



FUTURE CIRCULAR COLLIDERS



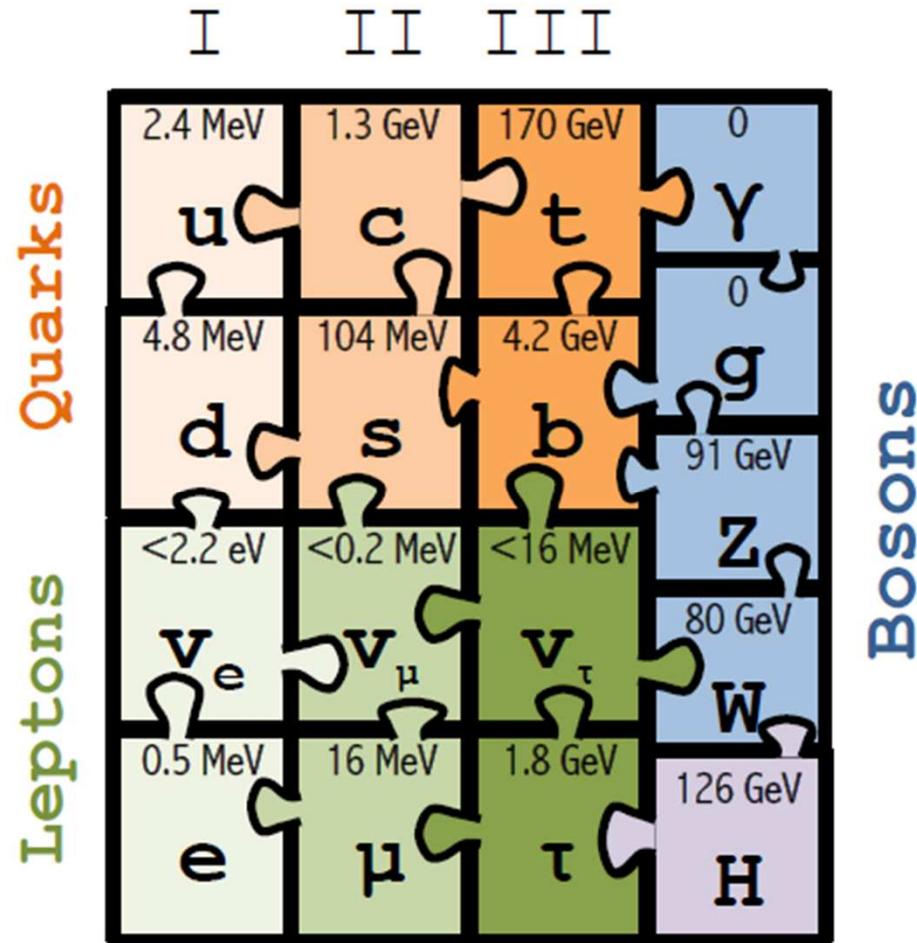
join us! <http://cern.ch/fcc-ee>
<http://espace2013.cern.ch/fcc/Pages/Science.aspx>

4/30/2015

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1997-2013 Higgs boson mass cornered (LEP H, M_Z etc +Tevatron m_t , M_W)
Higgs Boson discovered (LHC)
Englert and Higgs get Nobel Prize



(c) Sfyrla

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Asymptotic safety of gravity and the Higgs boson mass

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12 January 2010

Abstract

There are indications that gravity is asymptotically safe. The Standard Model (SM) plus gravity could be valid up to arbitrarily high energies. Supposing that this is indeed the case and assuming that there are no intermediate energy scales between the Fermi and Planck scales we address the question of whether the mass of the Higgs boson m_H can be predicted. For a positive gravity induced anomalous dimension $A_\lambda > 0$ the running of the quartic scalar self interaction λ at scales beyond the Planck mass is determined by a fixed point at zero. This results in $m_H = m_{\min} = 126$ GeV, with only a few GeV uncertainty. This prediction is independent of the details of the short distance running and holds for a wide class of extensions of the SM as well. For $A_\lambda < 0$ one finds m_H in the interval $m_{\min} < m_H < m_{\max} \simeq 174$ GeV, now sensitive to A_λ and other properties of the short distance running. The case $A_\lambda > 0$ is favored by explicit computations existing in the literature.

Key words:

Asymptotic safety, gravity, Higgs field, Standard Model

PACS: 04.60.-m 11.10.Hi 14.80.Bn

Detecting the Higgs scalar
with mass around 126 GeV at the LHC could give a
strong hint for the absence of new physics influencing
the running of the SM couplings between the Fermi and
Planck/unification scales.





Is it the end?

- Dark matter
- Baryon Asymmetry in Universe
- Neutrino masses
- and... why are the charges of e and p identical to 21 significant digits?

are *experimental* proofs that there is more to understand.

We must continue our quest

«We can extrapolate to the Planck scale»

and

«There MUST be new physics at TeV scale»

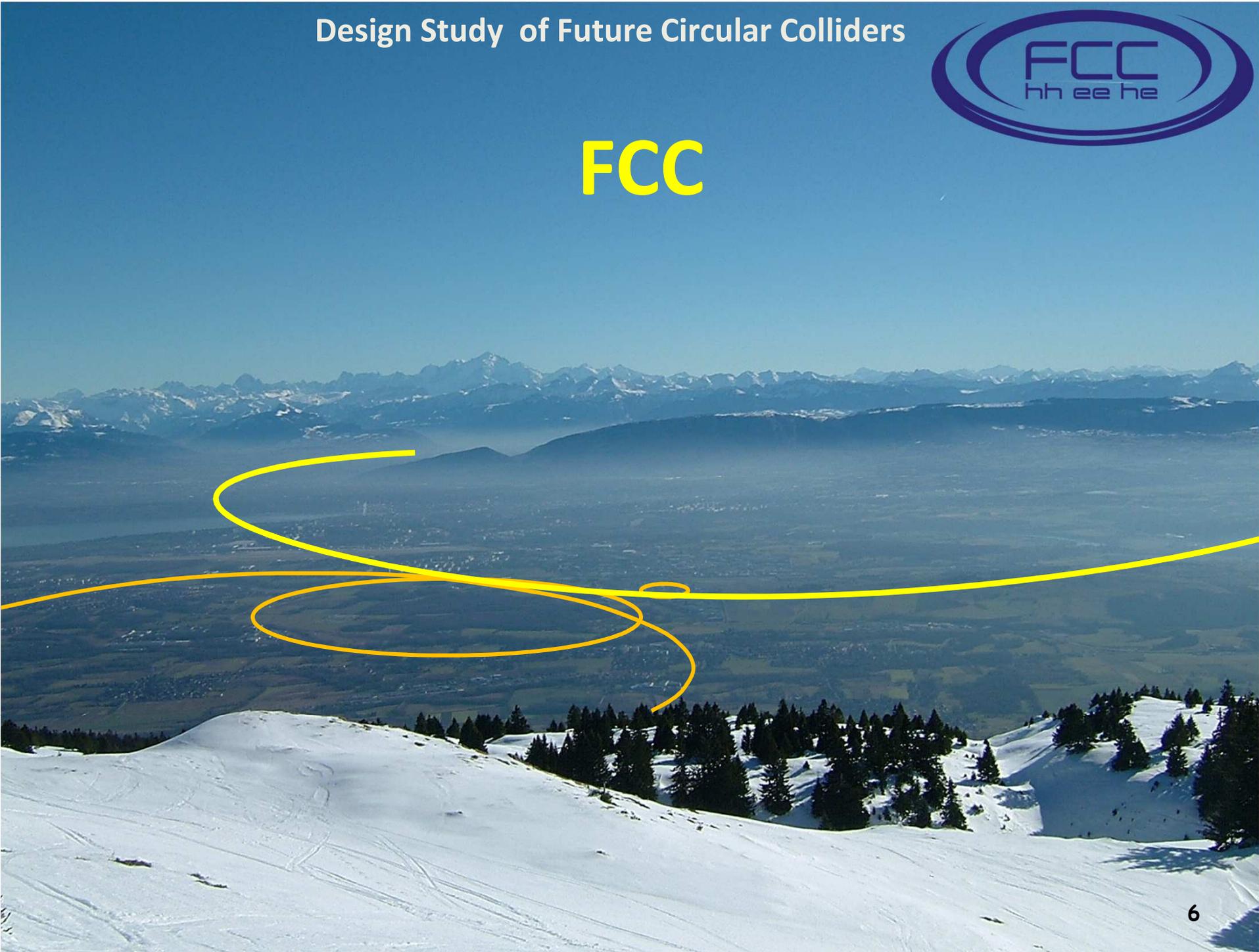
mutually exclusive?

There is one way to find out: go look!





FCC



Future Circular Collider Study - SCOPE

CDR and cost review for the next ESU (2018)

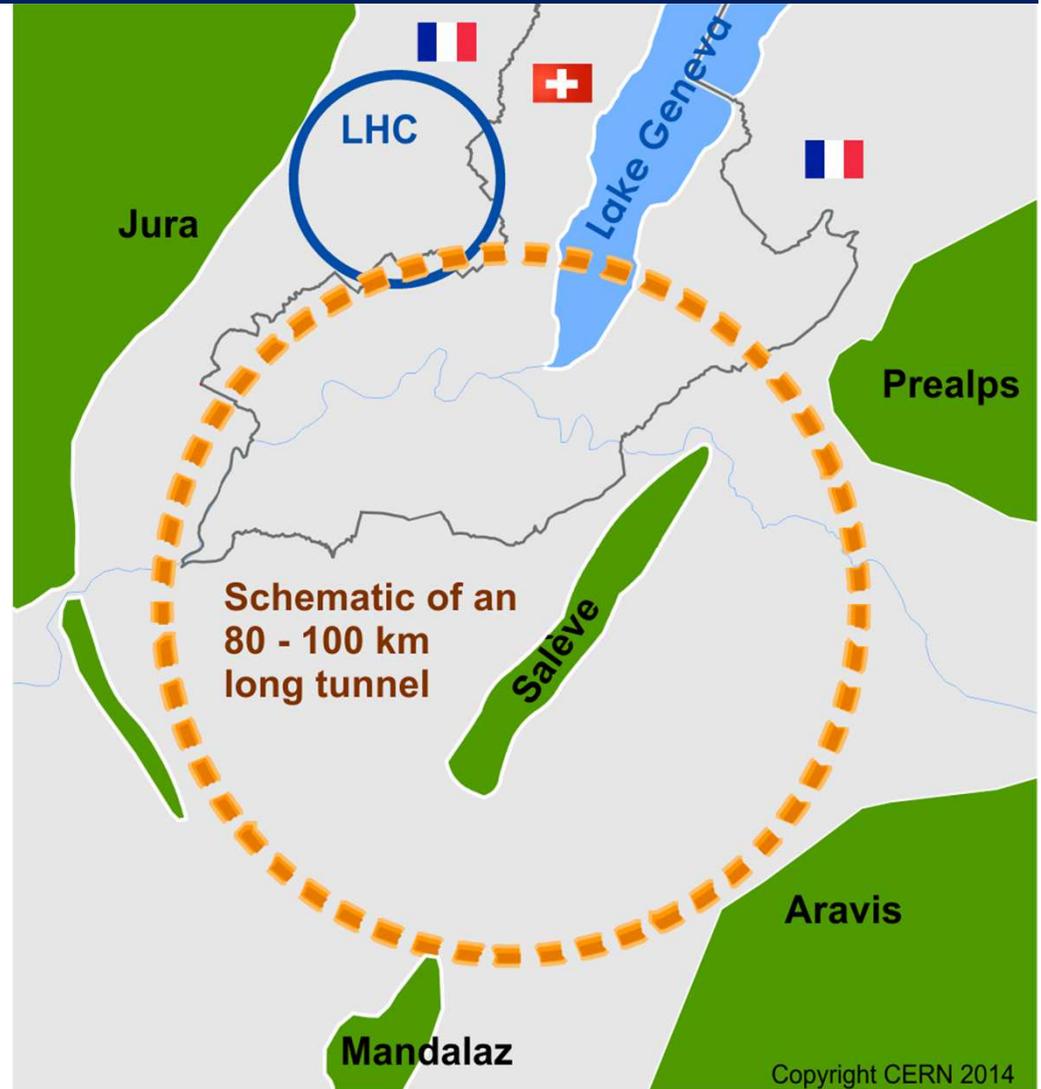
Forming an international collaboration to study:

- ***pp*-collider (*FCC-hh*)**
→ defining infrastructure

~16 T ⇒ 100 TeV *pp* in 100 km

~20 T ⇒ 100 TeV *pp* in 80 km

- ***e⁺e⁻* collider (*FCC-ee*) as potential intermediate step ECM=90-400 GeV**
- ***p-e* (*FCC-he*) option**
- **80-100 km infrastructure in Geneva area**



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Future Circular Collider Study - FCC

Mandate

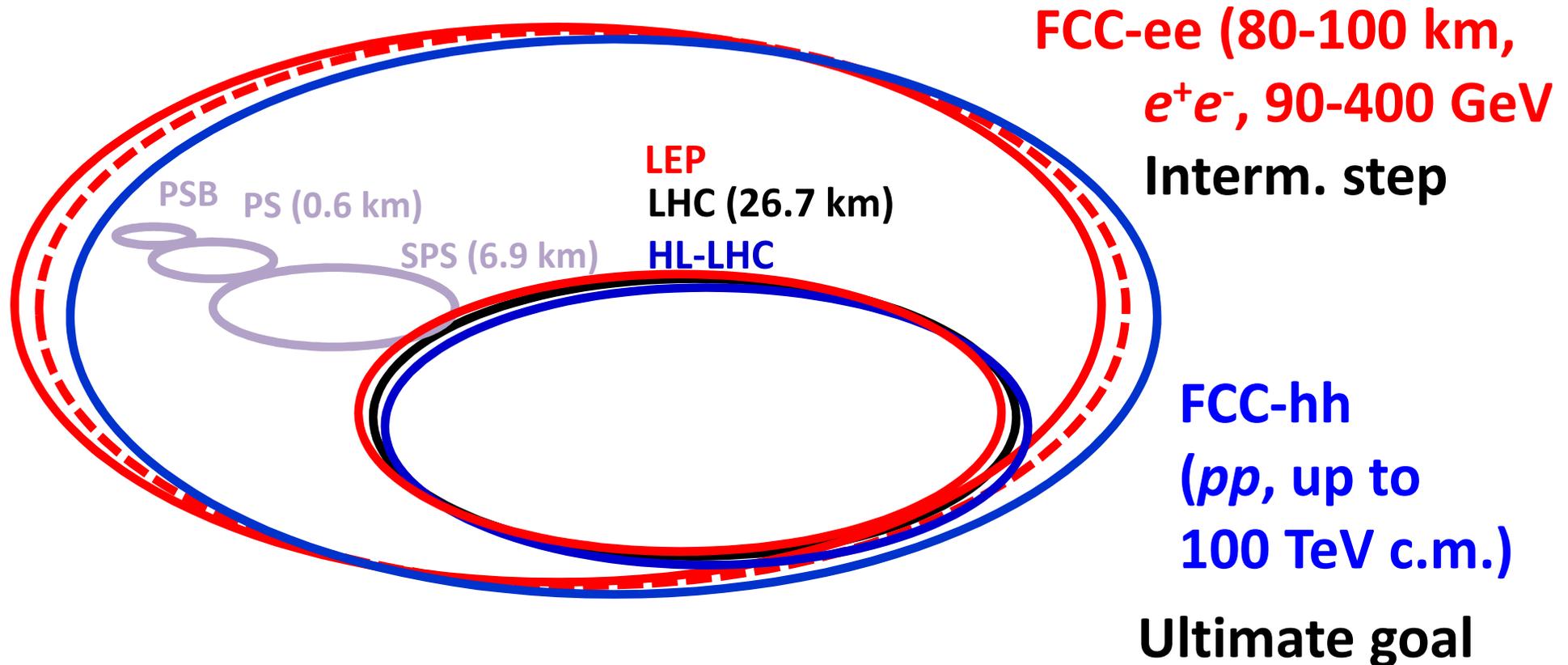
Scope

The main emphasis of the conceptual design study shall be the long-term goal of a hadron collider with a centre-of-mass energy of the order of 100 TeV (currently referred to as VHE-LHC) in a new tunnel of 80-100 km circumference for the purposes of studying physics at the highest energies. The hadron collider and its detectors shall determine the basic requirements for the tunnel, surface and technical infrastructures. The corresponding hadron injector chain shall be included in the study, taking into account the existing CERN accelerator infrastructure and long-term accelerator operation plans. The performance and cost of the hadron collider shall be compared to a high-energy LHC based on the same high-field magnet technology and housed in the LHC tunnel.

The conceptual design study shall also include a lepton collider and its detectors (currently referred to as TLEP), as a potential intermediate step towards realization of the hadron facility. The design of the lepton collider complex shall be based on the hadron collider infrastructure and any substantial incompatibilities with respect to the hadron collider infrastructure requirements shall be analysed and quantified. Potential synergies with linear collider detector designs should be considered.



possible long-term strategy



& e^\pm (120 GeV)– p (7, 16 & 50 TeV) collisions FCC-eh)

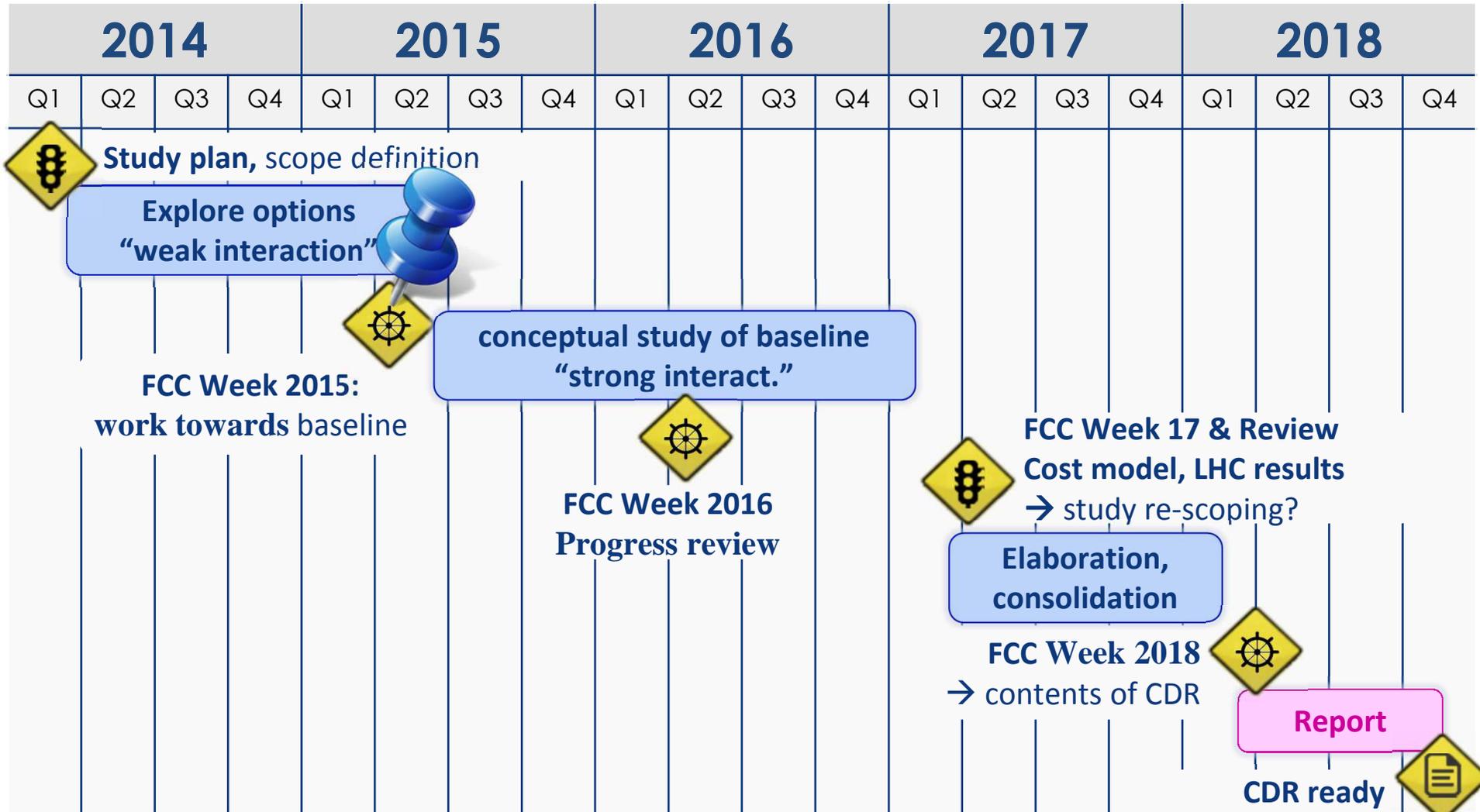
≥60 years of e^+e^- , pp , ep/A physics at highest energies

4/30/2015





Study time line towards CDR





93km "optimised" racetrack



PRELIMINARY

Alignment Shaft Tools

Choose alignment option
90km quasi-circular

Tunnel depth at centre: 236mASL

Gradient Parameters

Azimuth (°): -15
Slope Angle x-x(%): .3
Slope Angle y-y(%): 0

CAI CII ATF

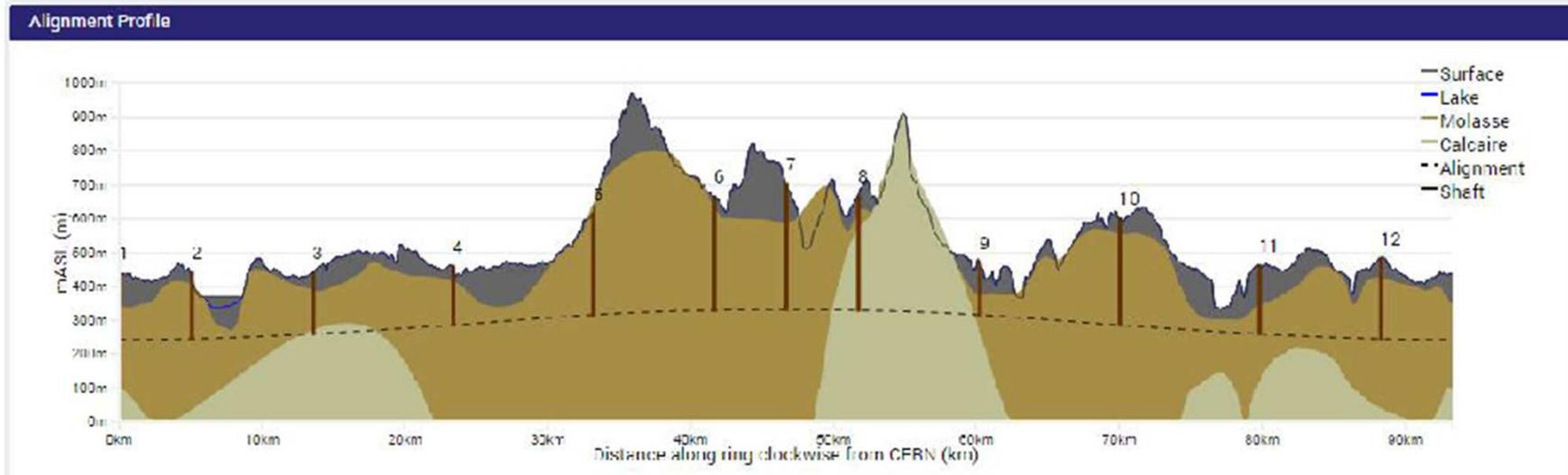
Alignment centre
X: 2493923 Y: 1105695

LHC Intersection	IP 1	IP 2
Angle	1°	-1°
Depth	542m	542m

Alignment Location

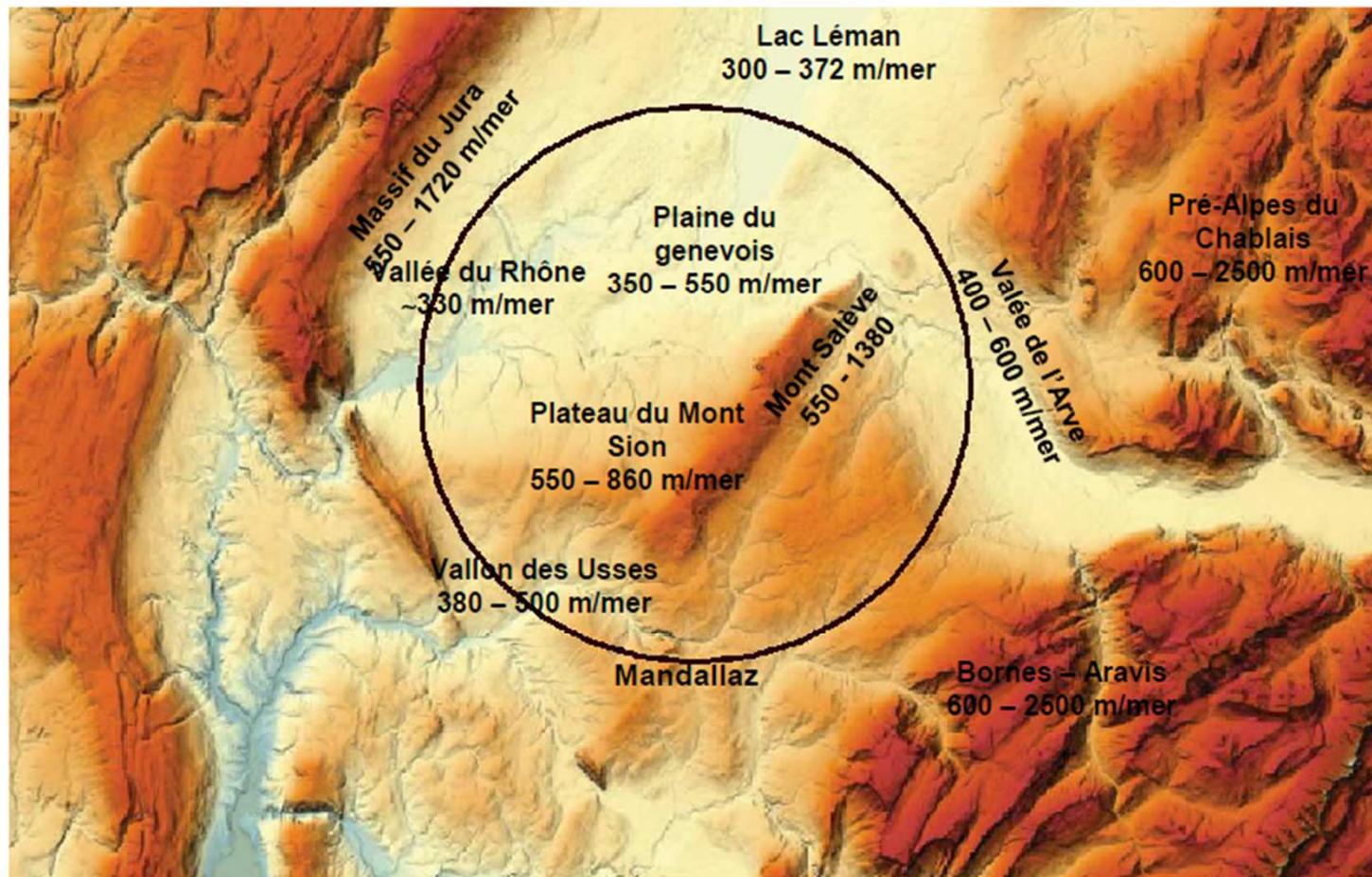
Geology Intersected by Shafts Shaft Depths

