

Proposal for an improved “ M_X method”

Graeme Watt

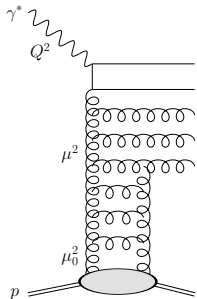
DESY Hamburg

ZEUS Diffractive Review Meeting
23rd November 2005

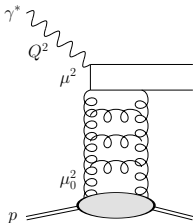
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Collinear factorisation in DDIS

Resolved \mathbb{P} :



Direct \mathbb{P} :



- **Pomeron** structure is analogous to **photon** structure: both **resolved** and **direct** contributions.
- $Q^2 \gg \dots \gg \mu^2 \gg \dots \gg \mu_0^2 \sim 1 \text{ GeV}^2$
- **Direct** Pomeron contribution responsible for **exclusive** diffractive processes, but also contributes to **inclusive** diffraction.

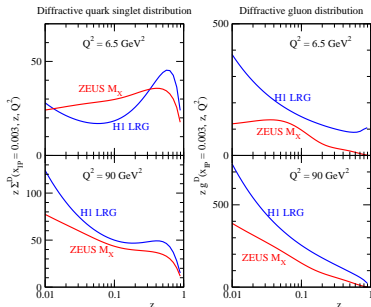
DDIS structure function:
$$F_2^{\text{D}(3)}(x_{\mathbb{P}}, \beta, Q^2) = \underbrace{\sum_{a=q,g} C_{2,a} \otimes a^{\text{D}}}_{\text{Resolved Pomeron}} + \underbrace{C_{2,\mathbb{P}}}_{\text{Direct Pomeron}}$$

DPDF evolution:
$$\frac{\partial a^{\text{D}}(x_{\mathbb{P}}, z, Q^2)}{\partial \ln Q^2} = \underbrace{\sum_{a'=q,g} P_{aa'} \otimes a'^{\text{D}}}_{\text{DGLAP term}} + \underbrace{P_{a\mathbb{P}}(z) f_{\mathbb{P}}(x_{\mathbb{P}}; Q^2)}_{\text{Inhomogeneous term}}$$

pQCD Pomeron flux factor:
$$f_{\mathbb{P}}(x_{\mathbb{P}}; Q^2) = \frac{1}{x_{\mathbb{P}} B_D} \left[R_g \frac{\alpha_S(Q^2)}{Q} x_{\mathbb{P}} g(x_{\mathbb{P}}, Q^2) \right]^2$$

Motivation for an improved “ M_X method”

- Diffractive gluon distribution sensitive to Q^2 dependence of DDIS data:



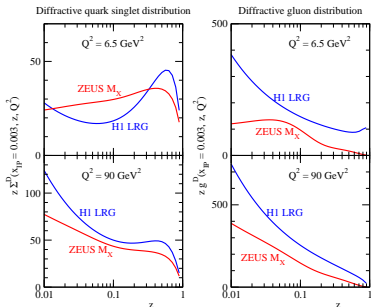
- See HERA-LHC proceedings.
- H1 LRG and ZEUS M_X data give very different DPDFs due to different Q^2 dependence of data sets.
- From results shown at last collaboration meeting:
 - New ZEUS LRG data also seem to have different Q^2 dependence than published ZEUS M_X data.
 - New ZEUS M_X (high Q^2) data seem to be compatible with published ZEUS M_X data.

Possible (tentative) explanations:

- 1 Amount of proton dissociation is Q^2 dependent?
But would be difficult to explain theoretically.
- 2 Significant non-diffractive contribution to LRG events (even at small $x_{\mathbb{P}}$)?
But should be exponentially suppressed compared to diffractive contribution.
- 3 Unjustified approximations made in “ M_X method” used to subtract non-diffractive events?

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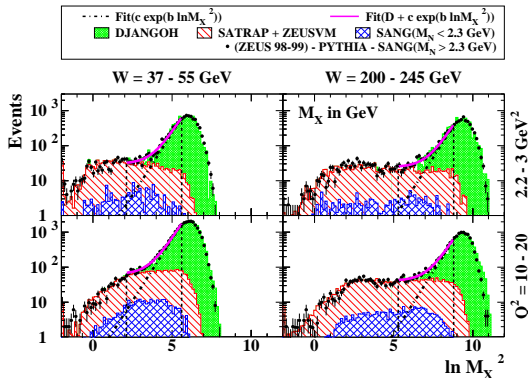
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Reminder of the “ M_X method”

