

$$i \partial_{\mu} \bar{\psi} \gamma^{\mu} \psi - \bar{\psi} m \psi = 0$$

TAKE DIFFERENCE

$$i \bar{\psi} \gamma^{\mu} \partial_{\mu} \psi + i \partial_{\mu} \bar{\psi} \gamma^{\mu} \psi = 0$$

$$= \partial_{\mu} (\bar{\psi} \gamma^{\mu} \psi) = 0$$

$$\text{SO } J^{\mu} = \bar{\psi} \gamma^{\mu} \psi$$

SATISFIES CONTINUITY EQN

$$\begin{aligned} P = J^0 &= \bar{\psi} \gamma^0 \psi = \psi^{\dagger} \gamma^0 \gamma^0 \psi = \psi^{\dagger} \psi \\ &= |\psi|^2 \end{aligned}$$

E. M. CURRENT DENSITY

$$J^{\mu} = -e \bar{\psi} \gamma^{\mu} \psi$$

FREE PARTICLE SOLUTION

$$\psi = U(p) e^{-i p x}$$

NEEDS TO
HAVE DIMENSION
4

$$(i \gamma^{\mu} \partial_{\mu} - m) U(p) e^{-i p x} = 0$$

$$(i \gamma^{\mu} (-i p_{\mu}) - m) U(p) e^{-i p x} = 0$$

$$(\gamma^{\mu} p_{\mu} - m) U(p) e^{-i p x} = 0$$

$$\Rightarrow (\gamma^{\mu} p_{\mu} - m) U(p) = 0$$

DEFINE $\not{p} = \gamma^{\mu} p_{\mu}$

$$(P - m)U(P) = 0$$

$$(\gamma^0 E - \gamma^k P_k - m)U(P) = 0$$

MULTIPLY BY γ^0

$$((\gamma^0)^2 E - \gamma^0 \gamma^k P_k - \gamma^0 m)U(P) = 0$$

$$(E - \alpha^k P_k - \beta m)U(P) = 0$$

FOR A PARTICLE AT
REST $\underline{P} = \underline{0}$

$$(E \underline{1} - \beta m)U(P) = 0$$

$$\begin{pmatrix} \underline{1} & 0 \\ 0 & \underline{1} \end{pmatrix} E U(P) = \begin{pmatrix} \underline{1} & 0 \\ 0 & -\underline{1} \end{pmatrix} m U(P)$$

SOLUTIONS

$$\begin{vmatrix} (m-E) & 0 \\ 0 & -(m+E) \end{vmatrix} = 0$$

OR IN 4-D

$$0 = \begin{vmatrix} m-E & 0 \\ 0 & m-E \\ 0 & -m-E \\ 0 & -m-E \end{vmatrix}$$

$$[(m-E)(m+E)]^2 = 0$$

SOLUTIONS : $E = \pm m$

IF PARTICLE NOT AT

REST $\underline{P} \neq \underline{0}$

$$\left[(\underline{\alpha} \cdot \underline{P}) + \beta M \right] U = 0$$

$$\left[\begin{pmatrix} 0 & \underline{\sigma} \\ \underline{\sigma} & 0 \end{pmatrix} \cdot \underline{P} + \begin{pmatrix} \mathbb{1} & 0 \\ 0 & -\mathbb{1} \end{pmatrix} M \right] \chi$$

$$\chi \begin{pmatrix} U_A \\ U_B \end{pmatrix} = E \begin{pmatrix} U_A \\ U_B \end{pmatrix}$$

$$\begin{cases} \underline{\sigma} \cdot \underline{P} U_A - M U_B = E U_B \\ \underline{\sigma} \cdot \underline{P} U_B + M U_A = E U_A \end{cases}$$

$$\left\{ \begin{array}{l} U_A = \frac{\sigma \cdot P}{E - M} U_B \\ U_B = \frac{\sigma \cdot P}{E + M} U_A \end{array} \right.$$

$$U_A = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \text{ SPIN UP}$$

$$U_B = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \text{ SPIN DOWN}$$

FOR $E > 0$

$$U_1 = \begin{pmatrix} 1 \\ 0 \\ \frac{\sigma \cdot P}{E + M} \\ 0 \end{pmatrix} N$$

$$U_2 = \begin{pmatrix} 0 \\ 0 \\ \frac{\sigma \cdot P}{E + M} \\ 1 \end{pmatrix} N$$