Shaft	Shaft Depth (m)				Geology (m)		
	Actual	Min	Mean	Max	Moreaine	Molasse	Calcaire
1	230	195	197	230	92	108	0
2	196	143	181	211	54	167	0
3	183	175	184	194	63	121	9
4	174	145	166	178	44	130	0
5	299	285	311	350	0	325	0
6	336	325	339	350	55	307	0
7	374	340	377	412	119	256	0
8	397	378	341	366	44	66	257
9	155	131	145	157	94	61	0
10	315	305	320	336	46	269	0
11	233	199	202	234	122	81	0
12	239	229	238	243	58	181	0
Total	3014	2801	3001	3211	711	2062	247



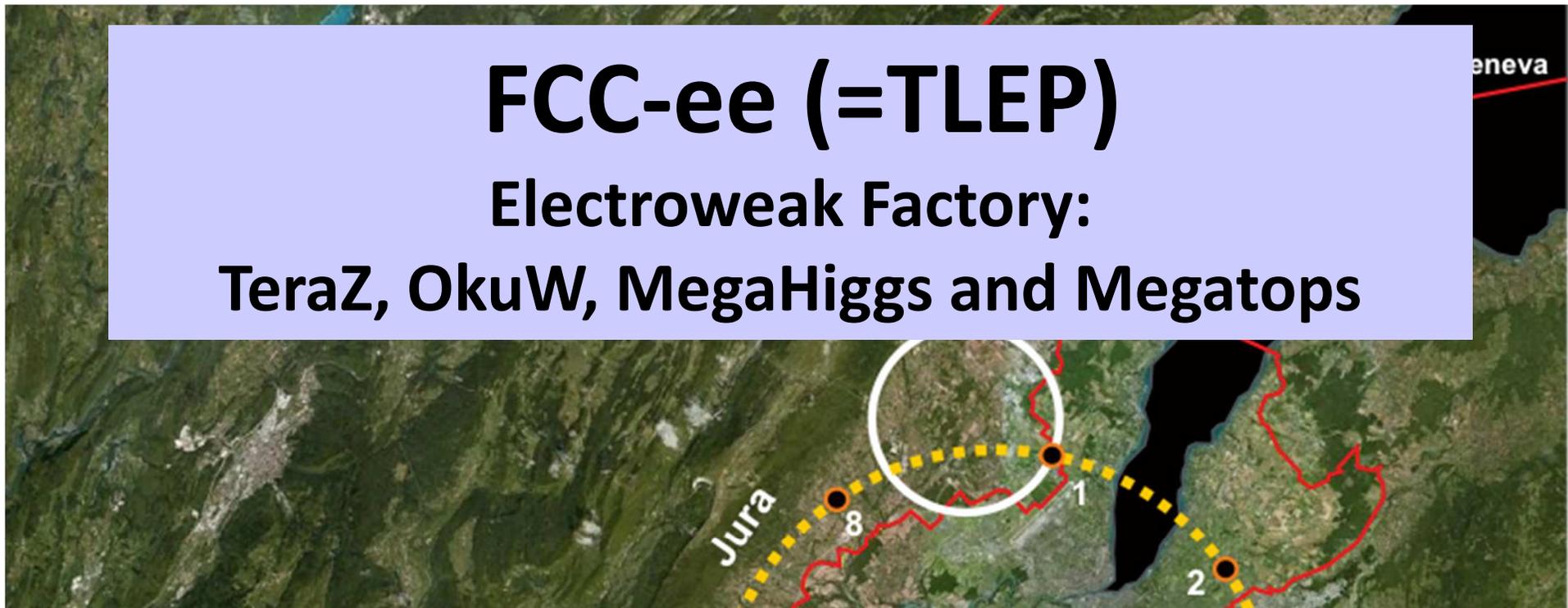
J. Osborne & C. Cook

- Minimize ground coverage
 - Hydrostatic pressure for TBM tunnelling
 - Shaft depth/cost



FCC-ee (=TLEP)

Electroweak Factory: TeraZ, OkuW, MegaHiggs and Megatops



Acknowledgments to all my FCC-ee colleagues for material and ideas (and hard work) in particular: J.Wenninger, F. Zimmermann, P. Lebrun, E. Jensen, R. Thomas, B. Harer, R. Martin, N. Bacchetta, P. Janot, B. Holzer, H. Burkhardt (CERN) M. Koratzinos (UNIGE), U. Wienands (SLAC) E. Gianfelice (FNAL), M. Boscolo (LNF) A.Bogomyagkov, I. Koop, E. Levichev, D. Shatilov, I. Telnov (BINP Novosibirsk) K. Ohmi, K. Oide (KEK)



Original motivation (end 2011): now that m_H and m_{top} are known, explore EW region with a high precision, affordable, high luminosity machine

→ Discovery of New Physics in rare phenomena or precision measurements

ILC studies → need increase over LEP 2 (average) luminosity by a factor 1000
How can one do that without exploding the power bill?

Answer is in the B-factory design: a low vertical emittance ring with higher intrinsic luminosity, and small β_y^* (1mm vs 5cm at LEP)

Electrons and positrons have a much higher chance of interacting

→ much shorter lifetime (few minutes)

→ top up continuously with booster ==> increase operation efficiency

Increase SR beam power to 50MW/beam

50

5

4

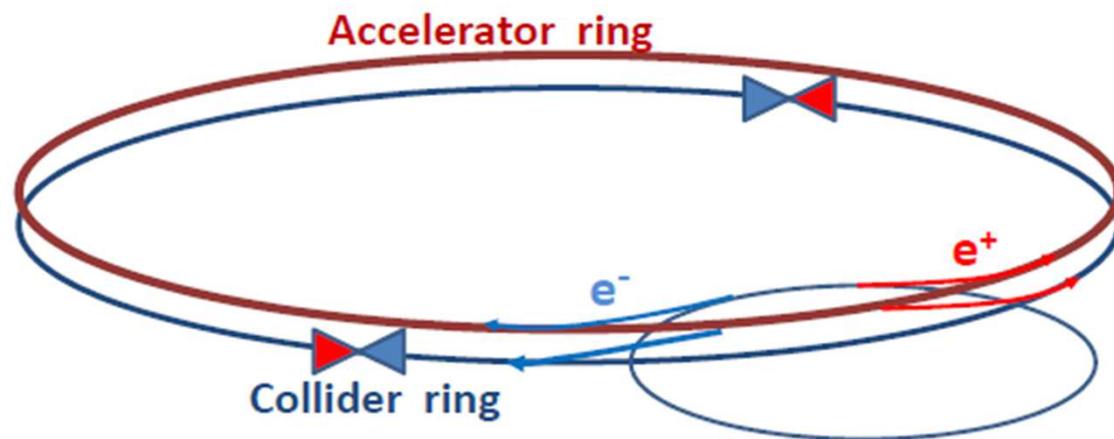
1000

at ZH threshold
in LEP/LHC tunnel

X 4 in FCC tunnel

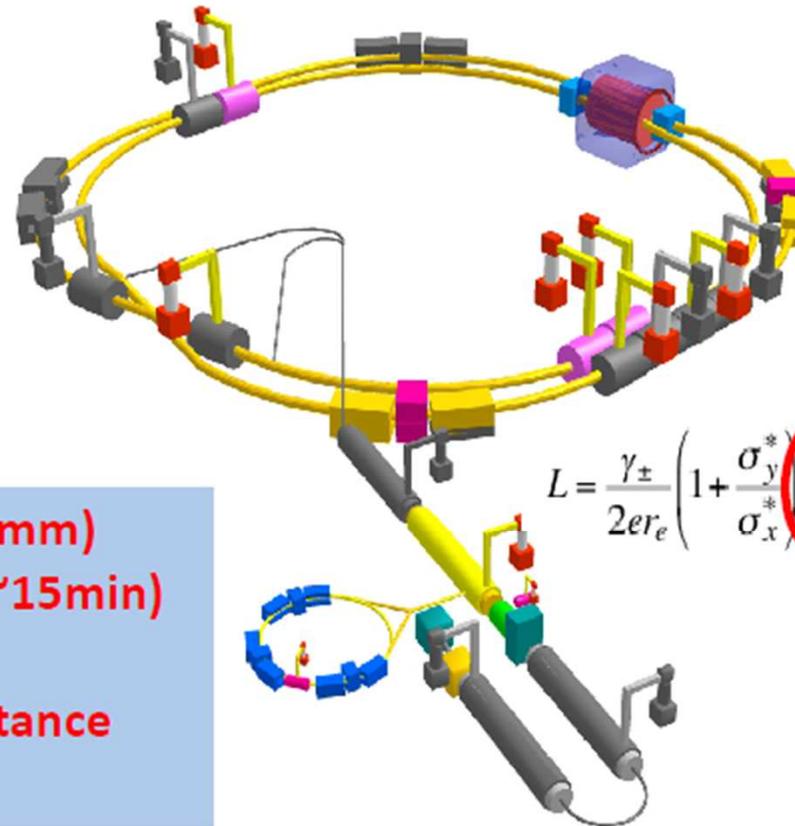
X 4 interaction points

EXCITING!



SuperKEKB – TLEP demonstrator!

beam
commissioning will
start in early 2015

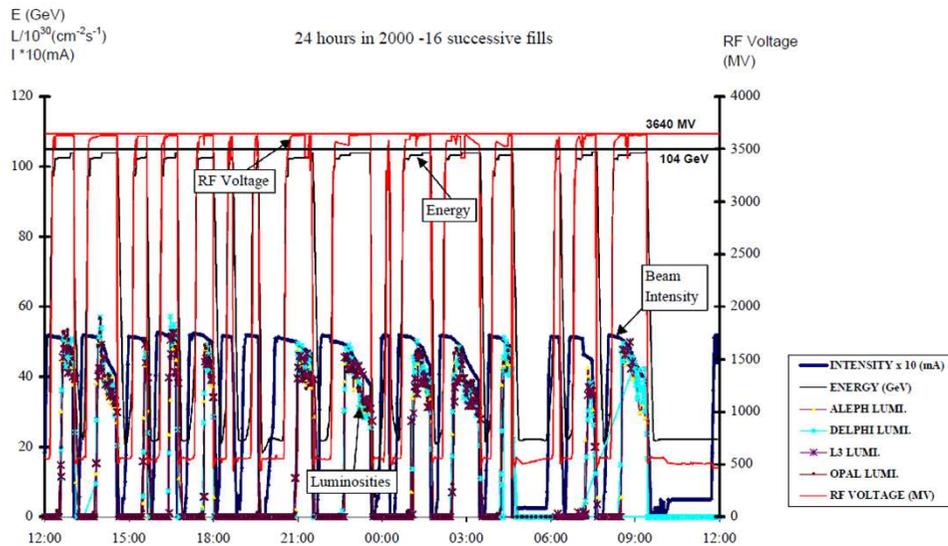


- $\beta_y^* = 300 \mu\text{m}$ (TLEP: 1 mm)
- lifetime 5 min (TLEP: ~15min)
- $\varepsilon_y/\varepsilon_x = 0.25\%$ (~TLEP)
- off momentum acceptance
- e^+ production rate

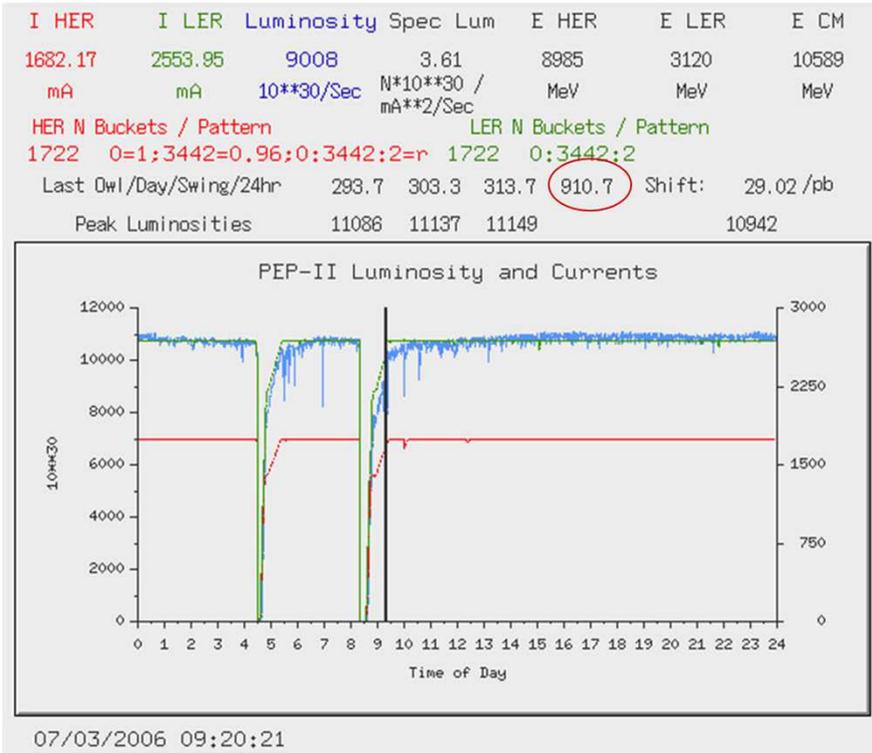
$$L = \frac{\gamma_{\pm}}{2e r_e} \left(1 + \frac{\sigma_y^*}{\sigma_x^*} \left(\frac{I_{\pm} \xi_{\pm y}}{\beta_y^*} \right) \left(\frac{R_L}{R_y} \right) \right)$$



Topping up ensures constant current, settings, etc... and greater reproducibility of system



LEP2 in 2000 (12th year!):
fastest possible turnaround but
average luminosity ~ 0.2 peak luminosity



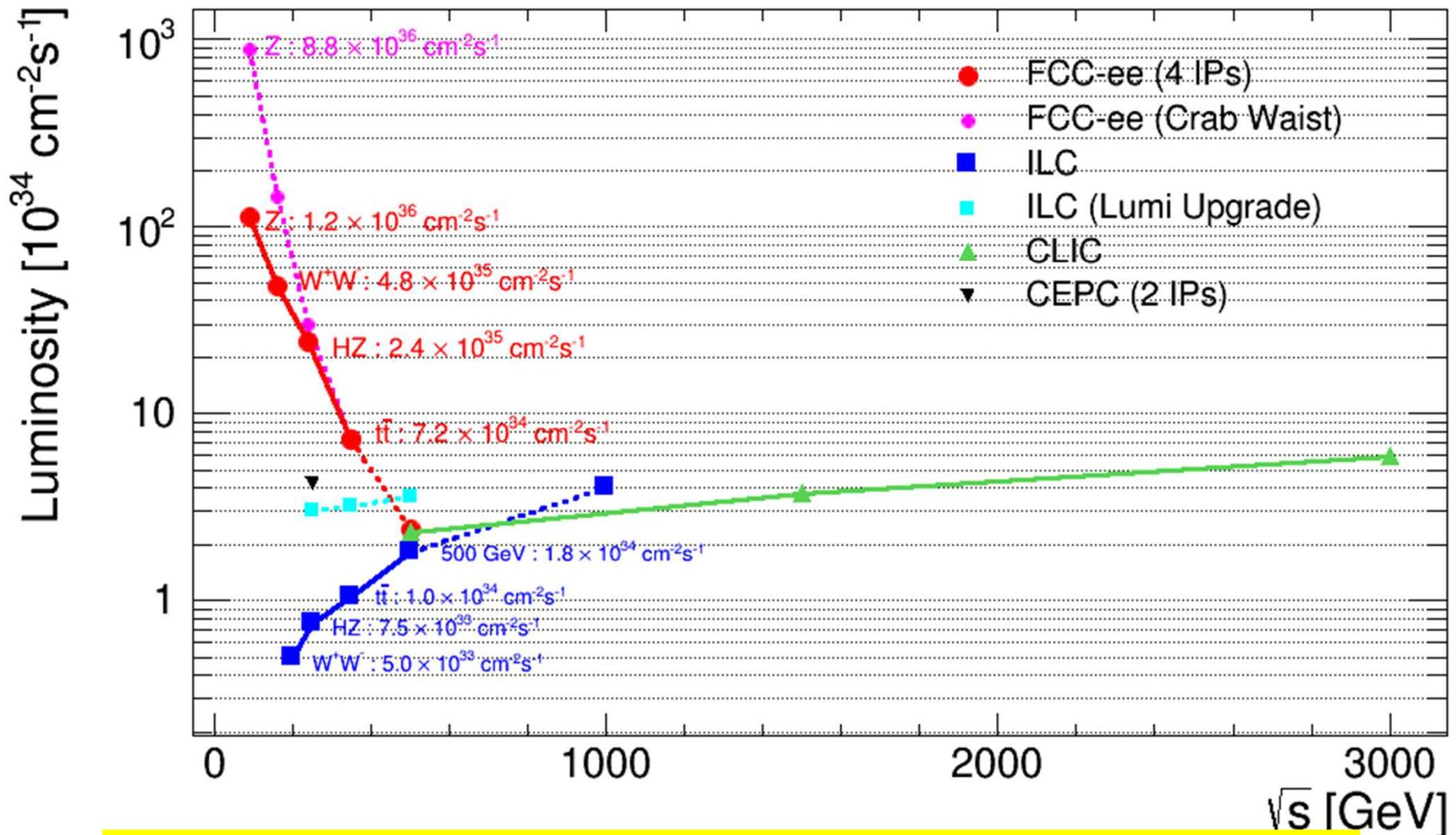
B factory in 2006 with topping up
average luminosity \approx peak luminosity



parameter	LEP2	FCC-ee				
		Z	Z (c.w.)	W	H	t
E_{beam} [GeV]	104	45	45	80	120	175
beam-beam par. ξ_y/IP	0.06	0.03	0.175	0.06	0.093	0.092
current [mA]	3.0	1450	1431	152	30	6.6
$P_{\text{SR,tot}}$ [MW]	22	100	100	100	100	100
no. bunches	4	16700	29791	4490	1360	98
N_b [10^{11}]	4.2	1.8	1.0	0.7	0.46	1.4
ϵ_x [nm]	22	29	0.14	3.3	0.94	2
ϵ_y [pm]	250	60	1	1	2	2
β_x^* [m]	1.2	0.5	0.5	0.5	0.5	1.0
β_y^* [mm]	50	1	1	1	1	1
σ_y^* [nm]	3500	250	32	84	44	45
$\sigma_{z,\text{SR}}$ [mm]	11.5	1.64	2.7	1.01	0.81	1.16
$\sigma_{z,\text{tot}}$ [mm] (w beamstr.)	11.5	2.56	5.9	1.49	1.17	1.49
hourglass factor F_{hg}	0.99	0.64	0.94	0.79	0.80	0.73
L/IP [$10^{34} \text{ cm}^{-2}\text{s}^{-1}$]	0.01	28	212	12	6	1.7
τ_{beam} [min]	434	298	39	73	29	21

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Colliders



Overlapp in Higgs/top region, but differences and complementarities between linear and circular machines



TLEP: PARAMETERS & STATISTICS

($e^+e^- \rightarrow ZH$, $e^+e^- \rightarrow W^+W^-$, $e^+e^- \rightarrow Z$, [$e^+e^- \rightarrow t\bar{t}$])

	TLEP-4 IP, per IP	statistics
circumference	80 km	
max beam energy	175 GeV	
no. of IPs	4	
Luminosity/IP at 350 GeV c.m.	$1.3 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$	10^6 $t\bar{t}$ pairs
Luminosity/IP at 240 GeV c.m.	$6.0 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$	2×10^6 ZH evts
Luminosity/IP at 160 GeV c.m.	$1.6 \times 10^{35} \text{ cm}^{-2}\text{s}^{-1}$	10^8 WW pairs
Luminosity/IP at 90 GeV c.m.	$2 \cdot 10^{35/36} \text{ cm}^{-2}\text{s}^{-1}$	$10^{12/13}$ Z decays

at the Z pole repeat the LEP physics programme in a few minutes...





First look at the physics case of TLEP

PUBLISHED



The TLEP Design Study Working Group

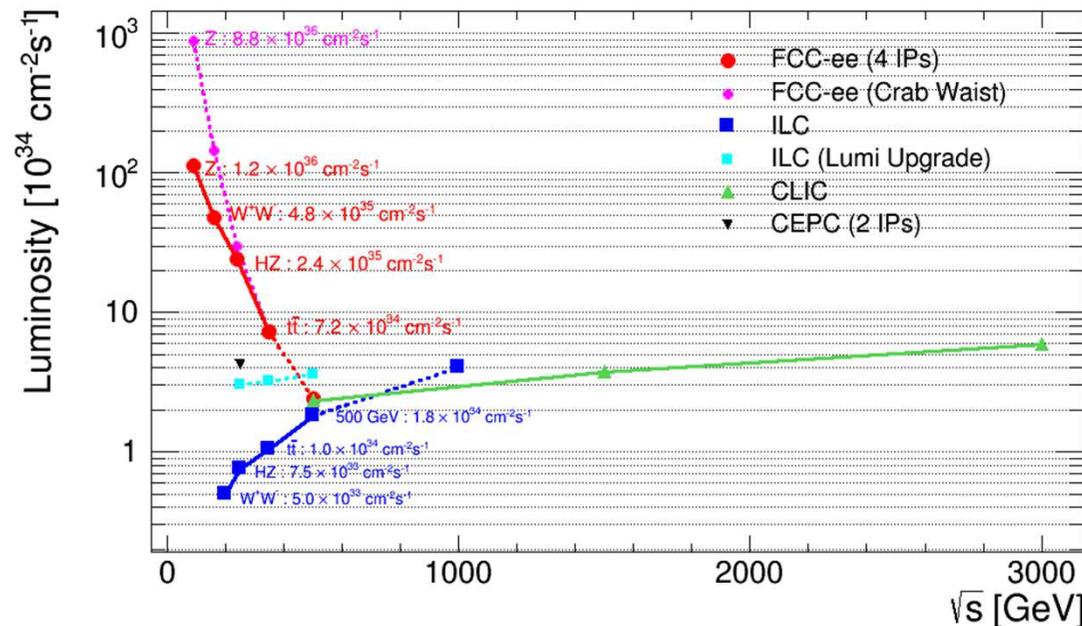
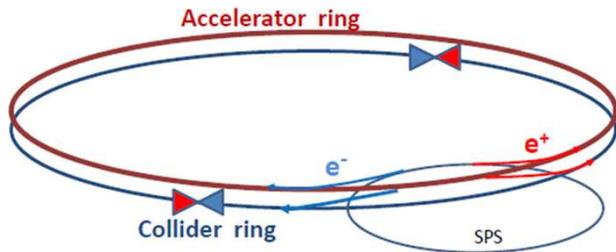
M. Bicer,^a H. Duran Yildiz,^b I. Yildiz,^c G. Coignet,^d M. Delmastro,^d T. Alexopoulos,^e C. Grojean,^f S. Antusch,^g T. Sen,^h H.-J. He,ⁱ K. Potamianos,^j S. Haug,^k A. Moreno,^l A. Heister,^m V. Sanz,ⁿ G. Gomez-Ceballos,^o M. Klute,^o M. Zanetti,^o L.-T. Wang,^p M. Dam,^q C. Boehm,^r N. Glover,^r F. Krauss,^r A. Lenz,^r M. Syphers,^s C. Leonidopoulos,^t V. Ciulli,^u P. Lenzi,^u G. Sguazzoni,^u M. Antonelli,^v M. Boscolo,^v U. Dosselli,^v O. Frasciello,^v C. Milardi,^v G. Venanzoni,^v M. Zobov,^v J. van der Bij,^w M. de Gruttola,^x D.-W. Kim,^y M. Bachtis,^z A. Butterworth,^z C. Bernet,^z C. Botta,^z F. Carminati,^z A. David,^z L. Deniau,^z D. d'Enterria,^z G. Ganis,^z B. Goddard,^z G. Giudice,^z P. Janot,^z J. M. Jowett,^z C. Lourenço,^z L. Malgeri,^z E. Meschi,^z F. Moortgat,^z P. Musella,^z J. A. Osborne,^z L. Perrozzi,^z M. Pierini,^z L. Rinolfi,^z A. de Roeck,^z J. Rojo,^z G. Roy,^z A. Sciabà,^z A. Valassi,^z C.S. Waaijer,^z J. Wenninger,^z H. Woehri,^z F. Zimmermann,^z A. Blondel,^{aa} M. Koratzinos,^{aa} P. Mermod,^{aa} Y. Onel,^{ab} R. Talman,^{ac} E. Castaneda Miranda,^{ad} E. Bulyak,^{ae} D. Porsuk,^{af} D. Kovalskyi,^{ag} S. Padhi,^{ag} P. Faccioli,^{ah} J. R. Ellis,^{ai} M. Campanelli,^{aj} Y. Bai,^{ak} M. Chamizo,^{al} R.B. Appleby,^{am} H. Owen,^{am} H. Maury Cuna,^{an} C. Gracious,^{ao} G. A. Munoz-Hernandez,^{ao} L. Trentadue,^{ap} E. Torrente-Lujan,^{aq} S. Wang,^{ar} D. Bertsche,^{as} A. Gramolin,^{at} V. Telnov,^{at} M. Kado,^{au} P. Petroff,^{au} P. Azzi,^{av} O. Nicrosini,^{aw} F. Piccinini,^{aw} G. Montagna,^{ax} F. Kapusta,^{ay} S. Laplace,^{ay} W. da Silva,^{ay} N. Gizani,^{az} N. Craig,^{ba} T. Han,^{bb} C. Luci,^{bc} B. Mele,^{bc} L. Silvestrini,^{bc} M. Ciuchini,^{bd} R. Cakir,^{be} R. Aleksan,^{bf} F. Couderc,^{bf} S. Ganjour,^{bf} E. Lançon,^{bf} E. Locci,^{bf} P. Schwemling,^{bf} M. Spiro,^{bf} C. Tanguy,^{bf} J. Zinn-Justin,^{bf} S. Moretti,^{bg} M. Kikuchi,^{bh} H. Koiso,^{bh} K. Ohmi,^{bh} K. Oide,^{bh} G. Pauletta,^{bi} R. Ruiz de Austri,^{bj} M. Gouzevitch,^{bk} and S. Chattopadhyay^{bl}

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JHEP01(2014)164





First look at the physics case of TLEP, arXiv:1308.6176v3 scoped the precision measurements:

- Model independent Higgs couplings and invisible width
 - Z mass (0.1 MeV), W mass (0.5 MeV) top mass ($\sim 10 \text{ MeV}$), $\sin^2_{\text{W}}^{\text{eff}}$, R_b , N_ν etc...
 - powerful exploration of new physics with EW couplings up to very high masses
 - importance of luminosity and E_{beam} calibration by beam depolarization up to W pair
- So far: simulations with CMS detector (Higgs) -- or «just» paper studies.