- Used in three ZEUS papers [Z. Phys. C **70** (1996) 391; Eur. Phys. J. C **6** (1999) 43; Nucl. Phys. B **713** (2005) 3] + **analysis in progress**.
- Subtract non-diffractive events in each (W, Q^2) bin by fitting (in a limited range of $\ln M_X^2$):

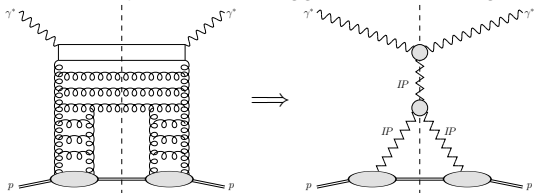
$$\frac{dN}{d \ln M_X^2} = D + \underbrace{c \exp(b \ln M_X^2)}_{\text{non-diffractive}}$$

ZEUS

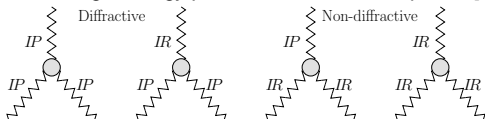


Regge theory of DDIS

- Replace pQCD ladders by “effective” Regge trajectories, e.g.



- But don't expect the “effective” trajectories to be universal, e.g. “effective” $\alpha_{\mathbb{P}}(0)$ greater than value for “soft” Pomeron and depends on Q^2 [NPB 713 (2005) 3].
- For $W^2 \gg M_X^2$ and $Q^2 \gg t$, consider triple Regge diagrams [see Barone & Predazzi, *High-energy particle diffraction*, Chap. 10.5]:



$$\frac{d\sigma_{\gamma^* p}}{d \ln x_{\mathbb{P}}} = \frac{|g_{\mathbb{P}}(\bar{t})|^2}{16\pi^2} x_{\mathbb{P}}^{2-2\alpha_{\mathbb{P}}(\bar{t})} \left[\mathcal{A}_{\mathbb{P}\mathbb{P}\mathbb{P}}(Q^2) \beta^{1-\alpha_{\mathbb{P}}(0)} + \mathcal{A}_{\mathbb{P}\mathbb{P}\mathbb{R}}(Q^2) \beta^{1-\alpha_{\mathbb{R}}(0)} \right] \\ + \frac{|g_{\mathbb{R}}(\bar{t})|^2}{16\pi^2} x_{\mathbb{P}}^{2-2\alpha_{\mathbb{R}}(\bar{t})} \left[\mathcal{A}_{\mathbb{R}\mathbb{R}\mathbb{P}}(Q^2) \beta^{1-\alpha_{\mathbb{P}}(0)} + \mathcal{A}_{\mathbb{R}\mathbb{R}\mathbb{R}}(Q^2) \beta^{1-\alpha_{\mathbb{R}}(0)} \right],$$

where $\alpha_{\mathbb{P}}(0) \approx 1.1-1.2$, $\alpha_{\mathbb{R}}(0) \lesssim 0.5$, and \bar{t} is some average value of t .

The “ $(M_X^2 + Q^2)$ method”

- Since $x_P = (M_X^2 + Q^2)/(W^2 + Q^2)$ and $\beta = Q^2/(M_X^2 + Q^2)$, rewrite as

$$\begin{aligned} \frac{d\sigma_{\gamma^*p}}{d\ln(M_X^2 + Q^2)} &= A_{\text{PPP}}(\bar{t})(W^2)^{2\alpha_P(\bar{t})-2}(M_X^2 + Q^2)^{1+\alpha_P(0)-2\alpha_P(\bar{t})} \\ &\quad + A_{\text{PPR}}(\bar{t})(W^2)^{2\alpha_P(\bar{t})-2}(M_X^2 + Q^2)^{1+\alpha_R(0)-2\alpha_P(\bar{t})} \\ &\quad + A_{\text{RRP}}(\bar{t})(W^2)^{2\alpha_R(\bar{t})-2}(M_X^2 + Q^2)^{1+\alpha_P(0)-2\alpha_R(\bar{t})} \\ &\quad + A_{\text{RRR}}(\bar{t})(W^2)^{2\alpha_R(\bar{t})-2}(M_X^2 + Q^2)^{1+\alpha_R(0)-2\alpha_R(\bar{t})}. \end{aligned}$$

- The “ M_X method” assumes that $M_X^2 \gg Q^2$ ($\Rightarrow \beta \ll 1$, so PPR and RRR contributions are negligible), and $\alpha_P(0) \approx \alpha_P(\bar{t}) \approx 1$. Then

$$\frac{d\sigma_{\gamma^*p}}{d\ln M_X^2} = D + c(M_X^2)^b = D + c \exp(b \ln M_X^2),$$

where $b = 1 + \alpha_P(0) - 2\alpha_R(\bar{t})$.

