Brunel University Queen Mary, University of London Royal Holloway, University of London University College London

Intercollegiate post-graduate course in High Energy Physics

Paper 1: The Standard Model

Friday, 29 January 2010

Time allowed for Examination: 3 hours

Answer 6 from 8 questions

Books and notes may be consulted

Question 1 (20 marks)

For a four-momentum, p_{μ} , show that $p_{\mu}p^{\mu}$ is a Lorentz invariant, by considering a Lorentz transformation along a spatial axis of your choice. [5]

At a collider, two high energy particles, A and B with energies E_A and E_B , which are much greater than their rest masses, collide head on. Derive the expression for the centre-of-mass energy. Using this expression, what would be the centre-of-mass energy of a proposed future facility ("LHeC") which is supposed to collide 7 TeV protons with 70 GeV electrons? Now consider particle B (the proton) to be at rest. Derive the formula for the centre-of-mass energy of such a fixed-target experiment. What electron beam energy would be required in the fixed-target experiment in order to achieve the same centre-of-mass energy as in the proposed LHeC facility?

[5]

Obtain the relation for the centre-of-mass energy in electron-neutrino scattering

$$s = m_e(2E_{\nu} + m_e).$$

[3]

A particle of mass M decays into two particles with masses m_1 and m_2 . Determine the energies of the decay products in the rest frame of the parent particle. [5]

Hence write down, in terms of masses and the centre-of-mass energy, the energy in the rest frame of particle A in a scattering, $AB \to CD$.

Question 2 (20 marks)

Given the cross section for the scattering of an electron from a fixed Coulomb potential of point charge Ze:

$$\frac{d\sigma}{d\Omega} = \frac{2(Z\alpha)^2 m^2}{|\mathbf{q}|^4} \operatorname{Tr} \left[\gamma_0 \frac{p_i + m}{2m} \gamma_0 \frac{p_f + m}{2m} \right] ,$$

where p_i and p_f are the initial and final momenta and $\mathbf{q} = \mathbf{p}_f - \mathbf{p}_i$, determine the Mott cross section :

$$\frac{d\sigma}{d\Omega} = \frac{(Z\alpha)^2}{4(\gamma\beta^2)^2(mc^2)^2\sin^4\theta/2} \left(1 - \beta^2\sin^2\frac{\theta}{2}\right).$$

Trace theorems used should be explicitly stated.

Show that in the non-relativistic limit

$$\frac{d\sigma}{d\Omega} = \frac{(Z\alpha)^2}{16E^2 \sin^4 \frac{\theta}{2}}$$

and in the extreme-relativistic limit

$$\frac{d\sigma}{d\Omega} = \frac{(Z\alpha)^2 \cos^2 \frac{\theta}{2}}{4E^2 \sin^4 \frac{\theta}{2}}.$$

[5]

[15]

Question 3 (20 marks)

Evaluate, in terms of the four-vectors (you do not have to convert to the Mandelstam variable),

Tr
$$\left[\gamma^{\mu} k \gamma^{\nu} (p + m) \gamma_{\nu} k' \gamma_{\mu} (p' + m) \right]$$

and

$$\operatorname{Tr}\left[\gamma^{\mu} k \gamma^{\nu} (p + m) \gamma_{\mu} k' \gamma_{\nu} (p' + m)\right].$$

that occur in the calculation of electron-photon scattering. Trace theorems and identities for γ matrices need not be derived, but should be quoted. [10]

In a massless limit, the terms in final squared transition amplitude for Compton scattering are

(a)
$$2e^4\left(-\frac{u}{s}\right)$$
 (b) $2e^4\left(-\frac{s}{u}\right)$ (c) $2e^4\frac{t}{us}(s+u+t)$.

Identify the Feynman diagram(s) which contribute to each term. Hence write down the final squared transition amplitudes when the incoming photon is real and when it is virtual. [5]

Given the Compton condition

$$\lambda' = \lambda + \frac{2\pi}{m}(1 - \cos\theta)$$

and the Klein-Nishina formula for Compton scattering

$$\frac{d\sigma}{d\Omega}(\lambda, \lambda') = \frac{\alpha^2}{4m^2} \left(\frac{\omega'}{\omega}\right)^2 \left[\frac{\omega'}{\omega} + \frac{\omega}{\omega'} + 4(\epsilon'^* \cdot \epsilon)^2 - 2\right]$$

determine the cross section in the low-energy limit, i.e. $\omega \to 0$, in terms of the fine structure constant, α , the mass of the electron, m, and the polarisation vectors of the photon, ϵ and ϵ'^* .