FOR $E < 0$

$$U_A = \frac{\sigma \cdot P}{E - M} \quad U_B =$$

$$= \frac{\sigma \cdot P}{-|E| - M} \quad U_B =$$

$$= - \frac{\sigma \cdot P}{|E| + M} \quad U_B$$

$$U_3 = N \begin{pmatrix} - \frac{\sigma \cdot P}{|E| + M} \\ 0 \\ 0 \\ 0 \end{pmatrix}$$

$$U_4 = N \begin{pmatrix} 0 \\ - \frac{\sigma \cdot P}{|E| + M} \\ 0 \\ 0 \end{pmatrix}$$

REDEFINE BASES

U_3, U_4 TO MAKE EXPLICIT

THAT THEY ARE NEGATIVE ENERGY SOLUTIONS

$$U^{(3,4)} = \begin{pmatrix} -i(-p)x & (2,1) \\ (p)0 & i(p)x \end{pmatrix} = N \begin{pmatrix} 1 \\ p \end{pmatrix} e$$

DEFINES $N^{2,1}$

$$N_2 = N \begin{pmatrix} \frac{0 \cdot p}{E+m} \\ 0 \\ | \\ 0 \end{pmatrix} \quad \begin{matrix} \text{SPIN} \\ \text{DOWN} \end{matrix}$$

$$N_1 = N \begin{pmatrix} 0 \\ \frac{0 \cdot p}{E+m} \\ 0 \\ | \end{pmatrix} \quad \begin{matrix} \text{SPIN} \\ \text{UP} \end{matrix}$$

CALCULATE N

$$\psi_1 = N \begin{pmatrix} 1 \\ 0 \\ \frac{\sigma \cdot p}{E+M} \\ 0 \end{pmatrix} e^{-i p x}$$

$$\psi_2 = N \begin{pmatrix} 0 \\ 1 \\ 0 \\ \frac{\sigma \cdot p}{E+M} \end{pmatrix} e^{-i p x}$$

ORTHOGONALITY

$$\int \psi_1^\dagger \psi_2 d^3x = 0$$

$$N^* N \begin{pmatrix} 1 & 0 & \frac{(\sigma \cdot p)^\dagger}{E+M} & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \\ 0 \\ \frac{\sigma \cdot p}{E+M} \end{pmatrix} = 0$$

IS SATISFIED

NORMALISATION

$$\int \psi_1^* \psi_1 d^3x = Z E$$

$$Z E = \int N^* N \left(1, 0, \left(\frac{\sigma \cdot p}{E + m} \right), 0 \right) \begin{pmatrix} 1 \\ 0 \\ \frac{\sigma \cdot p}{E + m} \\ 0 \end{pmatrix} d^3x$$

$$Z E = N^* N \left[1 + \left(\frac{\sigma \cdot p}{E + m} \right)^2 \right] d^3x$$

FOR A PARTICLE MOVING
ALONG Y

$$\sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\begin{aligned} \left(\frac{\sigma_y p_y}{E + m} \right)^2 &= \begin{pmatrix} 0 & -i p_y \\ i p_y & 0 \end{pmatrix}^2 = \begin{pmatrix} p_y^2 & 0 \\ 0 & p_y^2 \end{pmatrix} = \\ &= p_y^2 \mathbb{1} \end{aligned}$$

$$2E = N^* N \int \left(1 + \frac{p^2}{(E+m)^2} \right)^{1/3} d^3 p$$

$$= N^* N \int \frac{E^2 + 2Em + m + E - m}{(E+m)^2} d^3 p$$

$$= 2E N^* N \int \frac{E+m}{(E+m)^2} d^3 p$$

$$2E = 2E N^* N \frac{V}{E+m}$$

$$N = \sqrt{\frac{E+m}{V}}$$