Snapshot of novelties appeared in recent workshops

Higher luminosity prospects at W, Z with **crab-waist**

- sensitivity to right handed (sterile) neutrinos
- s-channel $e^+e^- \rightarrow H(125.2)$ production almost possible (→ monochromators?)
- rare Higgs Z W and top decays, FCNCs etc...
- discovery potential for very small couplings
- precision event generators (Jadach et al)

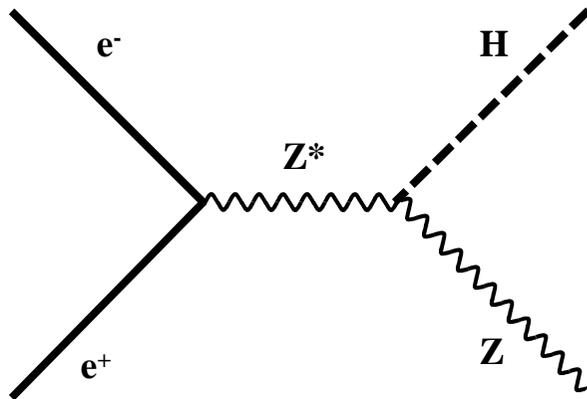
Higgs production mechanism

“higgstrahlung” process close to threshold

Production xsection has a maximum at near threshold ~ 200 fb

$10^{34}/\text{cm}^2/\text{s} \rightarrow 20'000$ HZ events per year. (\sim ILC, muon collider)

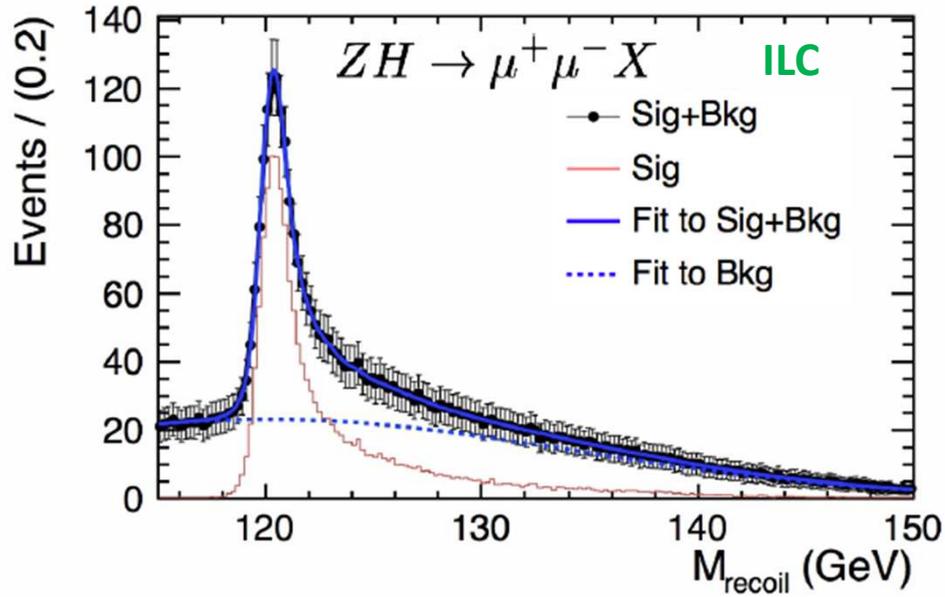
FCC-ee $\rightarrow 400'000$ HZ events a year.



**Z – tagging
by missing mass**

For a Higgs of 125GeV, a centre of mass energy of 240GeV is sufficient

\rightarrow kinematical constraint near threshold for high precision in mass, width, selection purity



Z – tagging by missing mass

total rate $\propto g_{\text{HZZ}}^2$

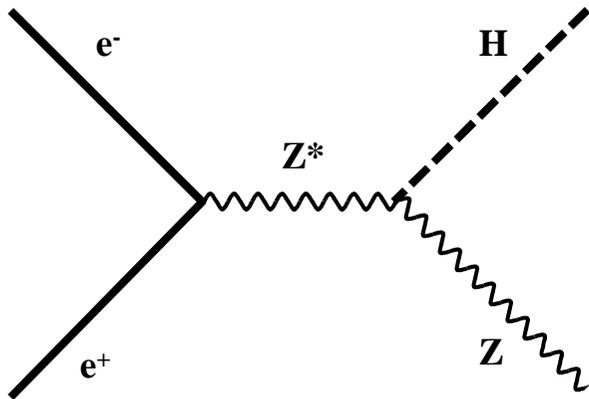
ZZZ final state $\propto g_{\text{HZZ}}^4 / \Gamma_{\text{H}}$

→ measure total width Γ_{H}

empty recoil = invisible width

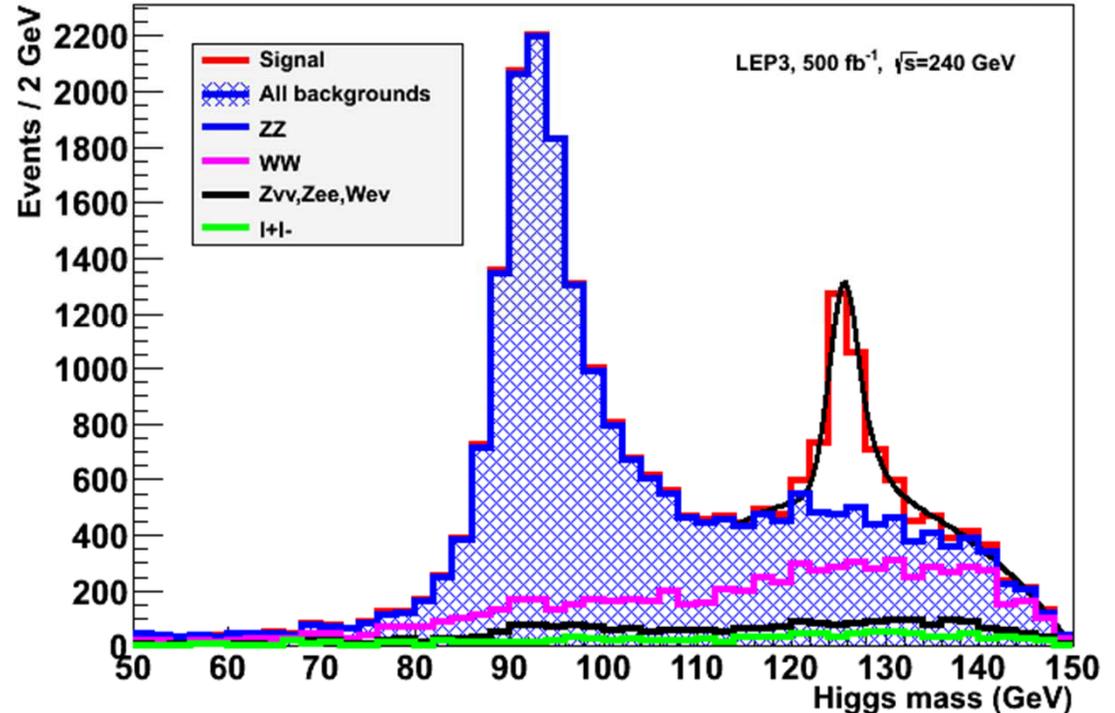
‘funny recoil’ = exotic Higgs decay

easy control below threshold



Z → l+l- with H → anything

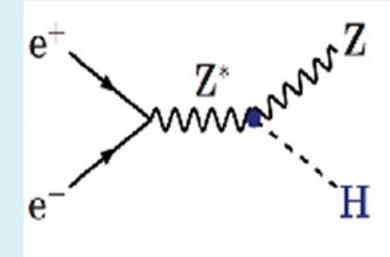
CMS Simulation



Higgs factory

(constrained fit including 'exotic')

	4 IPs	TLEP (2 IPs)
g_{HZZ}	0.05%	(0.06%)
g_{HWW}	0.09%	(0.11%)
g_{Hbb}	0.19%	(0.23%)
g_{Hcc}	0.68%	(0.84%)
g_{Hgg}	0.79%	(0.97%)
$g_{H\tau\tau}$	0.49%	(0.60%)
$g_{H\mu\mu}$	6.2%	(7.6%)
$g_{H\gamma\gamma}$	1.4%	(1.7%)
BR_{exo}	0.16%	(0.20%)



2 10^6 ZH events in 5 years

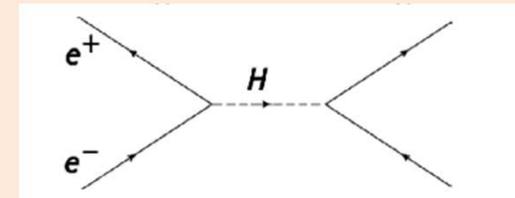
«A tagged Higgs beam».

sensitive to new physics in loops

incl. invisible = (dark matter?)

A big challenge, but unique:

Higgs s-channel production at $\sqrt{s} = m_H$



10^4 events per year.

Very difficult because huge background and beam energy spread $\sim 10 \times \Gamma_H$ limits or signal? monochromators?

Aleksan, D'Enterria, Wojcik

→ **total width**

<1%

HHH (best at FCC-hh)

28% → from HZ thresh

Htt (best at FCC-hh)

13% → from tt thresh

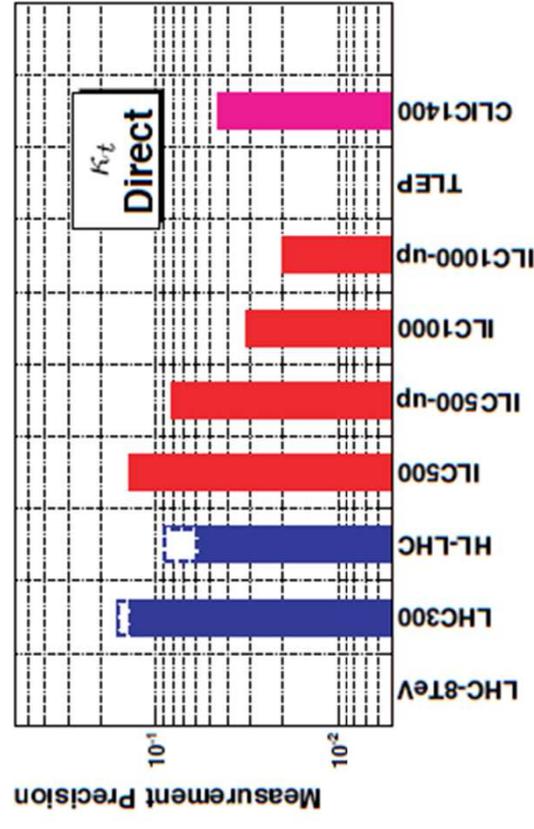
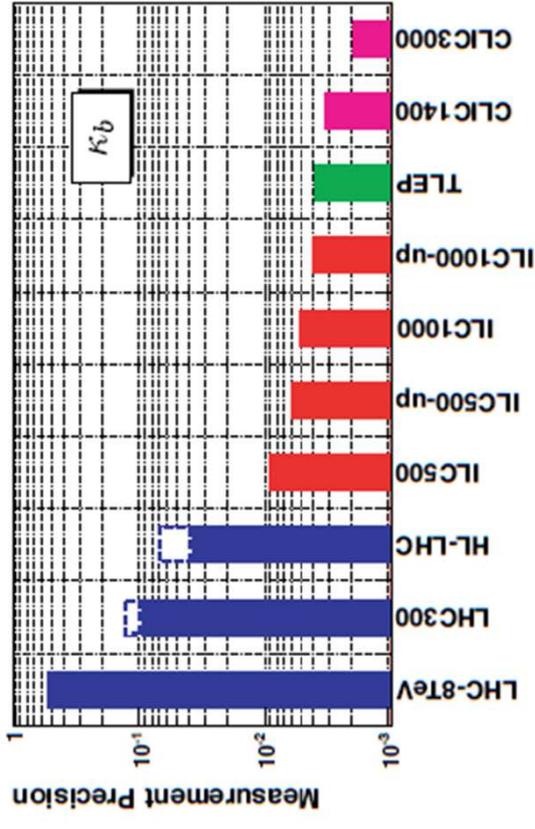
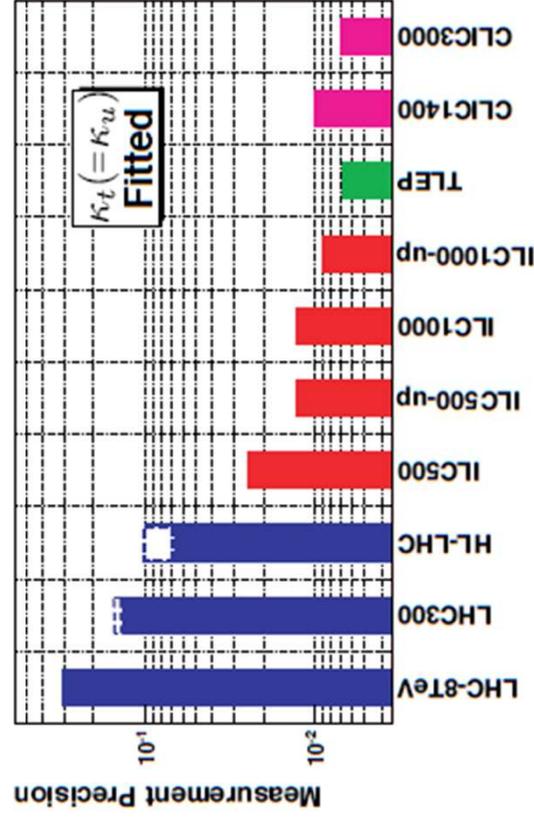
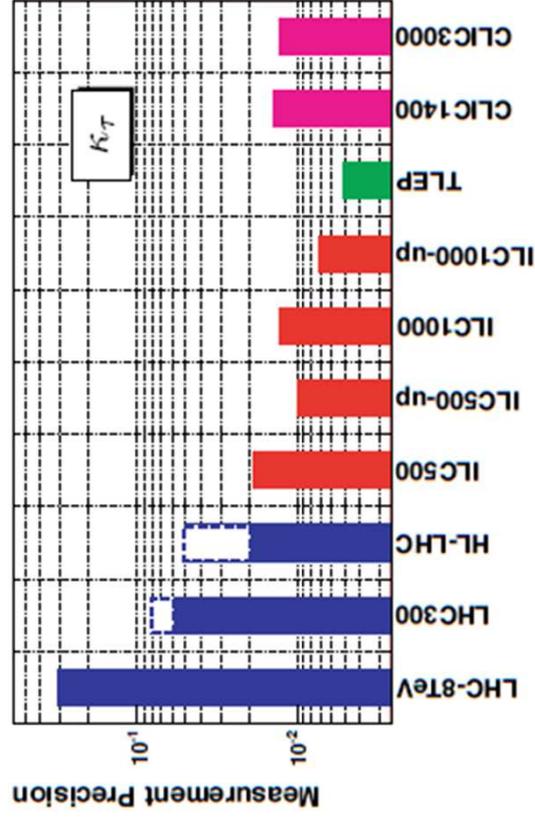


Figure 1-4. Measurement precision on κ_b , κ_τ , and κ_t measured both directly via $t\bar{t}H$ and through global fits at different facilities.

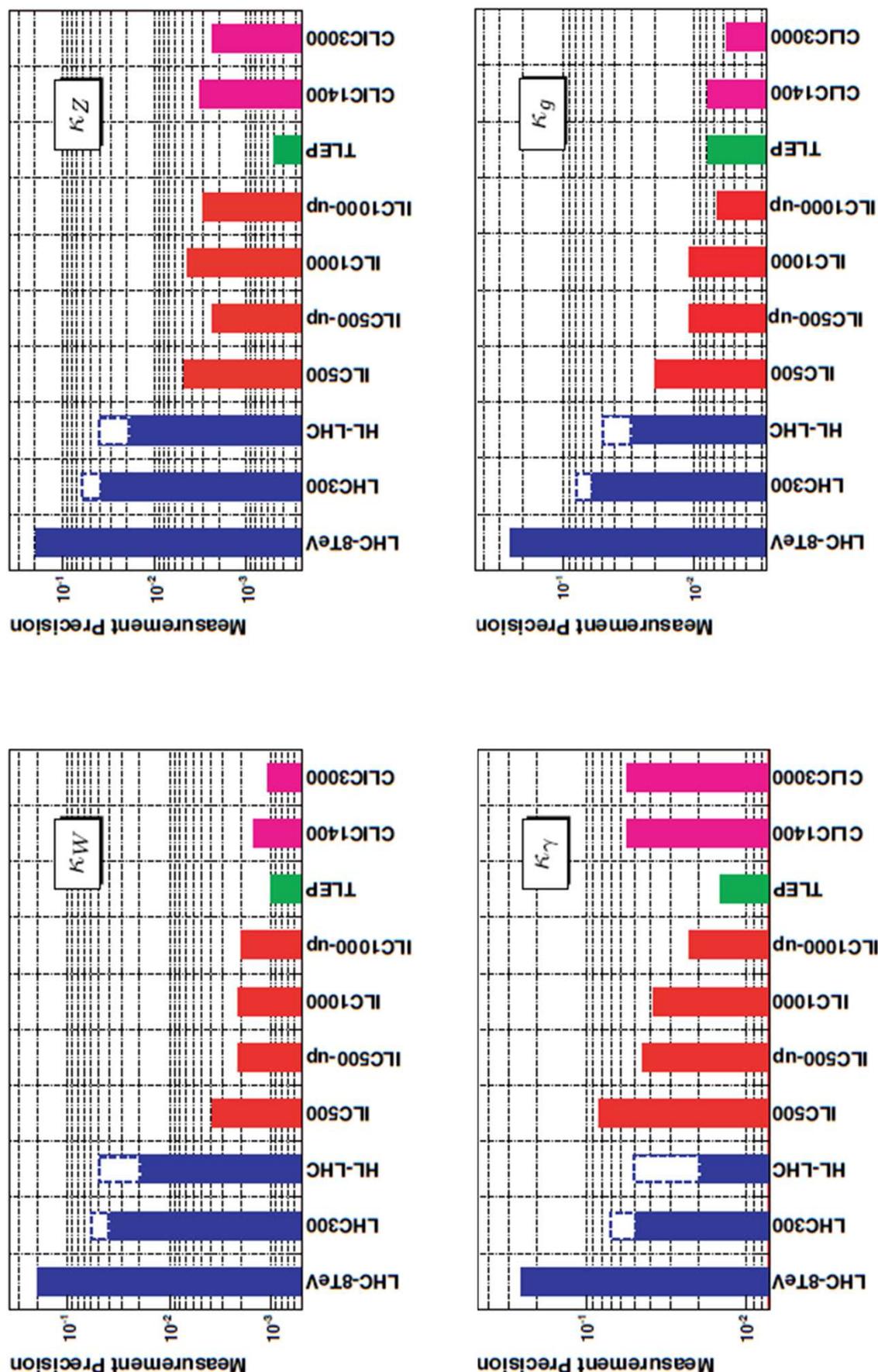
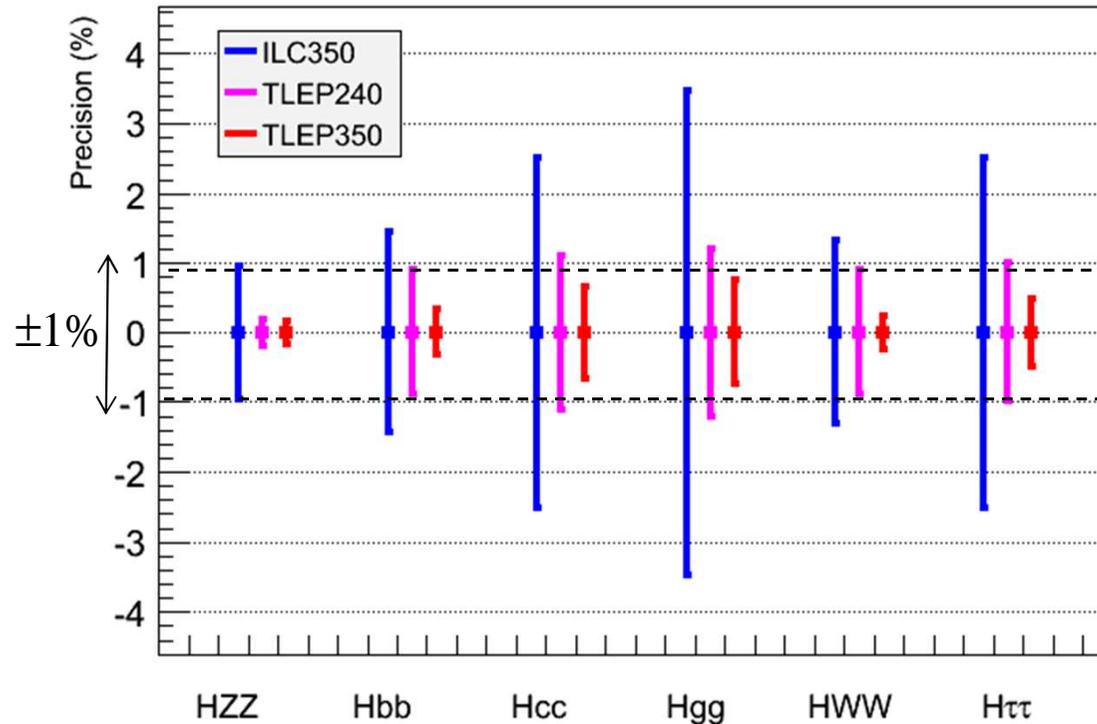


Figure 1-3. Measurement precision on κ_W , κ_Z , κ_γ , and κ_g at different facilities.