[5]

Question 4 (20 marks)

What property of the EM interaction means that photons do not self-couple? [2]

Explain the four terms in the Lagrangian of QED:

$$\mathcal{L} = \bar{\psi}(i\gamma_{\mu}\partial^{\mu} - m)\psi + e\bar{\psi}\gamma_{\mu}A^{\mu}\psi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$
[4]

Briefly explain the concept of a "running" coupling constant in QED where the variation is with the scale of the process, Q^2 . Draw two Feynman diagrams, one for QED and and for QCD, to illustrate the effect. And draw a diagram which leads to the QCD coupling having a different dependence.

State what are meant by global and local gauge transformations. [2]

Given the phase transformations of the wave function and the electromagnetic field:

$$\phi(x) \to \phi'(x) = \exp(iq\alpha)\phi(x)$$
 $A^{\mu}(x) \to A'^{\mu}(x) = A^{\mu}(x) - \partial^{\mu}\alpha(x)$

and the gauge-covariant derivative,

$$\partial^{\mu} \to D^{\mu} = \partial^{\mu} + iqA^{\mu}$$

show that the Klein-Gordon equation is invariant under these transformations. [7]

Question 5 (20 marks)

In the decay of a π^- at rest, $\pi^- \to e^- + \bar{\nu_e}$, show that

$$\frac{1}{2}\left(1 - \frac{v_e}{c}\right) = \frac{m_e^2}{m_\pi^2 + m_e^2} \,.$$

where v_e is the velocity of the electron.

[5]

To lowest order, the partial decays rate for pions are:

$$\frac{1}{\tau(\pi \to e\bar{\nu_e})} = \frac{\alpha_\pi^2}{4\pi} \left(1 - \frac{v_e}{c} \right) p_e^2 E_e , \qquad \frac{1}{\tau(\pi \to \mu\bar{\nu_\mu})} = \frac{\alpha_\pi^2}{4\pi} \left(1 - \frac{v_\mu}{c} \right) p_\mu^2 E_\mu .$$

where α_{π} is an effective coupling constant and E_e , E_{μ} and p_e , p_{μ} are the charged lepton's energy and momentum. Hence show:

$$\frac{\tau(\pi \to \mu \bar{\nu_{\mu}})}{\tau(\pi \to e \bar{\nu_{e}})} = \frac{m_e^2 (m_{\pi}^2 - m_e^2)^2}{m_{\mu}^2 (m_{\pi}^2 - m_{\mu}^2)^2}.$$

[5]

Use the analogue of the above equation for the decay of the K^- to estimate the ratio

$$\frac{\tau(K \to \mu \bar{\nu_{\mu}})}{\tau(K \to e \bar{\nu_{e}})}$$

and compare with the observed value $(2.4 \pm 0.1) \times 10^{-5}$.

Given the lifetimes $\tau(K \to \mu \bar{\nu_{\mu}}) = 1.948 \times 10^{-8} \text{s}$ and $\tau(\pi \to \mu \bar{\nu_{\mu}}) = 2.603 \times 10^{-8} \text{s}$, estimate α_K/α_{π} .

$$(m_K = 493.67 \,\text{MeV}, \, m_\pi = 139.57 \,\text{MeV}, \, m_\mu = 105.66 \,\text{MeV}, \, m_e = 0.511 \,\text{MeV}.)$$
 [5]

Draw quark model diagrams for the decays $\pi^- \to \mu^- + \bar{\nu_\mu}$ and $K^- \to \mu^- + \bar{\nu_\mu}$, stating which element of the CKM matrix is involved in each.

Neglecting masses, the ratio of the CKM elements is equal to α_K/α_{π} . Hence estimate $\sin \theta_{12}$.