- More generally, look at $(M_X^2 + Q^2)$ distribution:

$$\frac{d\sigma_{\gamma^*p}}{d\ln(M_X^2 + Q^2)} = \underbrace{c_P(M_X^2 + Q^2)^{-b_P}}_{\text{diffractive}} + \underbrace{c_R(M_X^2 + Q^2)^{b_R}}_{\text{non-diffractive}},$$

where $-b_P \leq 1 + \alpha_P(0) - 2\alpha_P(\bar{t}) \approx -(0.1-0.2)$ and $b_R \geq 1 + \alpha_R(0) - 2\alpha_R(\bar{t}) \gtrsim 0.5$.

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where $-b_P \leq 1 + \alpha_P(0) - 2\alpha_P(\bar{t}) \approx -(0.1-0.2)$ and $b_R \geq 1 + \alpha_R(0) - 2\alpha_R(\bar{t}) \gtrsim 0.5$.

The “ $x_{\mathbb{P}}$ method”

- The “($M_X^2 + Q^2$) method” is better than the “ M_X method”, but it uses a combination of two powers of ($M_X^2 + Q^2$) to approximate a combination of four powers.
- Much more direct way of subtracting the non-diffractive contribution is to look at the $x_{\mathbb{P}}$ distribution:

$$\frac{d\sigma_{\gamma^*p}}{d \ln x_{\mathbb{P}}} = \frac{|g_{\mathbb{P}}(\bar{t})|^2}{16\pi^2} x_{\mathbb{P}}^{2-2\alpha_{\mathbb{P}}(\bar{t})} \left[\mathcal{A}_{\text{PPP}}(Q^2)\beta^{1-\alpha_{\mathbb{P}}(0)} + \mathcal{A}_{\text{PPR}}(Q^2)\beta^{1-\alpha_{\mathbb{R}}(0)} \right] \\ + \frac{|g_{\mathbb{R}}(\bar{t})|^2}{16\pi^2} x_{\mathbb{P}}^{2-2\alpha_{\mathbb{R}}(\bar{t})} \left[\mathcal{A}_{\text{RRP}}(Q^2)\beta^{1-\alpha_{\mathbb{P}}(0)} + \mathcal{A}_{\text{RRR}}(Q^2)\beta^{1-\alpha_{\mathbb{R}}(0)} \right]$$

$$\Rightarrow \boxed{\frac{d\sigma_{\gamma^*p}}{d \ln x_{\mathbb{P}}} = \underbrace{c_{\mathbb{P}}(x_{\mathbb{P}})^{-b_{\mathbb{P}}}}_{\text{diffractive}} + \underbrace{c_{\mathbb{R}}(x_{\mathbb{P}})^{b_{\mathbb{R}}}}_{\text{non-diffractive}},}$$

where $-b_{\mathbb{P}} = 2 - 2\alpha_{\mathbb{P}}(\bar{t}) \approx -(0.2-0.4)$ and $b_{\mathbb{R}} = 2 - 2\alpha_{\mathbb{R}}(\bar{t}) \gtrsim 1$.

- In most general case, fit four parameters $c_{\{\mathbb{P},\mathbb{R}\}} \geq 0$ and $b_{\{\mathbb{P},\mathbb{R}\}} \geq 0$ separately in each (β, Q^2) bin, in some limited range of $x_{\mathbb{P}}$.

Summary

- “ M_X method” justified if $M_X^2 \gg Q^2$ and $\alpha_P(0) = 1$:

$$\frac{dN}{d \ln M_X^2} = \underbrace{D}_{\text{diffractive}} + \underbrace{c(M_X^2)^b}_{\text{non-diffractive}}$$

- “ $(M_X^2 + Q^2)$ method” more general:

$$\frac{dN}{d \ln(M_X^2 + Q^2)} = \underbrace{c_P(M_X^2 + Q^2)^{-b_P}}_{\text{diffractive}} + \underbrace{c_R(M_X^2 + Q^2)^{b_R}}_{\text{non-diffractive}}$$

- “ x_P method” even better:

$$\frac{dN}{d \ln x_P} = \underbrace{c_P(x_P)^{-b_P}}_{\text{diffractive}} + \underbrace{c_R(x_P)^{b_R}}_{\text{non-diffractive}}$$

- More details given in write-up uploaded to ZEMS.