Performance Comparison

$$\sigma_{HZ} \propto g_{HZZ}^2, \text{ and } \sigma_{HZ,WW \rightarrow H} \times \text{BR}(H \rightarrow XX) \propto g_{HZZ, HWW}^2 g_{HXX}^2 / \Gamma_H$$

- Same conclusion when Γ_H is a free parameter in the fit



Expected precision on the total width

$\mu^+\mu^-$	ILC350	ILC1000	TLEP240	TLEP350
5%	5%	3%	2%	1%

TLEP : sub-percent precision, BSM Physics sensitivity beyond several TeV

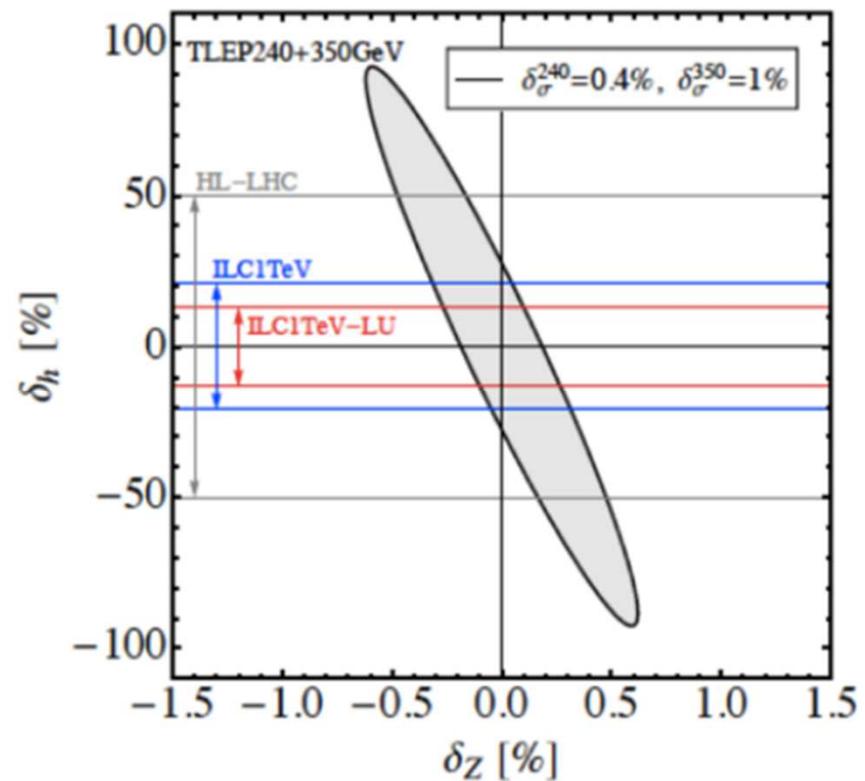
very accurate precision on threshold cross-section sensitive to loop corrections

$$\sigma_{Zh} = \left| \begin{array}{c} e \\ \diagup \\ \text{---} \\ \diagdown \\ e \end{array} \right|^2 + 2 \operatorname{Re} \left[\begin{array}{c} \text{---} \\ \diagup \\ z \\ \diagdown \\ h \end{array} \cdot \left(\begin{array}{c} e^+ \\ \diagup \\ \text{---} \\ \diagdown \\ e^- \end{array} \right) + \left(\begin{array}{c} e^+ \\ \diagup \\ \text{---} \\ \diagdown \\ e^- \end{array} \right) \right]$$

$$\delta_{\sigma}^{240} = 100 (2\delta_Z + 0.014\delta_h) \%$$

[arxiv:1312.3322](https://arxiv.org/abs/1312.3322)

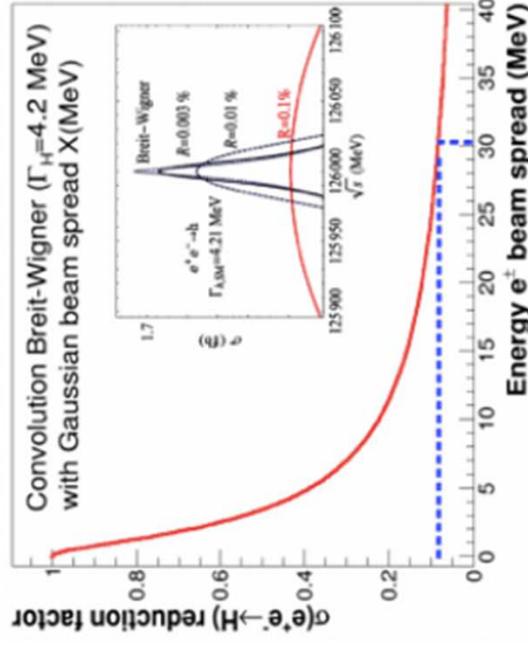
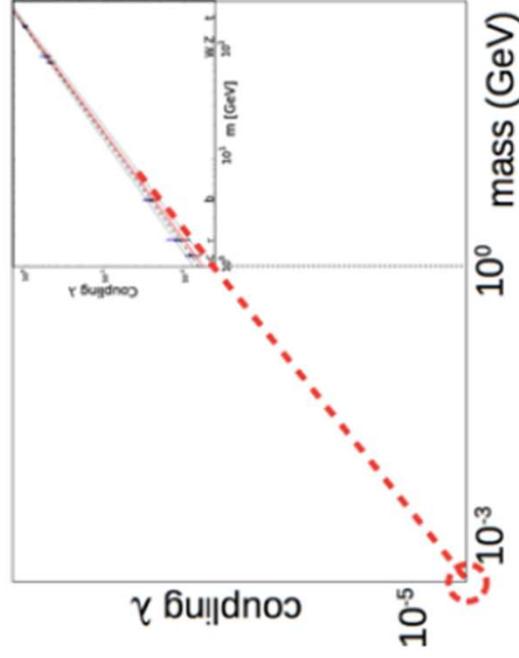
- ➔ Very large datasets at high energy allow extreme precision g_{ZH} measurements
- ➔ Indirect and model-dependent probe of Higgs self-coupling
- ➔ Note, the time axis is missing from the plot



First generation couplings

→ s-channel Higgs production

- Unique opportunity for measurement close to SM sensitivity
- Highly challenging; $\sigma(ee \rightarrow H) = 1.6 \text{ fb}$; 7 Higgs decay channels studied



Preliminary Results

$$L = 10 \text{ ab}^{-1}$$

$$K_e < 2.2 \text{ at } 3\sigma$$

→ Work in progress

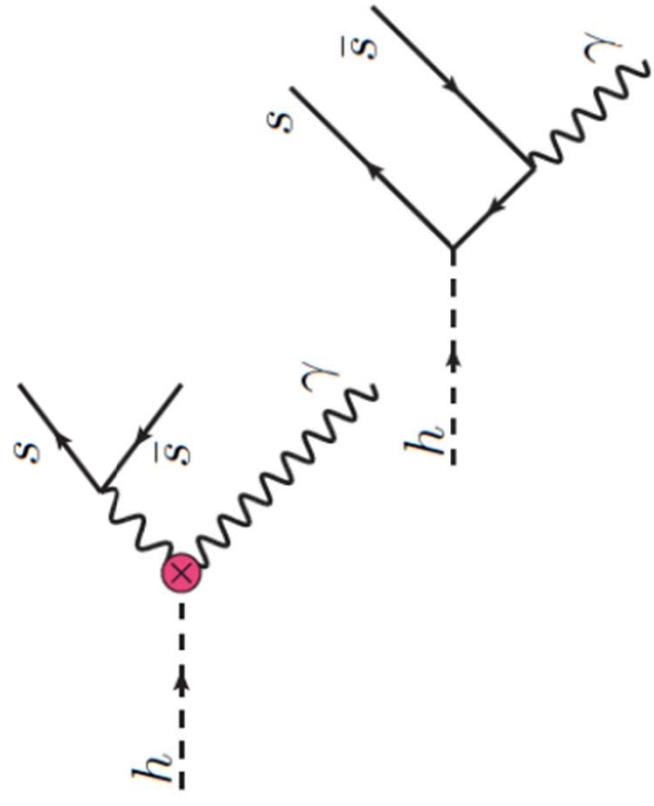
- How large are loop induced corrections? How large are BSM effects?
- Do we need an energy scan to find the Higgs?
- How much luminosity will be available for this measurement? By how much is the luminosity reduced by monochromators?

Exclusive Higgs boson decays

- ➔ First and second generation couplings accessible
 - Study of $\rho\gamma$ channel most promising; expect ~50 evts.
 - Sensitivity to u/d quark Yukawa coupling
 - Sensitivity due to interference

$$\frac{BR_{h \rightarrow \rho\gamma}}{BR_{h \rightarrow b\bar{b}}} = \frac{\kappa_\gamma [(1.9 \pm 0.15)\kappa_\gamma - 0.24\bar{\kappa}_u - 0.12\bar{\kappa}_d]}{0.57\bar{\kappa}_b^2} \times 10^{-5}$$

- ➔ Also interesting to FCC-hh program
- ➔ Alternative $H \rightarrow MV$ decays should be studied ($V = \gamma, W, \text{ and } Z$)

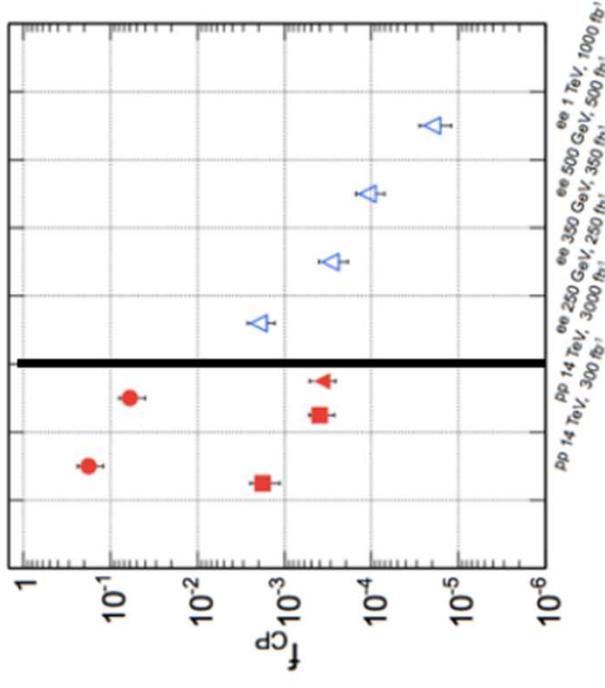


- $H \rightarrow J/\psi \gamma$ ➔ Y_c
- $H \rightarrow \phi \gamma$ ➔ Y_s
- $H \rightarrow \rho \gamma$ ➔ Y_u, Y_d
- $H \rightarrow \omega \gamma$

CP Measurements

- ➔ CP violation can be studied by searching for CP-odd contributions; CP-even already established
- ➔ Snowmass Higgs paper <http://arxiv.org/abs/1310.8361>
- ➔ Higgs to Tau decays of interest
- ➔ More detailed presentation by Felix Yu <http://arxiv.org/abs/1308.1094>

for HWV couplings



$$\mathcal{L}_{hff} \propto h\bar{f}(\cos \Delta + i\gamma_5 \sin \Delta)f$$

Colliders	LHC	HL-LHC	FCCee (1 ab ⁻¹)	FCCee (5 ab ⁻¹)	FCCee (10 ab ⁻¹)
-----------	-----	--------	-----------------------------	-----------------------------	------------------------------

Accuracy(1σ)	25°	8.0°	5.5°	2.5°	1.7°
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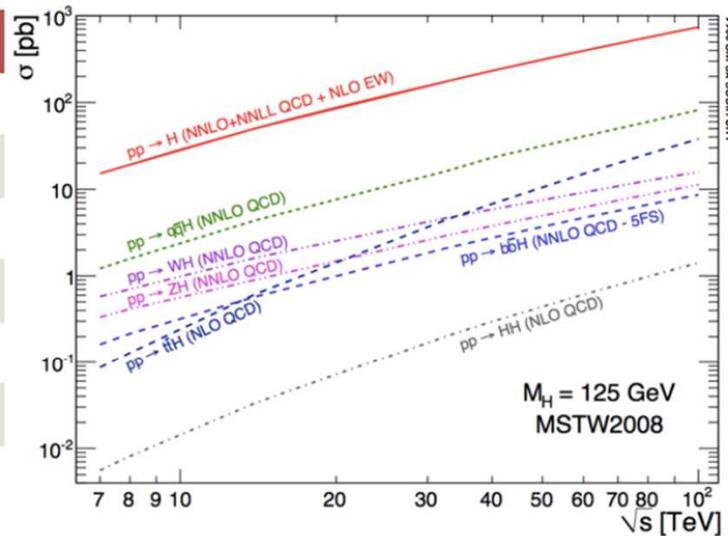
Rare and Exotics Higgs Bosons

- ➔ 2,000,000 ZH events allow for detailed studies of rare and exotic decays
 - requires hadronic and invisible Z decays
 - set requirements for FCC-ee detector
- ➔ Coupling measurements have sensitivity to BSM decays
- ➔ Dedicated studies using specific final states improve sensitivity
- ➔ Example: Higgs to invisible, flavor violating Higgs, and many more
- ➔ Potential at the LHC (and HL-LHC) currently not fully explored
- ➔ Modes with of limited LHC sensitivity are of particular importance to FCC-ee program
 - currently under study
- ➔ FCC-ee might allow precision measurement of exotic Higgs decays
- ➔ Detailed discussion of exotic Higgs decays at [Phys. Rev. D 90, 075004 \(2014\)](#) More from David Curtin

$h \rightarrow \cancel{\chi}_T$
$h \rightarrow 4b$
$h \rightarrow 2b2\tau$
$h \rightarrow 2b2\mu$
$h \rightarrow 4\tau, 2\tau 2\mu$
$h \rightarrow 4j$
$h \rightarrow 2\gamma 2j$
$h \rightarrow 4\gamma$
$h \rightarrow ZZ_{D}, Z_a \rightarrow 4\ell$
$h \rightarrow Z_D Z_D \rightarrow 4\ell$
$h \rightarrow \gamma + \cancel{\chi}_T$
$h \rightarrow 2\gamma + \cancel{\chi}_T$
$h \rightarrow 4 \text{ ISOLATED LEPTONS} + \cancel{\chi}_T$
$h \rightarrow 2\ell + \cancel{\chi}_T$
$h \rightarrow \text{ONE LEPTON-JET} + X$
$h \rightarrow \text{TWO LEPTON-JETS} + X$
$h \rightarrow b\bar{b} + \cancel{\chi}_T$
$h \rightarrow \tau^+ \tau^- + \cancel{\chi}_T$

HIGGS AT FCC-pp

Process	8 TeV	14 TeV	100 TeV
gF	0.38	1	14.7
VBF	0.38	1	18.6
WH	0.43	1	9.7
ZH	0.47	1	12.5
ttH	0.21	1	61
bbH	0.34	1	15
gF to HH	0.24	1	42



Proton-proton
Higgs datasets

LHC
Run I

→
x300-600

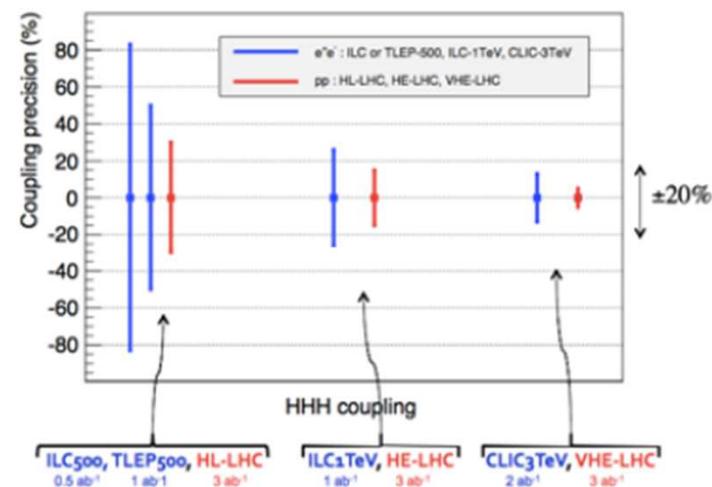
HL
LHC

→
x10-400

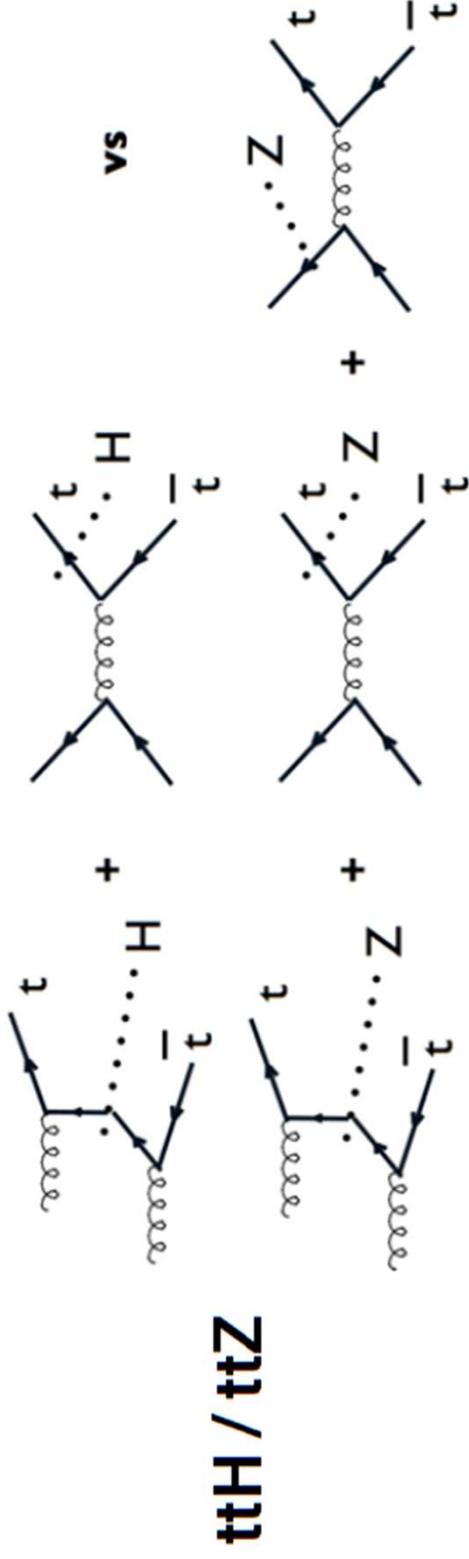
FCC
pp

	HL-LHC	HE-LHC	VLHC
\sqrt{s} (TeV)	14	33	100
$\int \mathcal{L} dt$ (fb ⁻¹)	3000	3000	3000
$\sigma \cdot \text{BR}(pp \rightarrow HH \rightarrow bb\gamma\gamma)$ (fb)	0.089	0.545	3.73
S/\sqrt{B}	2.3	6.2	15.0
λ (stat)	50%	20%	8%

arXiv:1310.8361



➔ ... but also new measurements not possible at the LHC/HL-LHC



➔ Theoretical uncertainties cancel mostly

- PDF (CTEQ 6.6) $\pm 0.5\%$

- Missing higher orders $\pm 1.2\%$

➔ One can not conclude that one can measure the cross section ratio with $\sim 2\%$ ($\delta\lambda_{\text{top}} \approx 1\%$) precision. **More detailed studies are ongoing.**

Table from D. Curtin FCC workshop, Washington, 23-27 March 2015)

- Both lepton and 100 TeV pp colliders are vital for this effort!

Observables at Current + Future Colliders	100 TeV	ILC/TLEP
● producing extra higgs states (incl. superpartners)	✓	
● Exotic Higgs Decays	✓	✓
● Electroweak Precision Observables		✓
● Higgs coupling measurements	✓	✓
● Higgs portal direct production of new states	✓	
● Higgs self coupling measurements	✓	✓
● Zh cross section measurements		✓

Higgs invisible decays

Right handed Neutrinos

etc.. etc..

because of Luminosity FCC-ee (in combination with HL-LHC and/or FCC-hh)
is a very powerful Higgs Factory, but....

FCC-ee is MUCH more than a Higgs Factory!



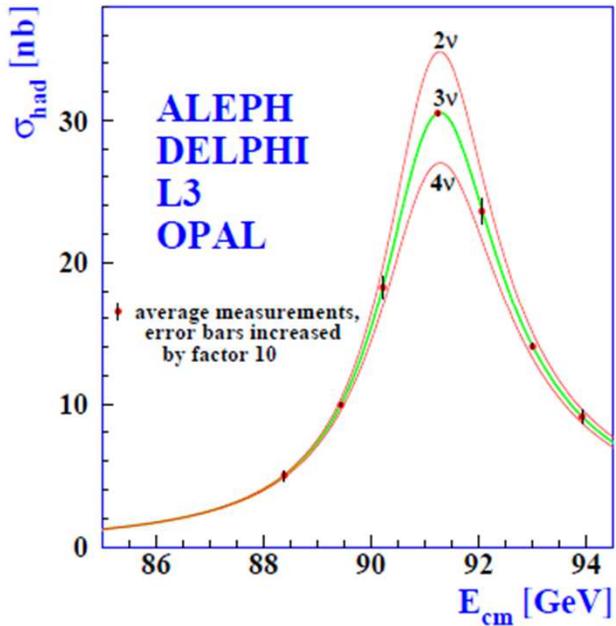
TERA-Z, Oku-W, Megatops

Precision tests of the closure of the Standard Model

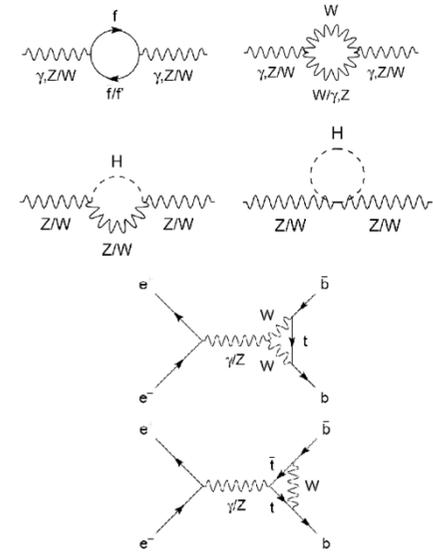
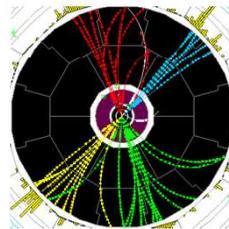
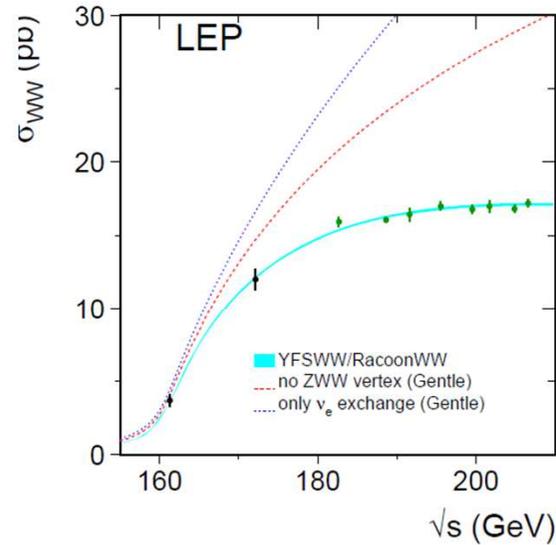


Precision tests of EWSB

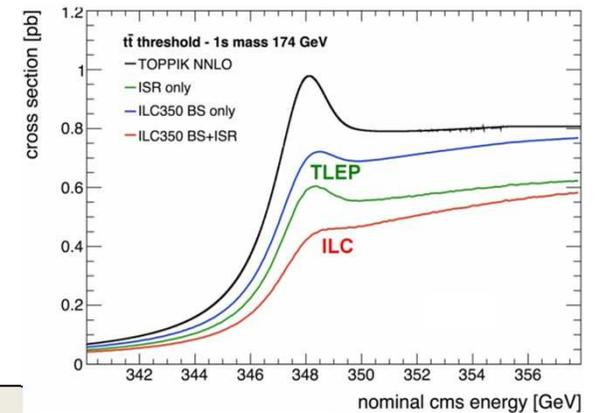
Z pole ssymmetries, lineshape



WW threshold scan



tt threshold scan



TLEP : Repeat the LEP1 physics programme every 15 mn

Transverse polarization up to the WW threshold

➤ Exquisite beam energy determination (10 keV)

Longitudinal polarization at the Z pole

➤ Measure $\sin^2\theta_W$ to $2 \cdot 10^{-6}$ from A_{LR}

➤ Statistics, statistics: 10^{10} tau pairs, 10^{10} bb pairs, QCD and QED studies etc...

4/29/2015

Alain Blondel FCC Future Circular Colliders

Frank Simon



Beam polarization and E-calibration @ FCC-ee

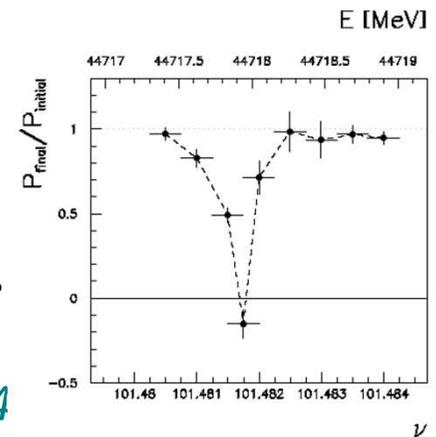
Precise meas of E_{beam} by resonant depolarization

~100 keV each time the meas is made

At LEP transverse polarization was achieved routinely at Z peak.

instrumental in 10^{-3} measurement of the Z width in 1993

led to prediction of top quark mass (179+- 20 GeV) in March 1994



Polarization in collisions was observed (*40% at BBTS = 0.04*)

At LEP beam energy spread destroyed polarization above 60 GeV

$\sigma_E \propto E^2/\sqrt{\rho} \rightarrow$ At FCC-ee transverse polarization up to at least 80 GeV
to go to much higher energies requires spin rotators and siberian snake

FCC-ee: use 'single' bunches to measure the beam energy continuously

no interpolation errors due to tides, ground motion or trains etc...

but saw-toothing must be well understood! require Wigglers to speed up pol. time

<< 100 keV beam energy calibration around Z peak and W pair threshold.

30.04.2015
 $\Delta m_Z \sim 0.1 \text{ MeV}, \Delta \Gamma_Z \sim 0.1 \text{ MeV}, \Delta m_W \sim 0.5 \text{ MeV}$

Example (from Erler & Freytag **PDG 2014**)

$$\Delta\rho = \varepsilon_1 = \alpha(M_Z) \cdot \mathbf{T}$$

$$\varepsilon_3 = 4 \sin^2\theta_W \alpha(M_Z) \cdot \mathbf{S}$$

$\Delta\rho$ today = 0.00040 +- 0.00024

-- is consistent with 0 at 1.7σ

-- is sensitive to non-conventional Higgs bosons (e.g. in SU(2) triplet with 'funny v.e.v.s')

-- is sensitive to Isospin violation such as $m_t \neq m_b$ or **ibid for stop-sbottom**

-- does not decouple!

$$\rho_0 = 1 + \frac{3 G_F}{8\sqrt{2}\pi^2} \sum_i \frac{C_i}{3} \Delta m_i^2, \quad (10.63)$$

where the sum includes fourth-family quark or lepton doublets, (t') or (E^0) , right-handed (mirror) doublets, non-degenerate vector-like fermion doublets (with an extra factor of 2), and scalar doublets such as (\tilde{t}) in Supersymmetry (in the absence of $L-R$ mixing).