[5]

Question 6 (20 marks)

In deep inelastic scattering at HERA, the four-momenta of the incoming and scattered electron are (E, \mathbf{p}) and (E', \mathbf{p}') , respectively. Show that the square of the four-momentum transfer is given by

$$Q^2 = 4EE'\sin^2\frac{\theta}{2}$$

where θ is the angle of the scattered electron. And show that the mass of the proton is related to the energies by

$$M(E - E') - 2EE'\sin^2\frac{\theta}{2} = 0.$$

[6]

Give a physical description of the kinematic variables, x, y and Q^2 , which describe deep inelastic scattering. [4]

The electron-quark cross section is:

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2 \hat{s}} 2e^4 q_i^2 \left[\frac{\hat{s}^2 + \hat{u}^2}{\hat{t}^2} \right] .$$

Show that this can be written in a more useful form as:

$$\frac{d\sigma}{dy} = \frac{2\pi\alpha^2}{Q^4}q_i^2s\left[1 + (1-y)^2\right].$$

[6]

The cross section for the QCD Compton and Boson-gluon fusion processes are :

$$\frac{d\sigma}{d\Omega} \sim -\frac{\hat{t}}{\hat{s}} - \frac{\hat{s}}{\hat{t}} + \frac{2\hat{u}Q^2}{\hat{s}\hat{t}}$$
$$\frac{d\sigma}{d\Omega} \sim \frac{\hat{u}}{\hat{t}} + \frac{\hat{t}}{\hat{u}} - \frac{2\hat{s}Q^2}{\hat{t}\hat{u}},$$

respectively. State where (e.g. with reference to Feynman diagrams or the Mandelstam channel) each term comes from. [4]

Question 7 (20 marks)

A possible decay of the W boson with associated four-momenta is

$$W(q) \rightarrow \mu(p) + \nu_{\mu}(k)$$
.

The transition amplitude squared is:

$$|T_{\rm fi}|^2 = \frac{g_W^2}{3} \left[p \cdot k + \frac{2}{M_W^2} (q \cdot k) (q \cdot p) \right].$$

In the W rest frame, show that :

$$q \cdot k = \frac{M_W^2}{2} \left(1 - \frac{m_\mu^2}{M_W^2} \right) \qquad q \cdot p = \frac{M_W^2}{2} \left(1 + \frac{m_\mu^2}{M_W^2} \right) \qquad p \cdot k = \frac{M_W^2}{2} \left(1 - \frac{m_\mu^2}{M_W^2} \right) . \tag{10}$$

Then using the relationship:

$$\Gamma = \frac{1}{16\pi M_W} |T_{\rm fi}|^2$$

derive the partial width in terms of the masses of the W and μ and show that this can be simplified to

$$\Gamma = \frac{G_F}{\sqrt{2}} \frac{M_W^3}{6\pi} \,.$$

$$(m_{\mu} = 105.66 \,\text{MeV}, \, M_W = 80.4 \,\text{GeV})$$
 [10]

Question 8 (20 marks)

In the Weinberg-Salam electroweak theory, the vacuum potential, boson masses, couplings, and mixing angle are related by :

$$M_Z = \frac{M_W}{\cos \theta_W}, \qquad \frac{g'}{g_W} = \tan \theta_W, \qquad M_W = \frac{vg_W}{2}$$

Hence show that

$$M_Z = \frac{1}{2} v \sqrt{g_W^2 + g'^2} \,.$$

[4]

From the Lagrangian

$$\frac{1}{8} \left[g_W^2 (v+h)^2 (W_\mu^1 - iW_\mu^2) (W_\mu^1 + iW_\mu^2) - (v+h)^2 (g'B_\mu - g_W W_\mu^3) (g'B^\mu - g_W W_3^\mu) \right]$$

derive the WWH and WWHH couplings and the ZZH and ZZHH couplings. (Simplify your answer to remove dependencies on both v and g'.)

Given Fermi's constant,
$$G_F = 1.16637 \times 10^{-5} \,\mathrm{GeV^{-2}}$$
, calculate the value of v . [4]

The Higgs Boson was searched for in e^+e^- collisions at LEP. Draw the Feynman diagram of the process for Higgs production. Masses of up to about 114 GeV were ruled out; hence give the approximate maximum centre-of-mass energy of LEP. [3]

[Total Marks = 120]

END OF PAPER