Present measurement implies

$$\sum_i \frac{C_i}{3} \Delta m_i^2 \leq (52 \text{ GeV})^2.$$

**Most e.g. SUSY models
have these symmetries
embedded from the start**

Similarly:

$$S = \frac{C}{3\pi} \sum_i \left(t_{3L}(i) - t_{3R}(i) \right)^2,$$



best-of ee-FCC/TLEP #2: Precision EW measts

Asset: -- high luminosity (10^{12} Z decays + 10^8 Wpairs + 10^6 top pairs)
-- exquisite energy calibration up and above WW threshold

target precisions

Quantity	Present precision	Measured from	Statistical uncertainty	Systematic uncertainty
m_Z (keV)	91187500 ± 2100	Z Line shape scan	5 (6) keV	< 100 keV
Γ_Z (keV)	2495200 ± 2300	Z Line shape scan	8 (10) keV	< 100 keV
R_ℓ	20.767 ± 0.025	Z Peak	0.00010 (12)	< 0.001
N_ν	2.984 ± 0.008	Z Peak	0.00008 (10)	< 0.004
N_ν	2.92 ± 0.05	$Z\gamma$, 161 GeV	0.0010 (12)	< 0.001
R_b	0.21629 ± 0.00066	Z Peak	0.000003 (4)	< 0.000060
A_{LR}	0.1514 ± 0.0022	Z peak, polarized	0.000015 (18)	< 0.000015
m_W (MeV)	80385 ± 15	WW threshold scan	0.3 (0.4) MeV	< 0.5 MeV
m_{top} (MeV)	173200 ± 900	$t\bar{t}$ threshold scan	10 (12) MeV	< 10 MeV

Also -- $\Delta \sin^2 \theta_W \approx 10^{-6}$

-- $\Delta \alpha_s = 0.0001$ from W and Z hadronic widths

-- orders of magnitude on FCNCs and rare decays etc. etc.

Design study to establish possibility of corresponding precision theoretical calculations.





best-of ee-FCC/TLEP #2: Precision EW measts

Asset: -- high luminosity (10^{12} Z decays + 10^8 Wpairs + τ pairs)
 -- exquisite energy calibration up and above threshold

Quantity	Present precision	Measured from	Systematic uncertainty
m_Z (keV)	91187500 ± 2100	Z Line scan	< 100 keV
Γ_Z (keV)	2495200 ± 2300	Z Peak	< 100 keV
R_ℓ	20.767 ± 0.025	Z Peak, polarized	0.00010 (12) < 0.001
N_ν	2.984 ± 0.004	WW threshold scan	0.00008 (10) < 0.004
N_ν	2.92 ± 0.01	tt threshold scan	0.0010 (12) < 0.001
R_b	0.21 ± 0.001	Z Peak	0.000003 (4) < 0.000060
A_{LR}		Z peak, polarized	0.000015 (18) < 0.000015
m_W (MeV)		WW threshold scan	0.3 (0.4) MeV < 0.5 MeV
m_{top} (GeV)	< 900	tt threshold scan	10 (12) MeV < 10 MeV

As another example of the importance of precision measurements, the LEP determination of $\alpha_s(m_Z)$ was already able, in association with $\sin^2 \theta_W^{\text{eff}}$, to distinguish between supersymmetric and non-supersymmetric models of grand unification [55–58]. The prospective TLEP accuracies on these quantities would take this confrontation between theory and experiments to a completely new level.

from W and Z hadronic widths
 magnitude on FCNCs and rare decays etc. etc.

study to establish possibility of corresponding precision theoretical calculations.





A Sample of Essential Quantities:

X	Physics	Present precision		TLEP stat Syst Precision	TLEP key	Challenge
M_Z MeV/c ²	Input	91187.5 ± 2.1	Z Line shape scan	0.005 MeV $< \pm 0.1$ MeV	E_cal	QED corrections
Γ_Z MeV/c ²	$\Delta\rho$ (T) (no $\Delta\alpha$!)	2495.2 ± 2.3	Z Line shape scan	0.008 MeV $< \pm 0.1$ MeV	E_cal	QED corrections
R_ℓ	α_s, δ_b	20.767 ± 0.025	Z Peak	0.0001 ± 0.002 - 0.0002	Statistics	QED corrections
N_ν	Unitarity of PMNS, sterile ν 's	2.984 ± 0.008	Z Peak Z+ γ (161 GeV)	0.00008 ± 0.004 0.0004-0.001	->lumi meast Statistics	QED corrections to Bhabha scat.
R_b	δ_b	0.21629 ± 0.00066	Z Peak	0.000003 $\pm 0.000020 - 60$	Statistics, small IP	Hemisphere correlations
A_{LR}	$\Delta\rho, \epsilon_3, \Delta\alpha$ (T, S)	0.1514 ± 0.0022	Z peak, polarized	± 0.000015	4 bunch scheme	Design experiment
M_W MeV/c ²	$\Delta\rho, \epsilon_3, \epsilon_2, \Delta\alpha$ (T, S, U)	80385 ± 15	Threshold (161 GeV)	0.3 MeV < 1 MeV	E_cal & Statistics	QED corections
m_{top} MeV/c ²	Input FCC Future Circular Colliders	173200 ± 900	Threshold scan	10 MeV	E_cal & Statistics	Theory limit at 100 MeV?



Input from Physics to the accelerator design

0. Nobody complains that the luminosity is too high (the more you get, the more you want)

1. Do we need polarized beams?

-1- transverse polarization:

continuous beam Energy calibration with resonant depolarization

central to the precision measurements of m_Z , m_W , Γ_Z

requires 'single bunches'

a priori doable up to W energies -- workarounds exist above (e.g. γZ events)

large ring with small emittance offers *a priori* excellent prospects

need wigglers; simulations ongoing (E. Gianfelice, M. Koratzinos)

-2- longitudinal polarization requires spin rotators and is very difficult at high energies

-- We recently found that it is not necessary to extract top couplings (Janot, Azzi)

-- improves Z peak measurements *if loss in luminosity is not too strong*
but brings no information that is not otherwise accessible

2. What energies are necessary?

-- in addition to Z, W, H and top listed the following are being considered

-- $e^+e^- \rightarrow H(125.2)$ (requires monochromatization A. Faus) (under study)

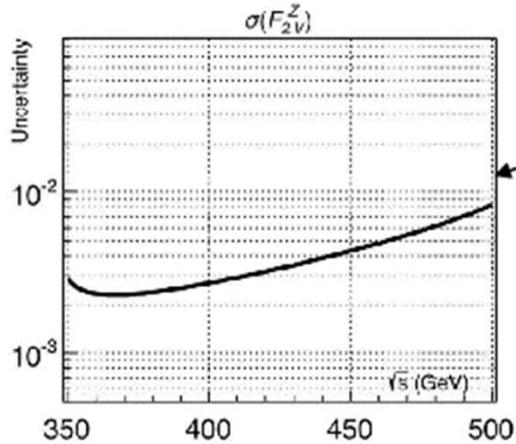
-- e^+e^- at ~ 70 GeV (Z- γ interference)

-- e^+e^- at top threshold + $< \sim 20$ GeV for top couplings (E_{max} up to 180 -185 GeV)

-- no obvious case for going to 500 GeV

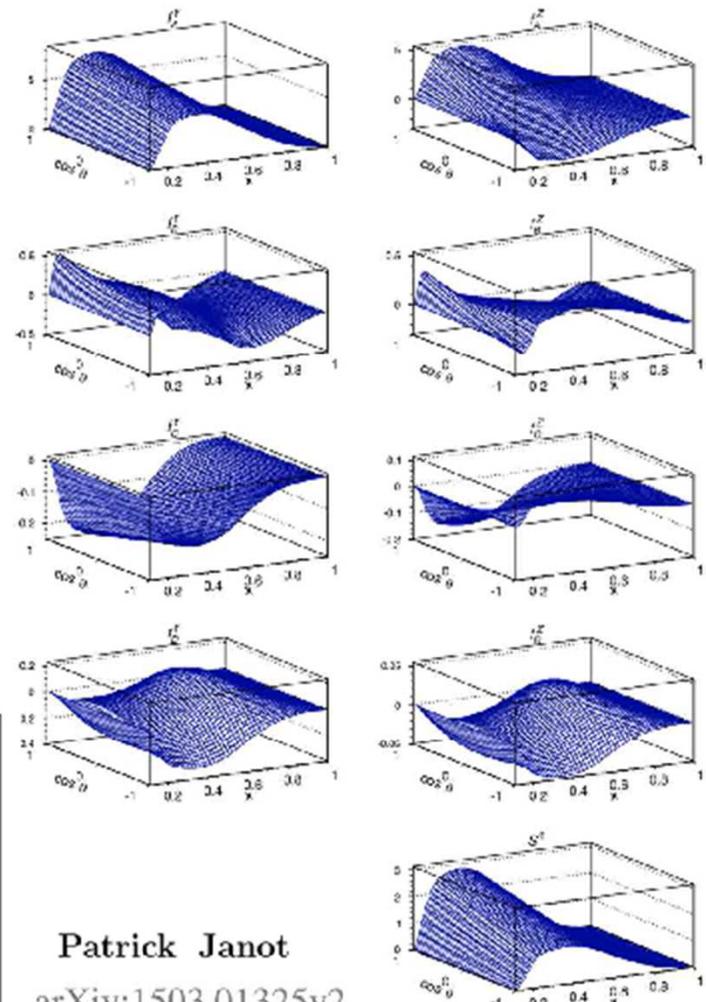
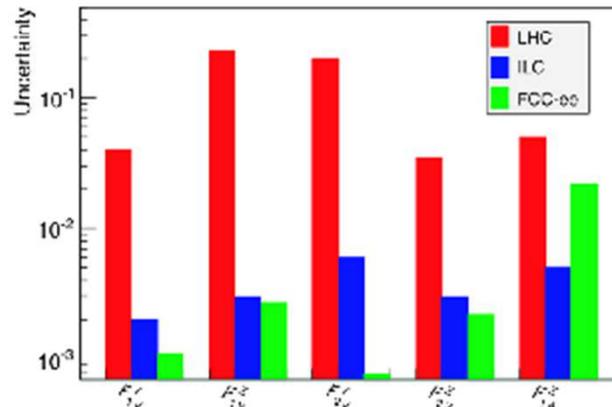


Determination of top-quark EW couplings via measurement of **top-quark polarization**.
 In semileptonic decays, fit to lepton momentum vs scattering angle



Typically best sensitivity just above production threshold
 Momenta up to: 175 GeV

Patrizia Azzi:
 Top physics at FCC-ee



Patrick Janot
 arXiv:1503.01325v2





Next plans for FCC-ee

-- quality of FCC-ee experiments is intimately related to accelerator performance

-- available energy points

-- Luminosities

reinforce work hand-in-hand

-- beam polarization and energy calibration

-- knowledge of other beam parameters (e.g. energy spread vs Z width)

-- we can (mostly out of LEP experience) project fairly well the experimental precisions

-- sometimes they are vertiginously small

$\Delta \sin^2\theta_W^{\text{eff}} = 5 \cdot 10^{-6}$, $\Delta m_Z = 0.1 \text{ MeV}$ $\Delta \Gamma_Z = 0.1 \text{ MeV}$ $\Delta m_W = 0.5 \text{ MeV}$ $\Delta \sigma_{ZH} / \sigma_{ZH} \sim 10^{-3}$ etc...

careful revisiting will be necessary.

-- full use of precision measurements requires a considerable improvement in the theory calculations

-- for the measurements themselves (e.g. Full two loops exponentiated for the QED ISR)

-- for the interpretation; full three loop calculations for EWRCs

and on inputs ($\Delta \alpha_{\text{QED}}(m_Z)$ *Was, Gluza, Heynemeyer, Kuhn, Frietas, Jadach, Ward..*)





Rare decays

-- FCNC: $Z \rightarrow e + \tau$ $Z \rightarrow \mu + \tau$

-- Heavy neutrinos (*they must be somewhere!*)

neutrino counting and search for explicit $Z \rightarrow \nu\text{-}N$
(with $N \rightarrow \nu X$ or $e X'$ and possibly displayed vertices)

-- other final states with single or double photons and jets

-- flavour physics...

-- and many others ($Z \rightarrow \gamma\gamma$ etc)

-- How far can one go with 10^{12} or 10^{13} Z decays?



But Where Is Everybody?

Nima



But Where Is Everybody?

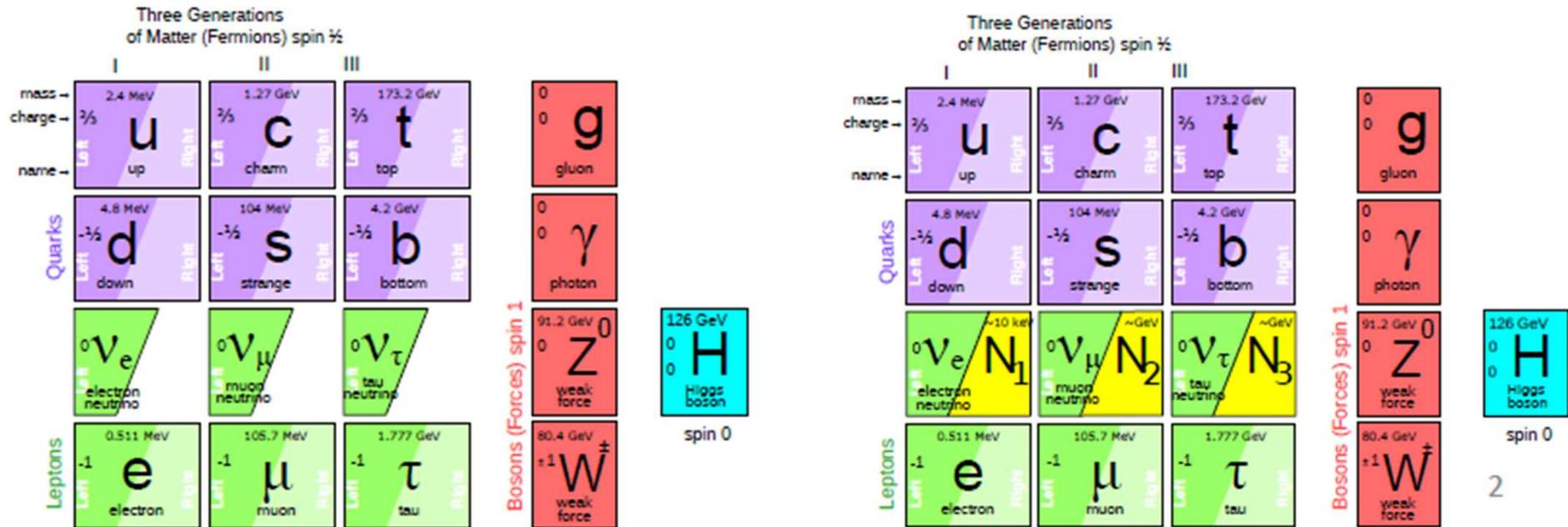
Nima

At higher masses -- or at smaller couplings?



THE STANDARD MODEL IS COMPLETE

but... at least 3 pieces are still missing!



neutrinos have mass...

and this very probably implies new degrees of freedom

➔ Right-Handed, Almost «Sterile» (*very small couplings*) Neutrinos completely unknown masses (meV to ZeV), nearly impossible to find.

.... but could perhaps explain all: DM, BAU, ν -masses



Electroweak eigenstates

$\begin{pmatrix} e \\ \nu_e \end{pmatrix}_L$	$\begin{pmatrix} \mu \\ \nu_\mu \end{pmatrix}_L$	$\begin{pmatrix} \tau \\ \nu_\tau \end{pmatrix}_L$	$(e)_R$	$(\mu)_R$	$(\tau)_R$	Q= -1
			$(\nu_e)_R$	$(\nu_\mu)_R$	$(\nu_\tau)_R$	Q= 0
$I = 1/2$			$I = 0$			

**Right handed neutrinos
 are singlets
 no weak interaction
 no EM interaction
 no strong interaction**

**can't produce them
 can't detect them
 -- so what? --**



Mass eigenstates

See-saw in a general way :

$$\mathcal{L} = \frac{1}{2} (\bar{\nu}_L, \bar{N}_R^c) \begin{pmatrix} 0 & m_D \\ m_D^T & M_R \end{pmatrix} \begin{pmatrix} \nu_L^c \\ N_R \end{pmatrix}$$

$$M_R \neq 0$$

$$m_D \neq 0$$

Dirac + Majorana
mass terms

$$\tan 2\theta = \frac{2m_D}{M_R - 0} \ll 1$$

$$m_\nu = \frac{1}{2} \left[(0 + M_R) - \sqrt{(0 - M_R)^2 + 4m_D^2} \right] \simeq -m_D^2/M_R$$

$$M = \frac{1}{2} \left[(0 + M_R) + \sqrt{(0 - M_R)^2 + 4m_D^2} \right] \simeq M_R$$

general formula

if $m_D \ll M_R$

$M_R = 0$
 $m_D \neq 0$
Dirac only, (like e- vs e+):

	ν_L	ν_R	$\bar{\nu}_R$	$\bar{\nu}_L$
$I_{\text{weak}} =$	$1/2$	0	$1/2$	0

4 states of equal masses
Some have $I=1/2$ (active)
Some have $I=0$ (sterile)

$M_R \neq 0$
 $m_D = 0$
Majorana only

	ν_L	$\bar{\nu}_R$
$I_{\text{weak}} =$	$1/2$	$1/2$

2 states of equal masses
All have $I=1/2$ (active)

$M_R \neq 0$
 $m_D \neq 0$
Dirac + Majorana see-saw

	ν_L	N_R	$\bar{\nu}_R$	N_L
$I_{\text{weak}} =$	$1/2$	0	$1/2$	0

4 states, 2 mass levels
m1 have $I=1/2$ (~active)
m2 have $I=0$ (~sterile)



There even exists a scenario that explains everything: the ν MSM

Shaposhnikov et al

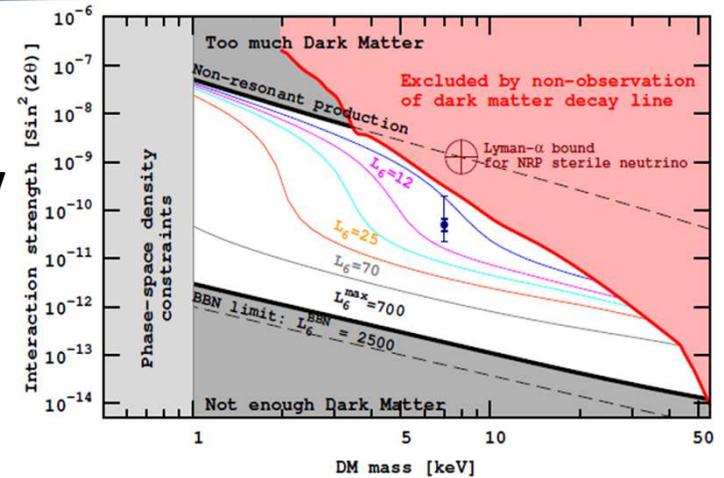


N_2, N_3

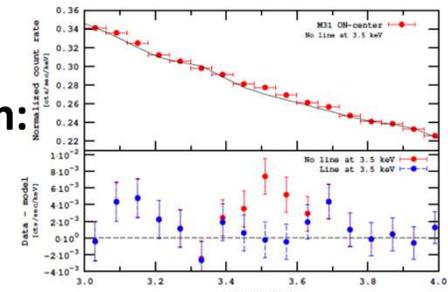
can generate Baryon Asymmetry of Universe
if $m_{N_2, N_3} > 140$ MeV

N_1

constrained:
mass: 1-50 keV
mixing :
 10^{-7} to 10^{-13}
decay time:
 $\tau_{N_1} > \tau_{\text{Universe}}$



Andromeda galaxy (zoom 3-4 keV)



$N_1 \rightarrow \nu \gamma$
may have been seen:
arxiv:1402:2301
arxiv:1402.4119

(or not)



Manifestations of right handed neutrinos

one family see-saw :

$$\theta \approx (m_D/M)$$

$$m_\nu \approx \frac{m_D^2}{M}$$

$$m_N \approx M$$

$$|U|^2 \propto \theta^2 \approx m_\nu / m_N$$

$$\nu = \nu_L \cos\theta - N^c_R \sin\theta$$

$$N = N_R \cos\theta + \nu_L^c \sin\theta$$

what is produced in W, Z decays is:

$$\nu_L = \nu \cos\theta + N \sin\theta$$

ν = light mass eigenstate

N = heavy mass eigenstate

$\neq \nu_L$, active neutrino

which couples to weak inter.

and $\neq N_R$, which doesn't.

- mixing with active neutrinos leads to various observable consequences
 - if very light (eV), possible effect on neutrino oscillations
 - if in keV region (dark matter), monochromatic photons from galaxies with $E=m_N/2$
- possibly measurable effects at High Energy
 - If N is heavy it will decay in the detector (not invisible)
 - PMNS matrix unitarity violation and deficit in Z «invisible» width
 - Higgs and Z visible exotic decays $H \rightarrow \nu_i \bar{N}_i$ and $Z \rightarrow \nu_i \bar{N}_i$, $W \rightarrow l_i \bar{N}_i$
 - also in charm and b decays via $W^* \rightarrow l_i \bar{N}_i$
 - violation of unitarity and lepton universality in Z, W or τ decays
 - etc... etc...
- Couplings are small (m_ν / m_N) (but who knows?) and generally out of reach of hadron colliders (but this deserves to be revisited for detached vertices @LHC, HL-LHC, FCC-hh)



At the end of LEP:

Phys.Rept.427:257-454,2006

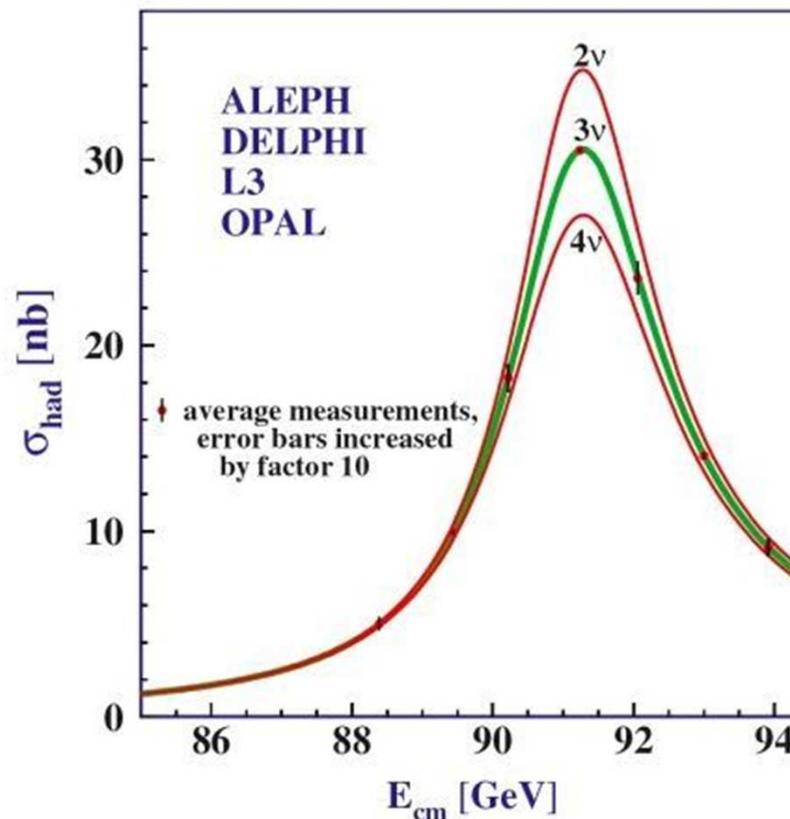
$$N_\nu = 2.984 \pm 0.008$$

- 2 σ :^) !!

This is determined from the Z line shape scan and dominated by the measurement of the hadronic cross-section at the Z peak maximum →

The dominant systematic error is the theoretical uncertainty on the Bhabha cross-section (0.06%) which represents an error of ± 0.0046 on N_ν

Improving on N_ν by more than a factor 2 would require a large effort to improve on the Bhabha cross-section calculation!



Neutrino counting at TLEP



given the very high luminosity, the following measurement can be performed

$$N_\nu = \frac{\frac{\gamma Z(inv)}{\gamma Z \rightarrow ee, \mu\mu}}{\frac{\Gamma_\nu}{\Gamma_{e, \mu}} (SM)}$$

The common **γ tag** allows cancellation of systematics due to photon selection, luminosity etc. The others are extremely well known due to the availability of $O(10^{12})$ Z decays.

The full sensitivity to the number of neutrinos is restored, and the theory uncertainty on $\frac{\Gamma_\nu}{\Gamma_e} (SM)$ is very very small.

A good measurement can be made from the data accumulated at the WW threshold where $\sigma(\gamma Z(inv)) \sim 4$ pb for $|\cos\theta_\gamma| < 0.95$

**161 GeV (10^7 s) running at $1.6 \times 10^{35}/\text{cm}^2/\text{s} \times 4$ exp $\rightarrow 3 \times 10^7$ $\gamma Z(inv)$ evts, $\Delta N_\nu = 0.0011$
 adding 5 yrs data at 240 and 350 GeV $\Delta N_\nu = 0.0008$**

A better point may be 105 GeV (20pb and higher luminosity) may allow $\Delta N_\nu = 0.0004$?



RHASnu's production in Z decays

Production:

$$BR(Z^0 \rightarrow \nu_m \bar{\nu}) = BR(Z^0 \rightarrow \nu \bar{\nu}) |U|^2 \left(1 - \frac{m_{\nu_m}^2}{m_{Z^0}^2}\right)^2 \left(1 + \frac{1}{2} \frac{m_{\nu_m}^2}{m_{Z^0}^2}\right)$$

multiply by 2 for anti neutrino and add contributions of 3 neutrino species (with different $|U|^2$)

Decay

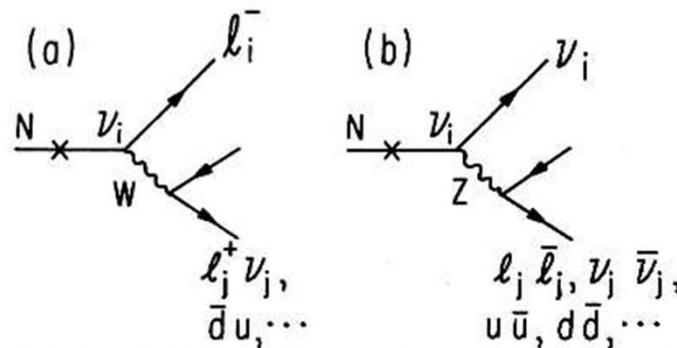


FIG. 2. Typical decays of a neutral heavy lepton via (a) charged current and (b) neutral current. Here the lepton l_i denotes $e, \mu, \text{ or } \tau$.

Decay length:

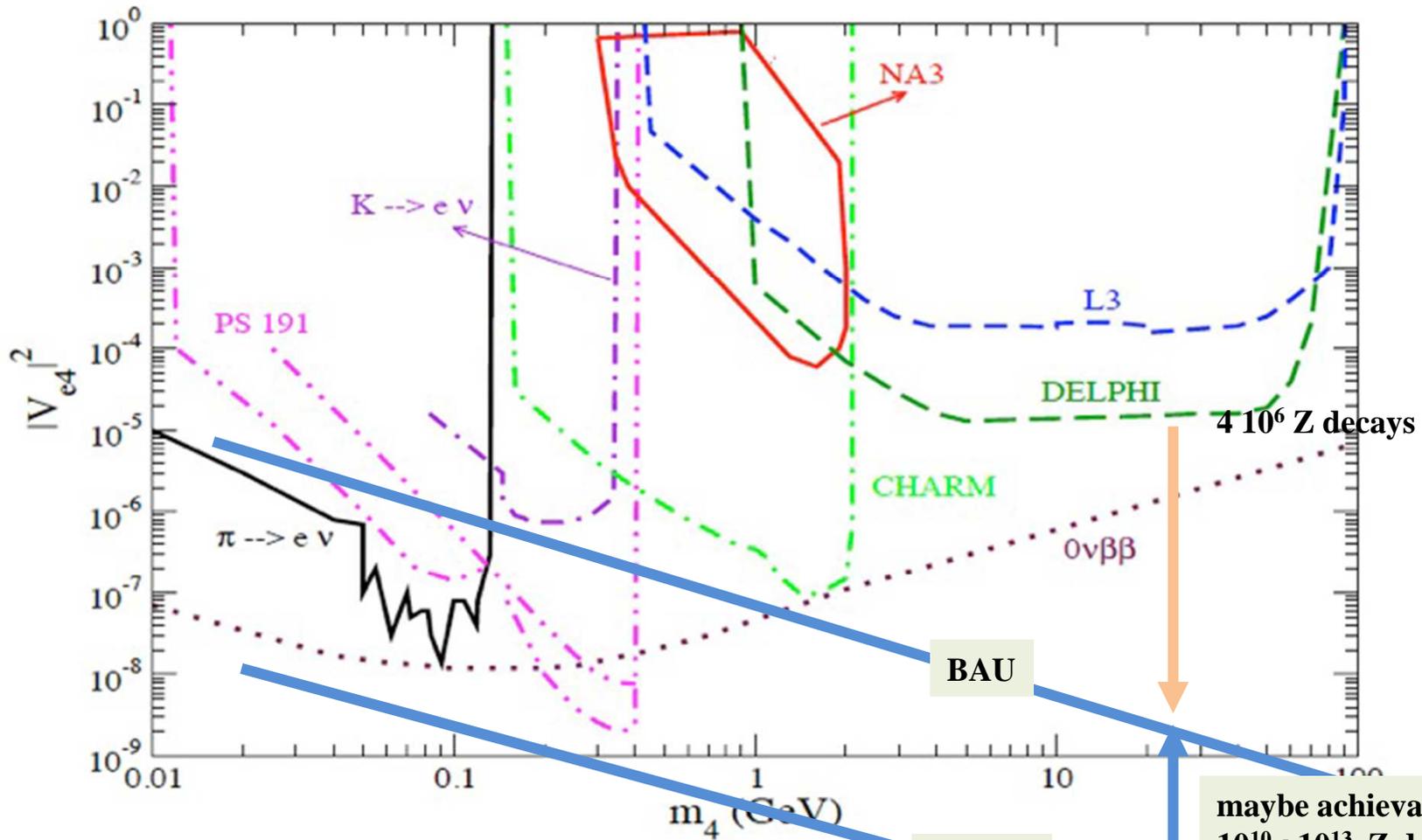
$$L \approx \frac{3 \text{ cm}}{|U|^2 (m_{\nu_m} (\text{GeV}/c^2))^6}$$

NB CC decay always leads to ≥ 2 charged tracks

Backgrounds : four fermion: $e^+e^- \rightarrow W^{*+} W^{*-}$ $e^+e^- \rightarrow Z^*(\nu\nu) + (Z/\gamma)^*$

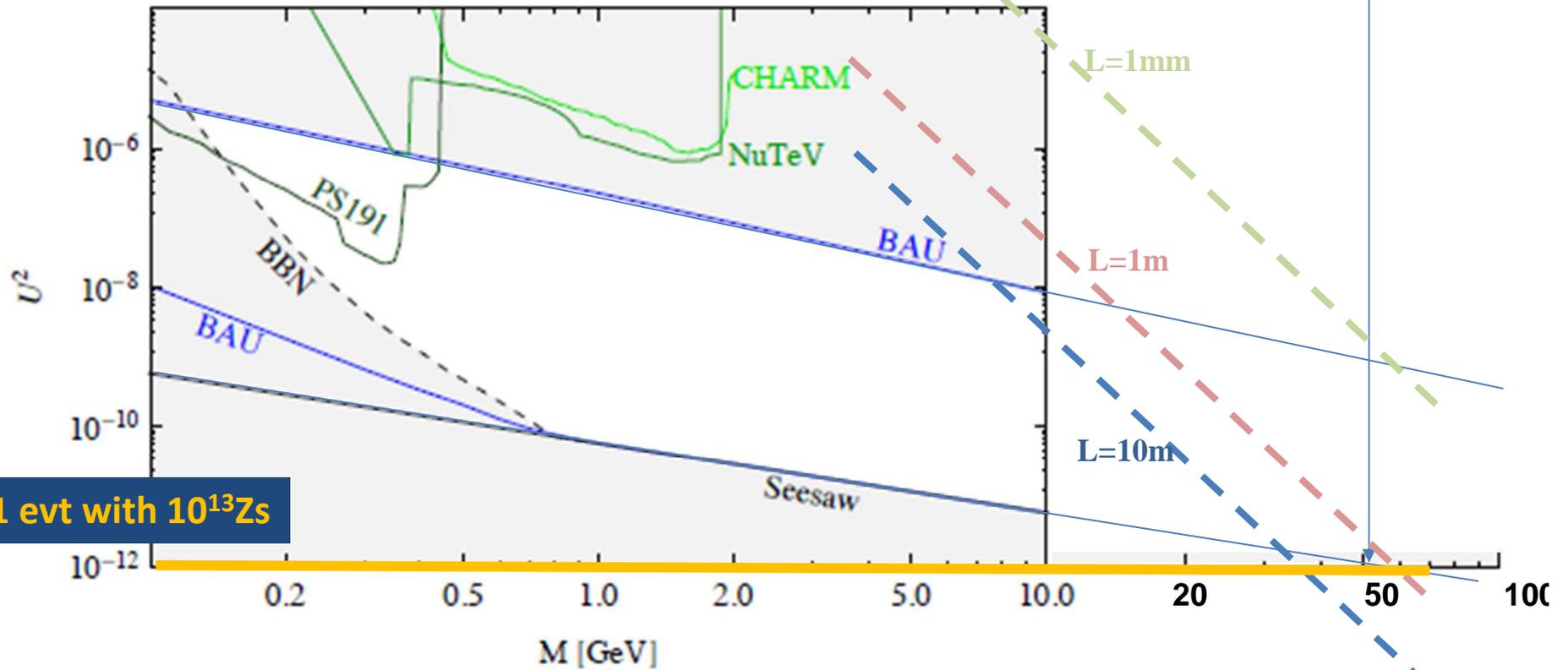


Order-of-magnitude extrapolation of existing limits



Decay length

Interesting region
 $|U|^2 \sim 10^{-9}$ to 10^{-12} @ 50 GeV



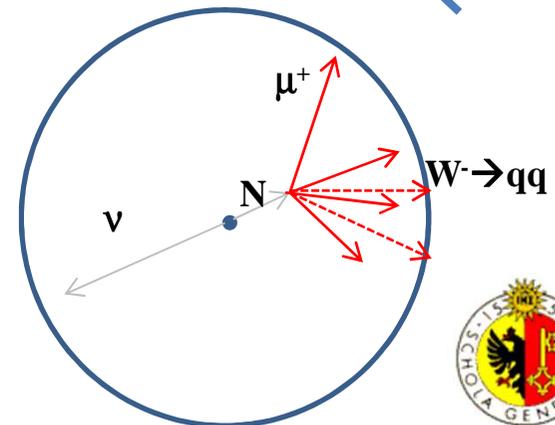
~1 evt with 10^{13} Zs

heavy neutrino mass $\sim M$

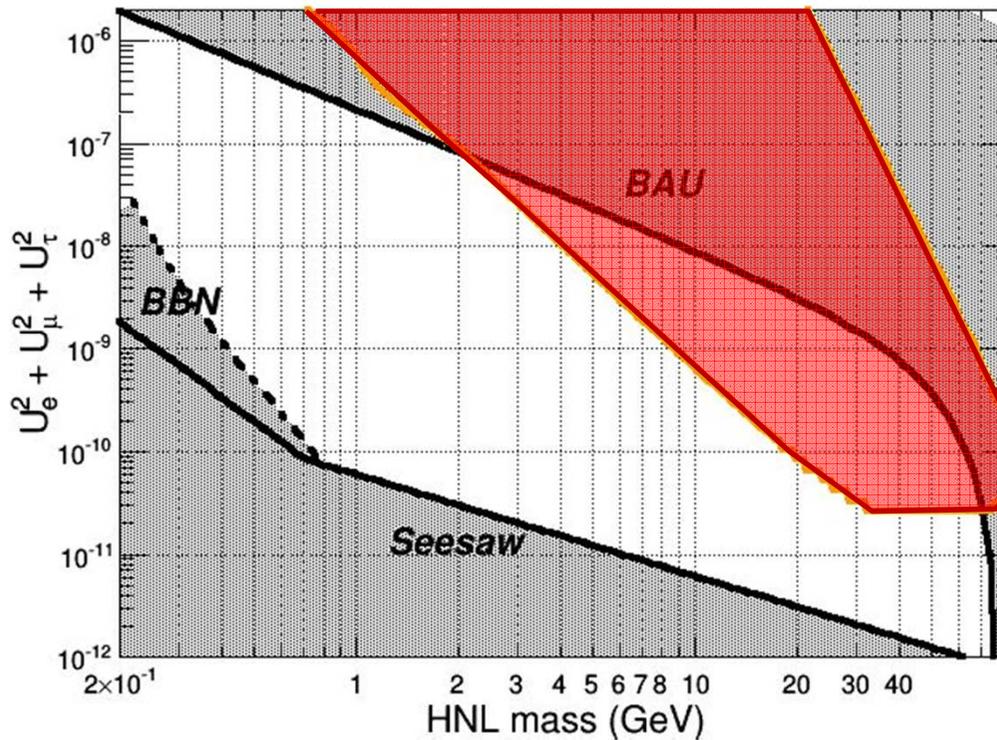
a large part of the interesting region will lead to detached vertices

... \rightarrow very strong reduction of background!

Exact reach domain will depend on detector size and details of displaced vertex efficiency & background



TLEP expected sensitivity to HNL (NH)



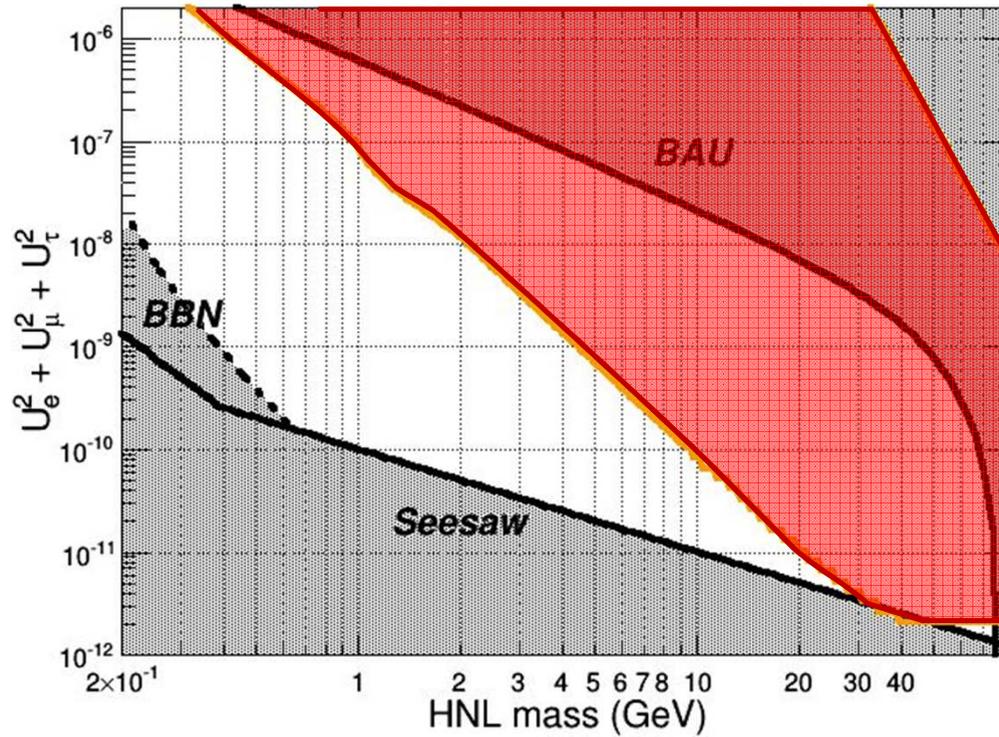
$N_2 = 10^{12}$ $1\text{mm} < L < 1\text{m}$

-  region of interest
-  FCC-ee sensitivity

A.B, Elena Graverini, Nicola Serra, Misha Shaposhnikov



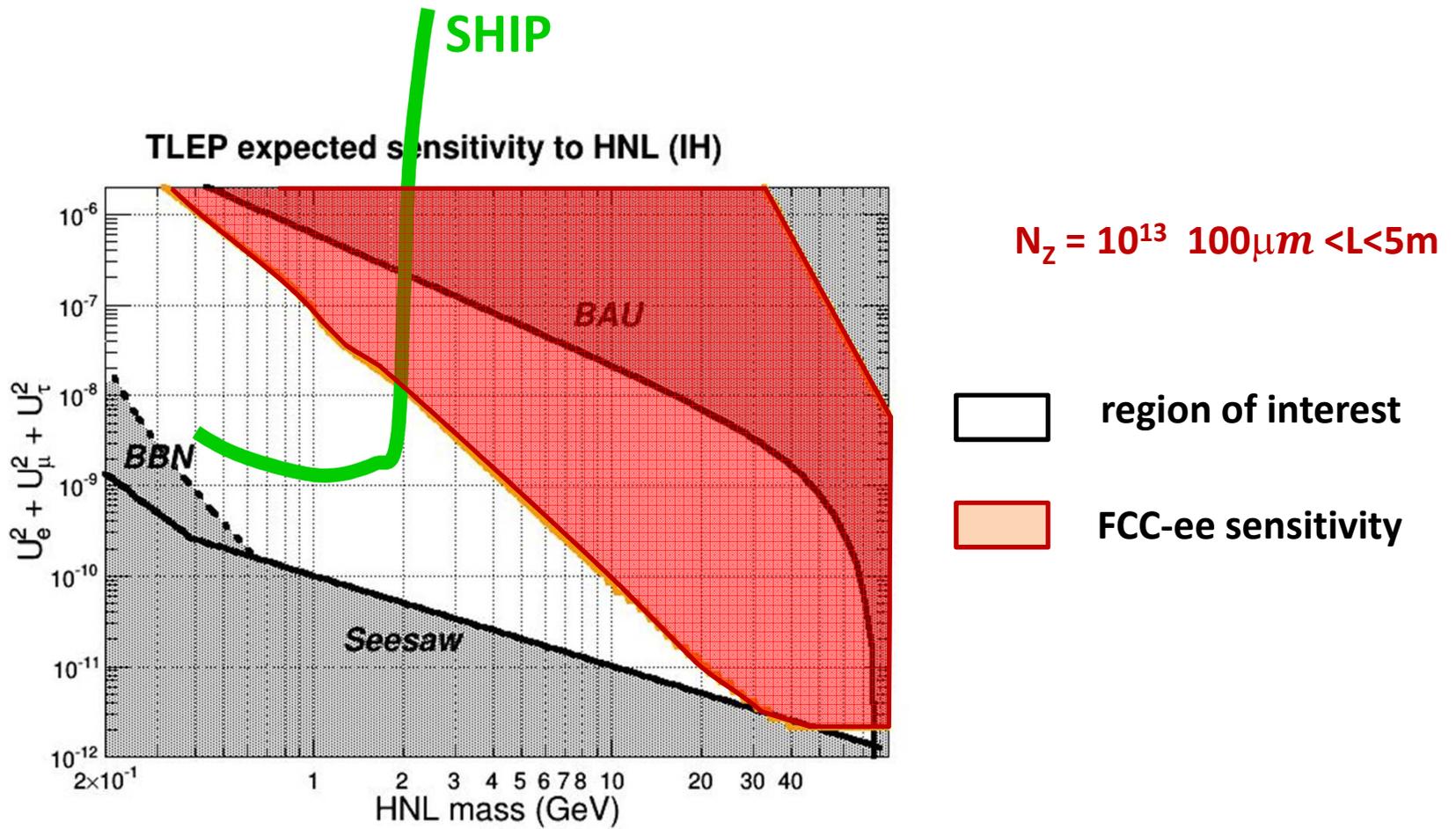
TLEP expected sensitivity to HNL (IH)



$N_z = 10^{13}$ $100\mu m < L < 5m$

- region of interest
- FCC-ee sensitivity







FCC-ee is a wonderful first step towards the Ultimate goal of a 100 TeV hadron collider and this is one of the reasons it is attractive.

But...

FCC-ee is MUCH more than a launching pad!





What we **believe** now and work to **demonstrate** in a few years:

The combination of the FCC machines offers outstanding discovery potential by exploration of new domains of

-- precision

and

-- direct search,

both at high energy and at very small couplings

join us!

<http://cern.ch/fcc-ee>

<http://espace2013.cern.ch/fcc/Pages/Science.aspx>





Kick-off Meeting of the Future Circular Colliders Design Study
12 - 15 February 2014, University of Geneva / Switzerland

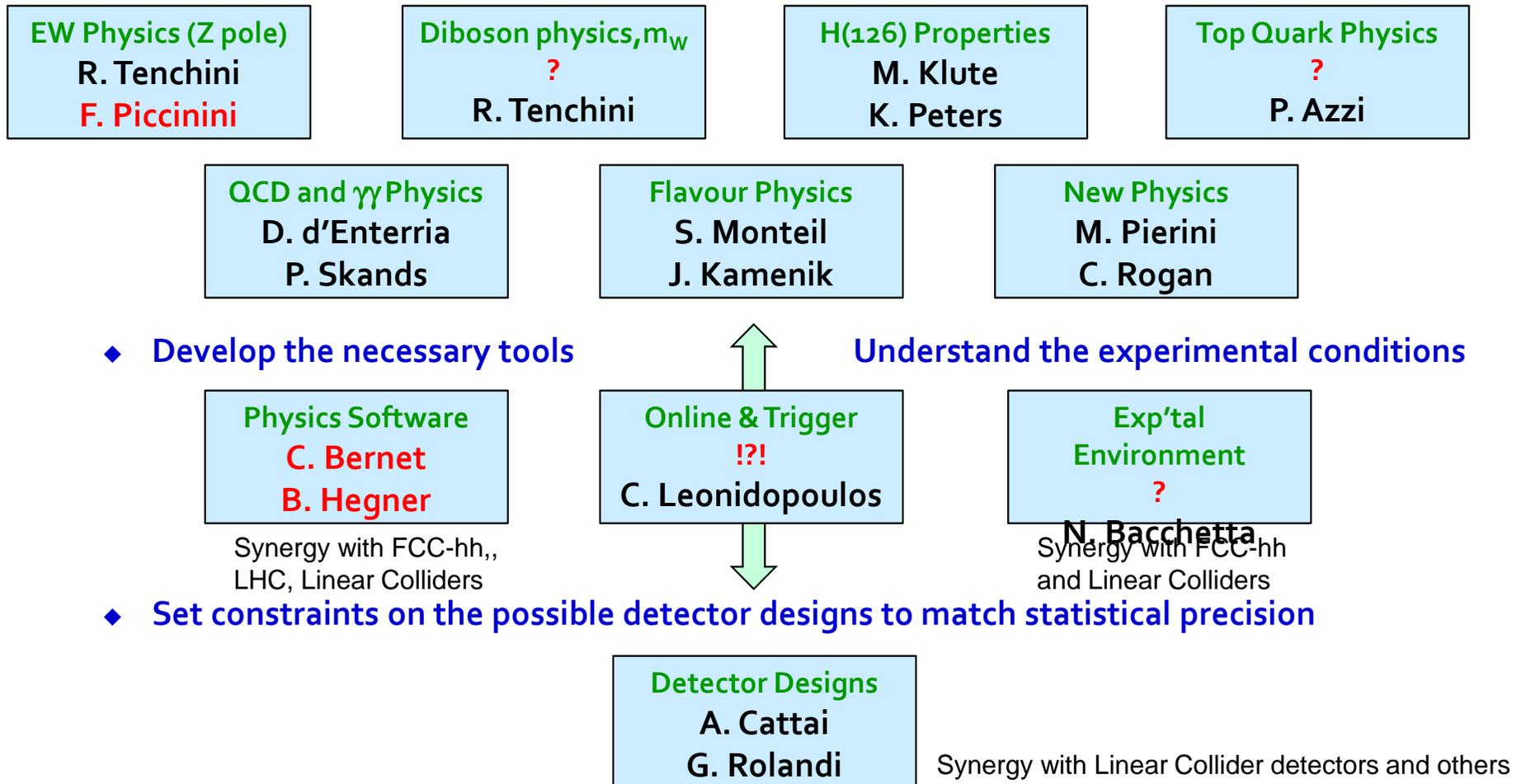
photo by Michael.Hoch@cern.ch

330 registered participants

Experimental Studies: Conveners

- Coordinators **A. Blondel, P. Janot**

- Study the properties of the Higgs and other particles with unprecedented precision



- Develop the necessary tools

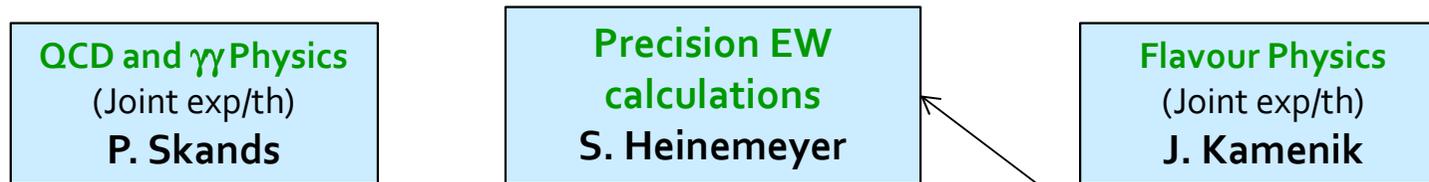
Understand the experimental conditions

- Set constraints on the possible detector designs to match statistical precision

Phenomenological Studies: Conveners

□ Coordinators: J. Ellis, C. Grojean

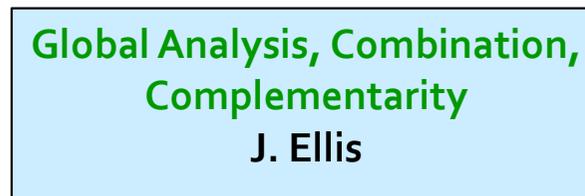
- ◆ Set up a long-term programme to match theory predictions to experimental precisions



- ◆ Understand how new physics would show up in precision measurements, and in searches for rare decays (Z , W , t , H , b , c , τ , ...) and rare processes



- ◆ Set up the framework for global fits and understand the complementarity with other colliders (LHC, FCC-hh, in particular)



some REFERENCES
for right handed
neutrino searches

PHYSICAL REVIEW D

VOLUME 29, NUMBER 11

1 JUNE 1984

Extending limits on neutral heavy leptons

Michael Gronau*

Department of Physics, Syracuse University, Syracuse, New York 132

FLAVOUR(267104)-ERC-23 TUM-HEP 850/12 SISSA 25/2012/EP CFTP/12-013

arxiv:1208.3654

Higgs Decays in the Low Scale
Type I See-Saw Model

C. Garcia Cely^{a)}, A. Ibarra^{a)}, E. Molinaro^{b)} and S. T. Petcov^{c,d)} 1

theories of the electroweak strong interactions. At present
and mixings with ordinary neutrinos of these leptons are ve

The Role of Sterile Neutrinos in Cosmology and
Astrophysics

Alexey Boyarsky[†], Oleg Ruchayskiy[†] and Mikhail Shaposhnikov

The ν MSM, Dark Matter and Neutrino Masses

Takehiko Asaka, Steve Blanchet, and Mikhail Shaposhnikov

Institut de Théorie des Phénomènes Physiques,

Phys.Lett.B631:151-156,2005

arXiv:hep-ph/0503065

CH-1015 Lausanne, Switzerland

(2005)

talks by Maurizio Pierini (BSM), Manqi Ruan (Higgs)
Roberto Tenchini (Top & Precision) tomorrow,
posters tonight at Future accelerator session



PUBLISHED FOR SISSA BY SPRINGER

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PUBLISHED: January 29, 2014

First look at the physics case of TLEP



arxiv:1308.6176

The TLEP Design Study Working Group

M. Bicer,^{a)} H. Duran Yildiz,^{b)} I. Yildiz,^{c)} G. Coignet,^{d)} M. Delmastro,^{d)} T. Alexopoulos,^{e)}
C. Grojean,^{f)} S. Antusch,^{g)} T. Sen,^{h)} H.-J. He,ⁱ⁾ K. Potamianos,^{j)} S. Haug,^{k)}
A. Moreno,^{l)} A. Heister,^{m)} V. Sanz,ⁿ⁾ G. Gomez-Ceballos,^{o)} M. Klute,^{p)} M. Zanetti,^{q)}
L.-T. Wang,^{r)} M. Dam,^{s)} C. Boehm,^{t)} N. Glover,^{u)} F. Krauss,^{v)} A. Lenz,^{w)} M. Syphers,^{x)}

CERN-PPE/96-195

18 December 1996

Search for Neutral Heavy Leptons
Produced in Z Decays

DELPHI Collaboration

FCC design study and FCC-ee <http://cern.ch/fcc-ee>
and presentations at FCC-ee physics workshop
<http://indico.cern.ch/event/313708/>
and arXiv:1411.5230v2 [hep-ex] 6 Dec 2014

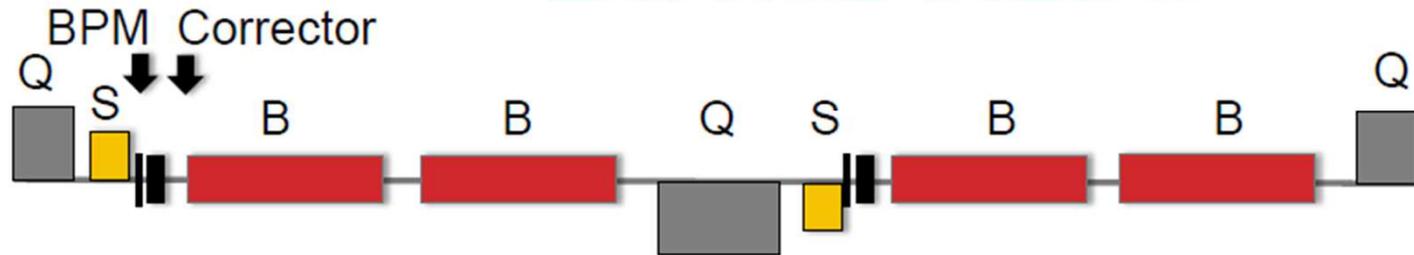
The Search for Heavy Majorana Neutrinos

Anupama Atre^{1,2}, Tao Han^{2,3,4}, Silvia Pascoli⁵, Bin Zhang^{4*}

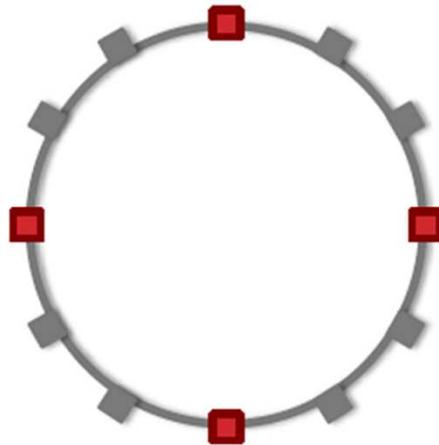


LATTICE V12B-S

arc cell layout



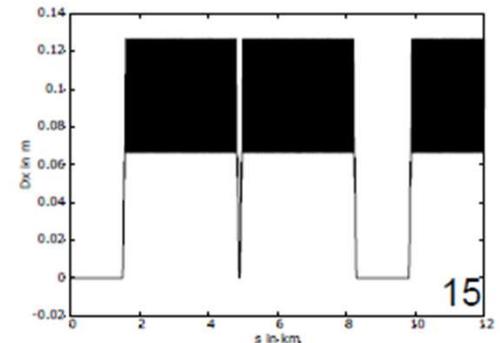
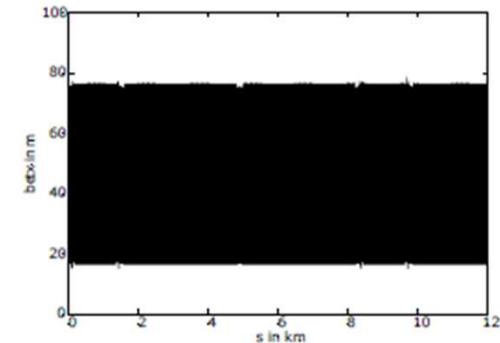
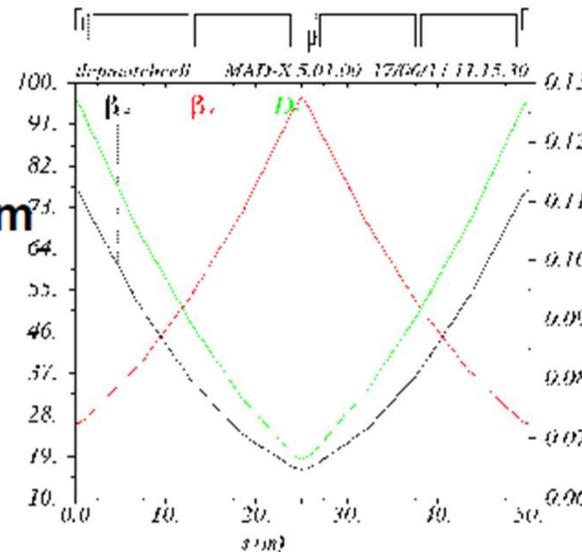
B = bending magnet, Q = quadrupole, S = sextupole



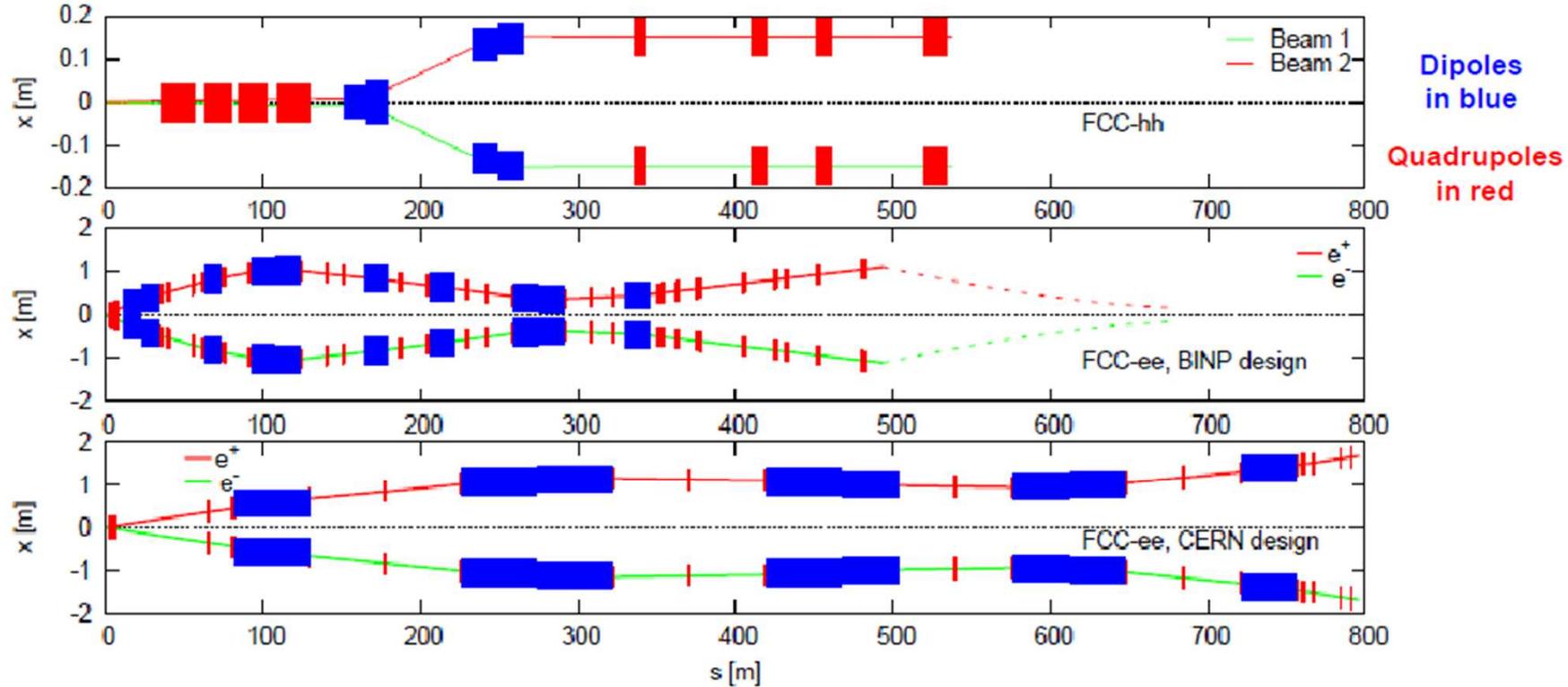
Circumference: 100 km
Arc length: 2 × 3.4 km
Straight section: 1.5 km

B. Harer, B. Holzer

FODO cell optics
 cell length 50 m



IR layouts



- Tunnel transverse width of both FCC-ee designs ~3-4 m.
- Additional length is required to bend beams back, plus room for RF.
- Synchrotron rad. power per IP: **CERN 140 kW, BINP 1400 kW.**
 - *Optimum between length and power loss to be identified !*
 - *93 km racetrack IR straights of 1400 m may be too short for ee !*



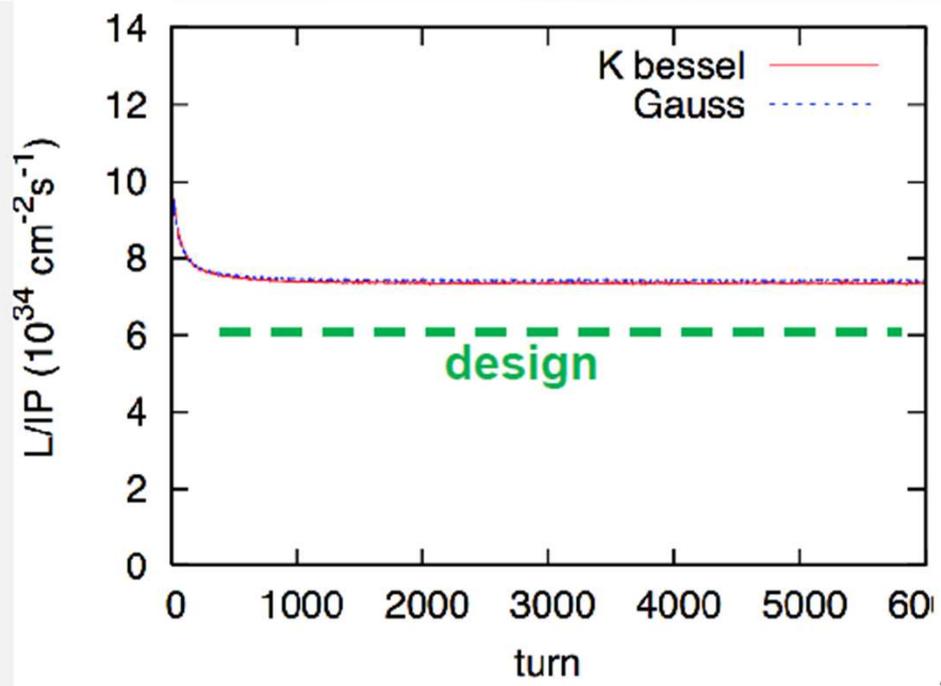
YFARS/AN S CERN

Beam-beam simulations



8th FCC-ee Physics Workshop - Paris - J. Wenninger

26/10/2014



BBSS strong-strong simulation with beamstrahlung

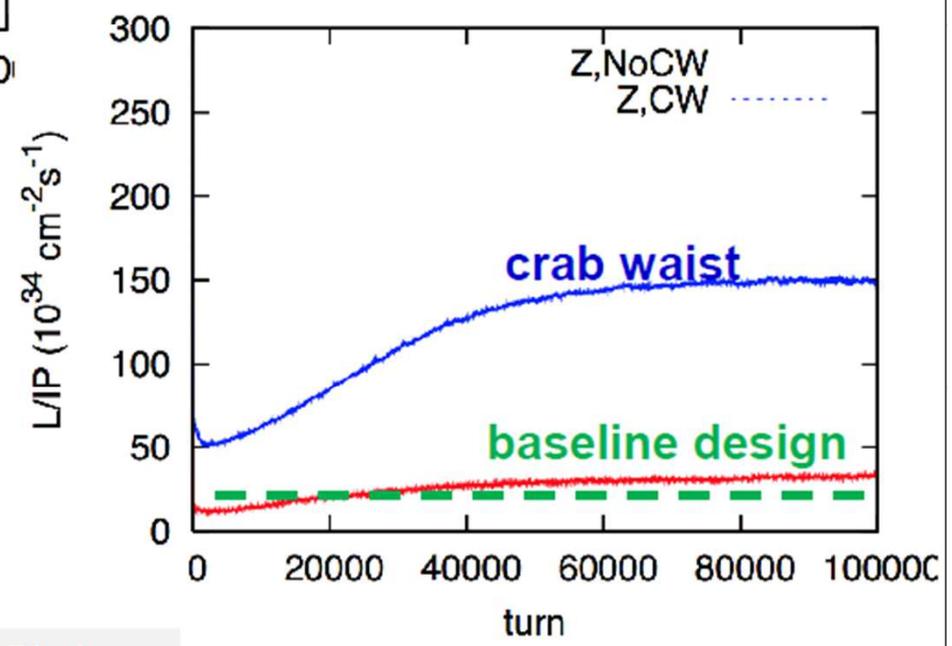
FCC-ee at 120 GeV:

$L \approx 7.5 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ per IP

FCC-ee in crab-waist mode at the Z pole (45.5 GeV):

$L \approx 1.5 \times 10^{36} \text{ cm}^{-2} \text{ s}^{-1}$ per IP

Tracking confirms assumptions!



K. Ohmi, A. Bogomyagkov, E. Levichev, P. Piminov

4/30/2015

Key Parameters FCC-hh

Parameter	FCC-hh	LHC
Energy [TeV]	100 c.m.	14 c.m.
Dipole field [T]	16	8.33
# IP	2 main, +2	4
Luminosity/IP _{main} [cm ⁻² s ⁻¹]	5 - 25 x 10 ³⁴	1 x 10 ³⁴
Stored energy/beam [GJ]	8.4	0.39
Synchrotron rad. [W/m/aperture]	28.4	0.17
Bunch spacing [ns]	25 (5)	25





FCC-hh: some design challenges

- **Stored beam energy: 8 GJ/beam (0.4 GJ LHC) = 16 GJ total**
➔ equivalent to an Airbus A380 (560 t) at full speed (850 km/h)

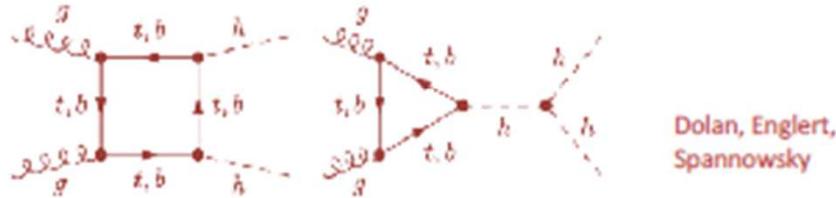


- **Collimation, beam loss control, radiation effects: very important**
- **Injection/dumping/beam transfer: very critical operations**
- **Magnet/machine protection: to be considered from early phase**

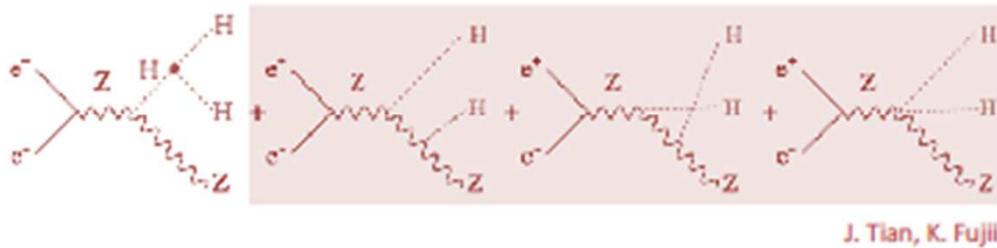


H³ @ TLEP

- At LHC (Requires $E_{CM} > 2 m_h$):



- At ILC (Requires $E_{CM} > 2 m_h + m_Z$):



- At TLEP 240 GeV: **M. McCullough '14**

$$\sigma_{Zh} = \left| \begin{array}{c} e \\ \downarrow \\ Z \\ \uparrow \\ e \end{array} \right|^2 + 2 \operatorname{Re} \left[\begin{array}{c} e \\ \downarrow \\ Z \\ \uparrow \\ h \end{array} \cdot \left(\begin{array}{c} e^+ \\ \downarrow \\ Z \\ \uparrow \\ h \end{array} + \begin{array}{c} e^+ \\ \downarrow \\ Z \\ \uparrow \\ h \end{array} \right) \right]$$

$$\delta_{\sigma}^{240} = 100 (2\delta_Z + 0.014\delta_h) \%$$

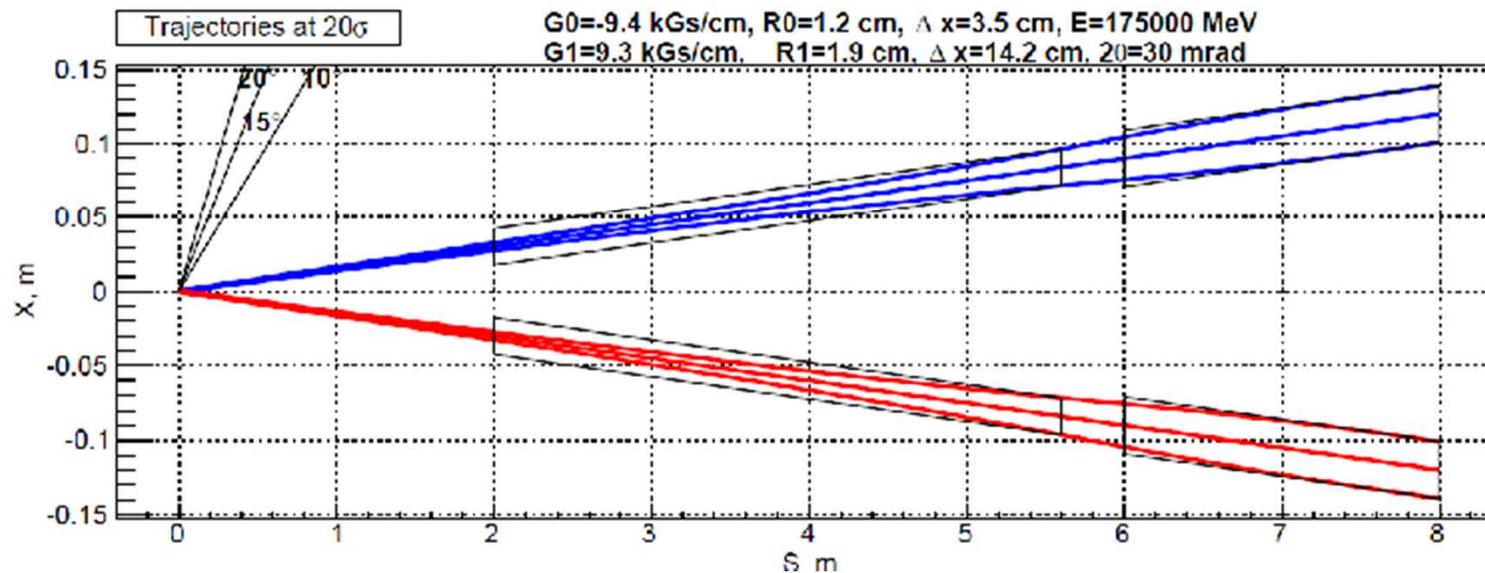
tiny effect but visible thanks to the extraordinary TLEP sensitivity on Zh (0.05%)

Table 1-16. Uncertainties on coupling scaling factors as determined in a completely model-independent fit for different e^+e^- facilities. Precisions reported in a given column include in the fit all measurements at lower energies at the same facility, and note that the model independence requires the measurement of the recoil HZ process at lower energies. [†]ILC luminosity upgrade assumes an extended running period on top of the low luminosity program and cannot be directly compared to TLEP and CLIC numbers without accounting for the additional running period. ILC numbers include a 0.5% theory uncertainty. For invisible decays of the Higgs, the number quoted is the 95% confidence upper limit on the branching ratio.

Facility		ILC		ILC(LumiUp)		TLEP (4 IP)		CLIC	
\sqrt{s} (GeV)	250	500	1000	250/500/1000	240	350	350	1400	3000
$\int \mathcal{L} dt$ (fb ⁻¹)	250	+500	+1000	1150+1600+2500 [†]	10000	+2600	500	+1500	+2000
$P(e^-, e^+)$	(-0.8, +0.3)	(-0.8, +0.3)	(-0.8, +0.2)	(same)	(0, 0)	(0, 0)	(-0.8, 0)	(-0.8, 0)	(-0.8, 0)
Γ_H	12%	5.0%	4.6%	2.5%	1.9%	1.0%	9.2%	8.5%	8.4%
κ_γ	18%	8.4%	4.0%	2.4%	1.7%	1.5%	—	5.9%	<5.9%
κ_g	6.4%	2.3%	1.6%	0.9%	1.1%	0.8%	4.1%	2.3%	2.2%
κ_W	4.9%	1.2%	1.2%	0.6%	0.85%	0.19%	2.6%	2.1%	2.1%
κ_Z	1.3%	1.0%	1.0%	0.5%	0.16%	0.15%	2.1%	2.1%	2.1%
κ_μ	91%	91%	16%	10%	6.4%	6.2%	—	11%	5.6%
κ_τ	5.8%	2.4%	1.8%	1.0%	0.94%	0.54%	4.0%	2.5%	<2.5%
κ_c	6.8%	2.8%	1.8%	1.1%	1.0%	0.71%	3.8%	2.4%	2.2%
κ_b	5.3%	1.7%	1.3%	0.8%	0.88%	0.42%	2.8%	2.2%	2.1%
κ_t	—	14%	3.2%	2.0%	—	13%	—	4.5%	<4.5%
BR_{inv}	0.9%	< 0.9%	< 0.9%	0.4%	0.19%	< 0.19%			

**the
10B\$ ILC**

Trajectories $L^*=2\text{m}$



Mogens Dam

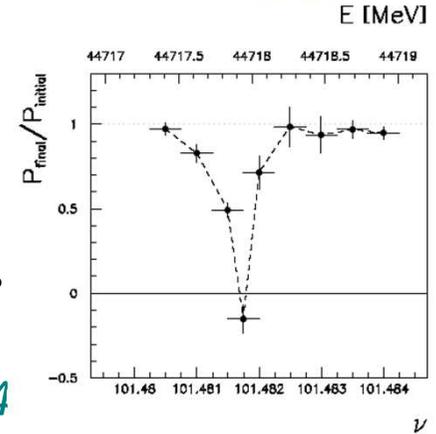
**example of challenge: crab crossing to
increase further luminosity? (Novosibirsk)
emittance and polarization compensation, etc**

Beam polarization and E-calibration @ TLEP



Precise meast of E_{beam} by resonant depolarization
~100 keV each time the meast is made

At LEP transverse polarization was achieved routinely at Z peak.
instrumental in 10^{-3} measurement of the Z width in 1993
led to prediction of top quark mass (179+- 20 GeV) in March 1994



Polarization in collisions was observed (*40% at BBTS = 0.04*)

At LEP beam energy spread destroyed polarization above 60 GeV
 $\sigma_E \propto E^2/\sqrt{\rho} \rightarrow$ At TLEP transverse polarization up to at least 80 GeV
to go to higher energies requires spin rotators and siberian snake

TLEP: use 'single' bunches to measure the beam energy continuously
no interpolation errors due to tides, ground motion or trains etc...

<< 100 keV beam energy calibration around Z peak and W pair threshold.

$\Delta m_Z \sim 0.1 \text{ MeV}$, $\Delta \Gamma_Z \sim 0.1 \text{ MeV}$, $\Delta m_W \sim 0.5 \text{ MeV}$

Asymptotic safety of gravity and the Higgs boson mass

Mikhail Shaposhnikov

Institut de Théorie des Phénomènes Physiques, École Polytechnique Fédérale de Lausanne, CH-1015 Lausanne, Switzerland

Christof Wetterich

Institut für Theoretische Physik, Universität Heidelberg, Philosophenweg 16, D-69120 Heidelberg, Germany

12 January 2010

Abstract

There are indications that gravity is asymptotically safe. The Standard Model (SM) plus gravity could be valid up to arbitrarily high energies. Supposing that this is indeed the case and assuming that there are no intermediate energy scales between the Fermi and Planck scales we address the question of whether the mass of the Higgs boson m_H can be predicted. For a positive gravity induced anomalous dimension $A_\lambda > 0$ the running of the quartic scalar self interaction λ at scales beyond the Planck mass is determined by a fixed point at zero. This results in $m_H = m_{\min} = 126$ GeV, with only a few GeV uncertainty. This prediction is independent of the details of the short distance running and holds for a wide class of extensions of the SM as well. For $A_\lambda < 0$ one finds m_H in the interval $m_{\min} < m_H < m_{\max} \simeq 174$ GeV, now sensitive to A_λ and other properties of the short distance running. The case $A_\lambda > 0$ is favored by explicit computations existing in the literature.

in 2010 Shaposhnikov and Wetterich predict $m_H=126$ GeV

if there is no intermediate energy scale between the Fermi and Planck scales...



FCC Work and Organisation (i)

Work/meeting structures established based on INDICO, see:

- **FCC Study:** <https://indico.cern.ch/category/5153/>
- <http://cern.ch/FCC-ee> (more developed, for FCC-ee)

In particular:

- **FCC-hh Hadron Collider Physics and Experiments VIDYO meetings**
 - <https://indico.cern.ch/category/5258/>
 - **Contacts:** michelangelo.mangano@cern.ch,
fabiola.gianotti@cern.ch, austin.ball@cern.ch
- **FCC-ee Lepton Collider (TLEP) Physics and Experiments VIDYO meetings**
 - <https://indico.cern.ch/category/5259/>
 - **Contacts:** alain.blondel@cern.ch, patrick.janot@cern.ch





FCC Work and Organisation (ii)

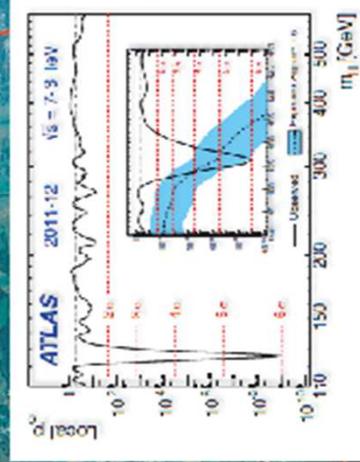
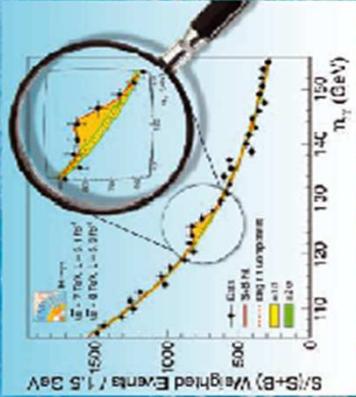
- **FCC-hh Hadron Collider VIDYO meetings**
 - <https://indico.cern.ch/category/5263/>
 - **Contacts:** daniel.schulte@cern.ch
- **FCC-hadron injector meetings**
 - <https://indico.cern.ch/category/5262/>
 - **Contacts:** brennan.goddard@cern.ch
- **FCC-ee (TLEP) Lepton Collider VIDYO meetings**
 - <https://indico.cern.ch/category/5264/>
 - **Contacts:** jorg.wenninger@cern.ch,
- **FCC infrastructure meetings**
 - <https://indico.cern.ch/category/5253/>
 - **Contacts:** philippe.lebrun@cern.ch, peter.sollander@cern.ch





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In praise of charter schools

Britain's banking scandal spreads

Volkswagen overtakes the rest

A power struggle at the Vatican

When Lonesome George met Nora

The
Economist

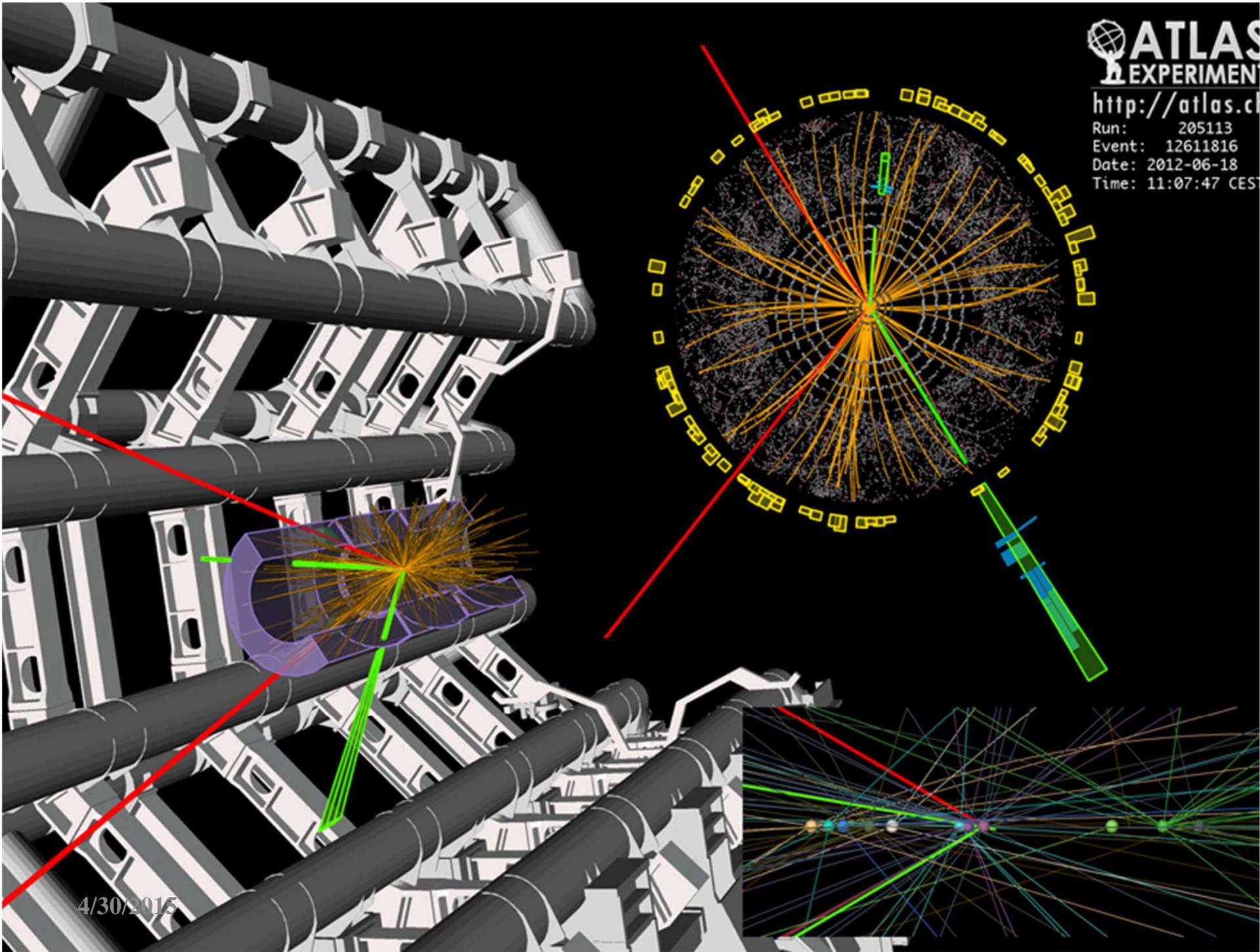
2012 776-1176 2012

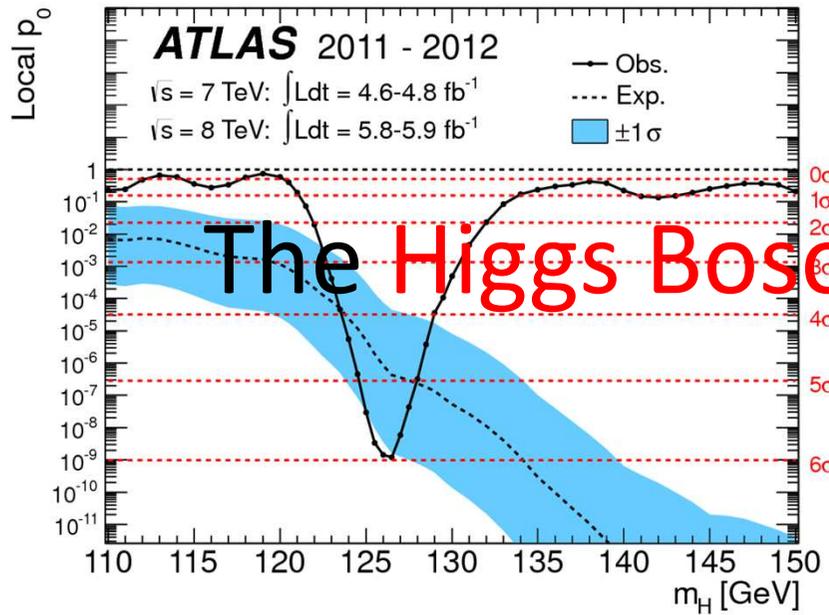
[Economist.com](http://economist.com)

A giant leap for science

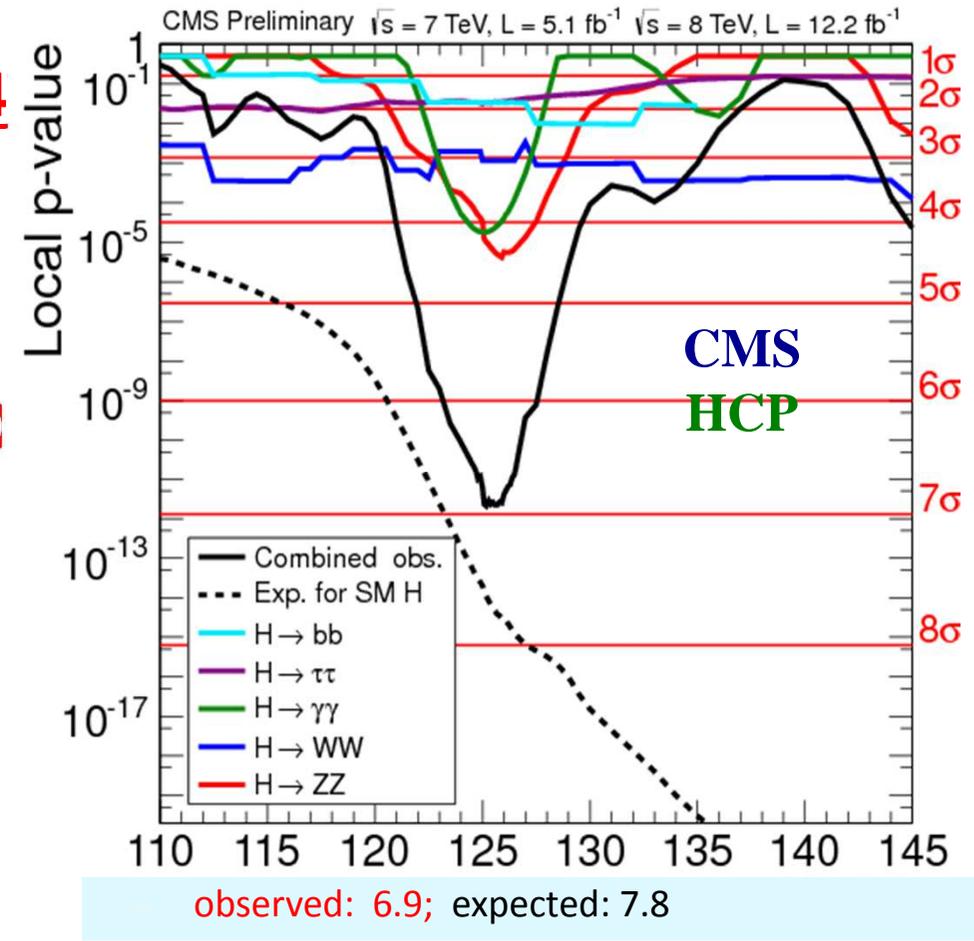


Finding the
Higgs boson





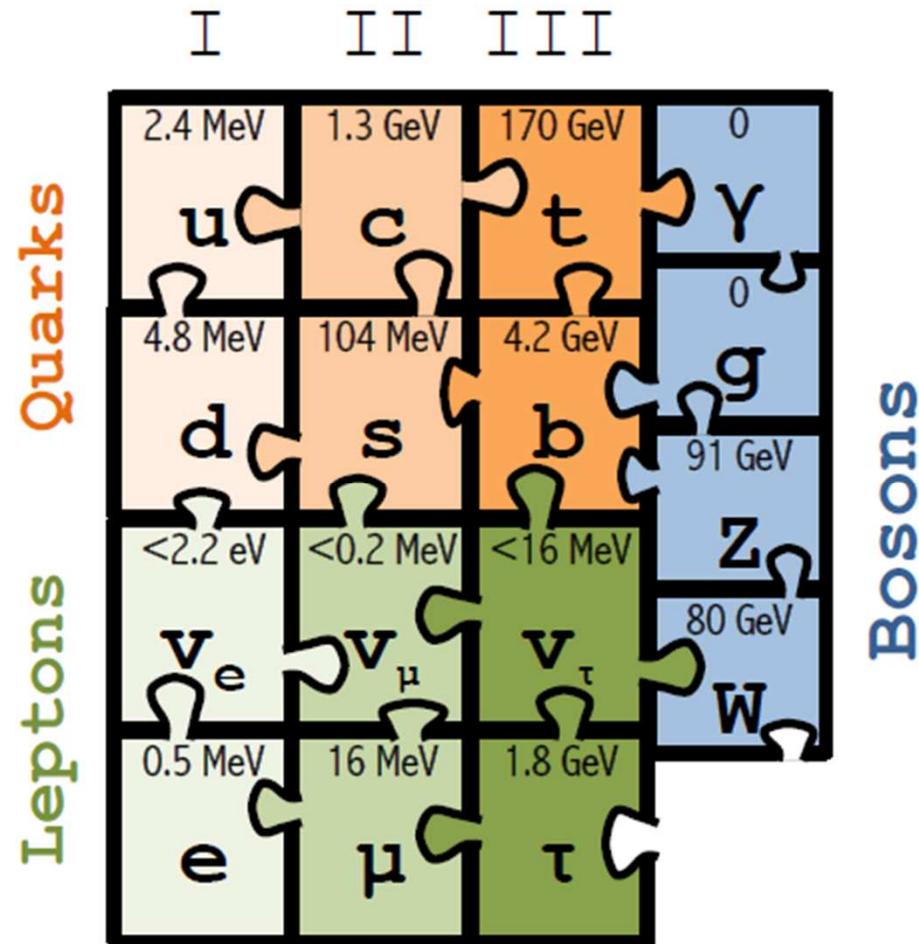
4



Discovered Higgs-like Boson: Clear mass peak in $\gamma\gamma$ and ZZ^* → 4l



1994-1999: top mass predicted (LEP, mostly Z mass&width)
top quark discovered (Tevatron)
t'Hooft and Veltman get Nobel Prize



(c) Sfyrla

4/30/2015

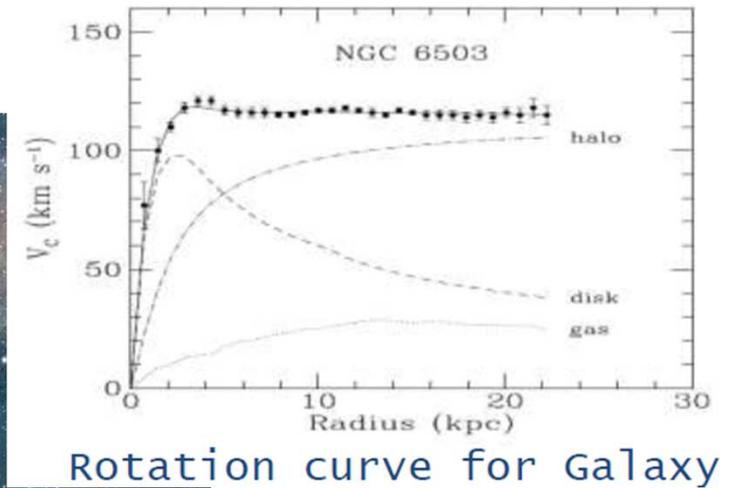
Alain Blondel FCC Future Circular Colliders



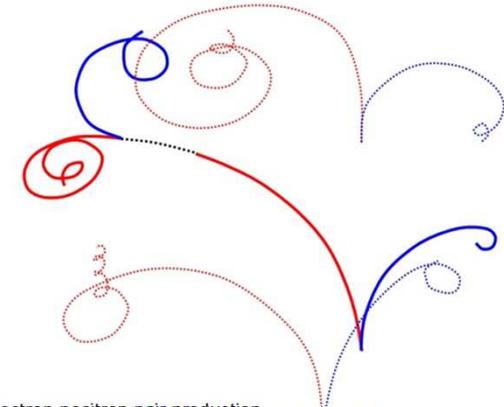
We cannot explain:

Dark matter

Standard Model particles constitute only 5% of the energy in the Universe

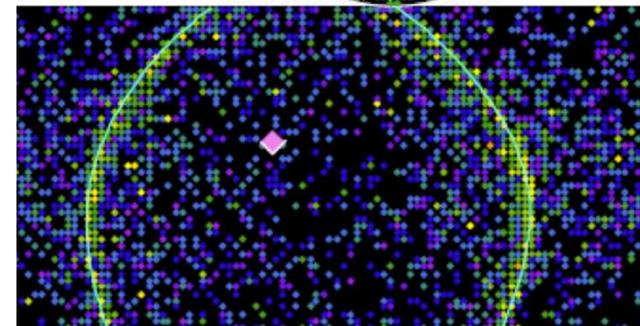


Where is antimatter gone?



What makes neutrino masses?

- Not a unique solution in the SM --
- Dirac masses (why so small?)
- Majorana masses (why not Dirac?)
- Both (the preferred scenarios, see-saw...)
- heavy right handed neutrinos?



we cannot explain:

charge of proton = - charge of electron

$$|q_p + q_e|/e < 1 \times 10^{-21}$$

we have no explanation for this, except ...

that it is necessary for the stability of

1. the universe

2. the Standard Model calculations





PARAMETERS FOR CRAB WAIST OPERATION

	Z	W	H	tt
Energy [GeV]	45	80	120	175
Perimeter [km]	100			
Crossing angle [mrad]	30			
Particles per bunch [10^{11}]	1	4	4.7	4
Number of bunches	29791	739	127	33
Energy spread [10^{-3}]	1.1	2.1	2.4	2.6
Emittance hor. [nm]	0.14	0.44	1	2.1
Emittance ver. [pm]	1	2	2	4.3
β_x^*/β_y^* [m]	0.5 / 0.001			
Luminosity / IP [$10^{34} \text{ cm}^{-2} \text{ s}^{-1}$]	Nominal : 212	28 36	12 9	6.0 1.3
Energy loss / turn [GeV]	0.03	0.3	1.7	7.7

Important scope for improvement in luminosity.





Luminosity optimisation

Ideal situation is that beam lifetime is driven by particle-particle interactions
-- dominated by radiative Bhabha scattering $e^+e^- \rightarrow e^+e^-\gamma$ (typically 150 mb)
with $e^{+/-}$ out of energy acceptance (improved with larger acceptance)

At high luminosity considered in FCC-ee, Beamstrahlung (particle-opp. beam interaction) becomes important.

- requires very flat beams and +/- 2% energy acceptance
- reduces beam lifetime
- increases energy spread and bunch length

This is the case in FCC-tt

At lower energy the beams are blowing each other (beam-beam interaction)

- this can be fought with 'crab waist' crossing

This is the case at all lower energies operating points

Numbers in main parameter list include beamstrahlung treatment, but have not considered crab waist operation.





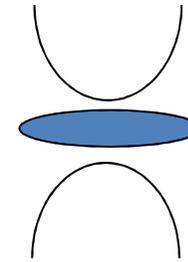
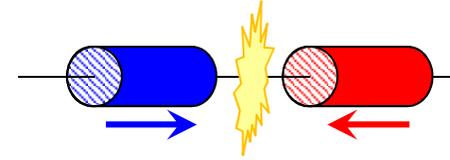
Luminosity

$$efkN = \text{beam current} \propto \frac{1}{E^4}$$

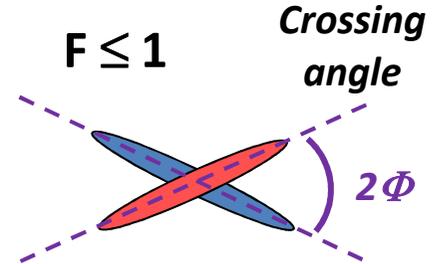
$$L = \frac{fkN^2}{4\pi\sigma_x\sigma_y} FH$$

$$\xi_y \propto \frac{\beta_y^* N}{E\sigma_x\sigma_y} \leq \xi_y^{\max}(E) \quad \text{Beam-beam parameter}$$

$$L \propto \frac{P_{SR}}{E^3} \frac{\xi_y}{\beta_y^*}$$



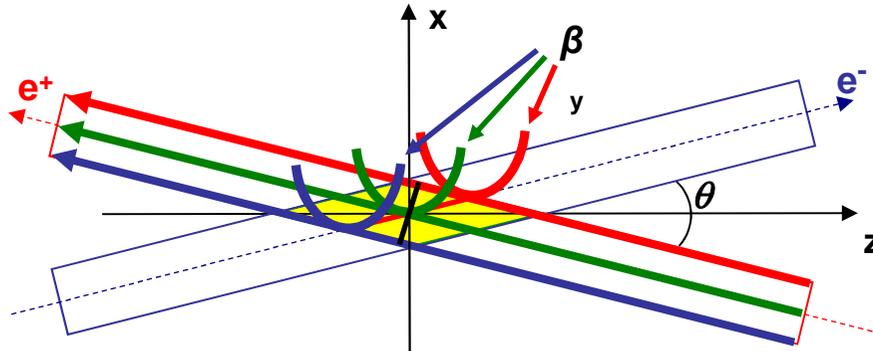
$H \leq 1$
Hour-glass



$F \leq 1$
Crossing angle

- σ = beam size
- k = no. bunches
- f = rev. frequency
- N = bunch population
- P_{SR} = synch. rad. power
- β^* = betatron fct at IP
(beam envelope)





$$\phi = \frac{\sigma_z}{\sigma_x} \operatorname{tg} \left(\frac{\theta}{2} \right) \quad \text{– Piwinski angle}$$

- 1) Large Piwinski angle: $\phi \gg 1$
- 2) β_y approx. equals to overlapping area: $\beta_y \sim \sigma_z / \phi$
- 3) Crab Waist: minimum of β_y along the axis of the opposite beam

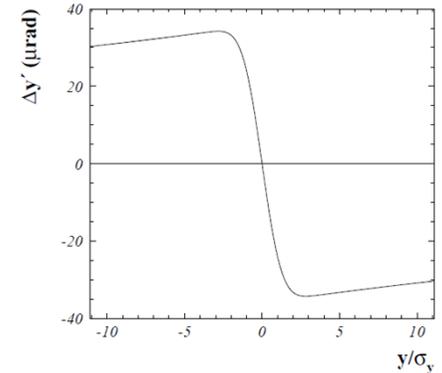
Advantages:

- ✓ Impact of hour-glass is small and does not depend on bunch lengthening
- ✓ Suppression of betatron coupling resonances allows to achieve $\xi_y \sim 0.2$
- ✓ As a result, luminosity can be significantly increased especially at Z, otherwise $\xi_y \sim 0.03$



Beam-beam parameter

- The beam-beam parameter ξ measures the strength of the field sensed by the particles due to the counter-rotating bunch.
- Beam-beam parameter limits are empirically scaled from LEP data (also 4 IPs).

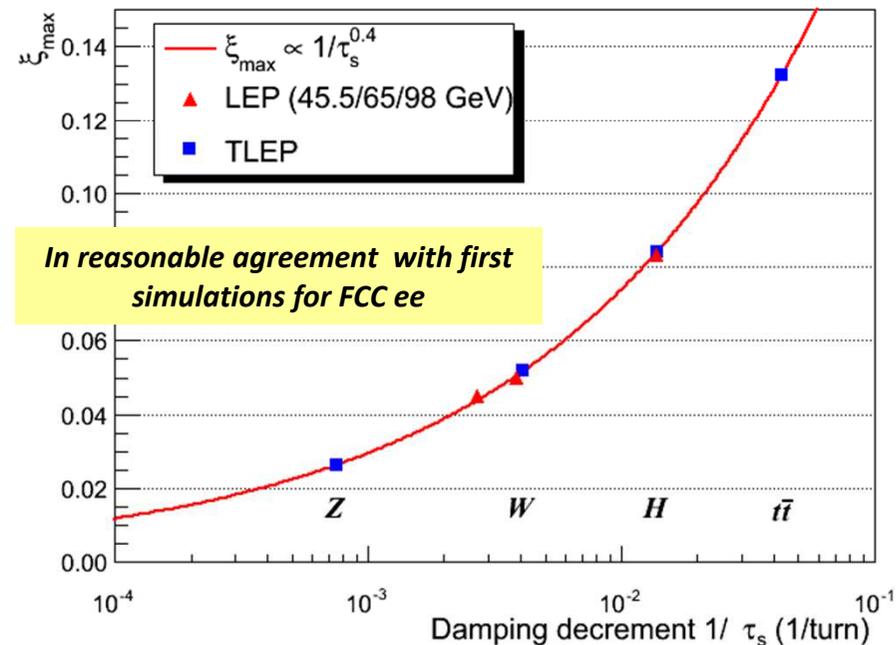


$$\xi_y \propto \frac{\beta_y^* N}{E \sigma_x \sigma_y} \leq \xi_y^{\max}(E)$$

$$\xi_y^{\max}(E) \propto \frac{1}{\tau_s^{0.4}} \propto E^{1.2}$$



$$L \propto \frac{P_{SR}}{E^{1.8}} \frac{1}{\beta_y^*}$$

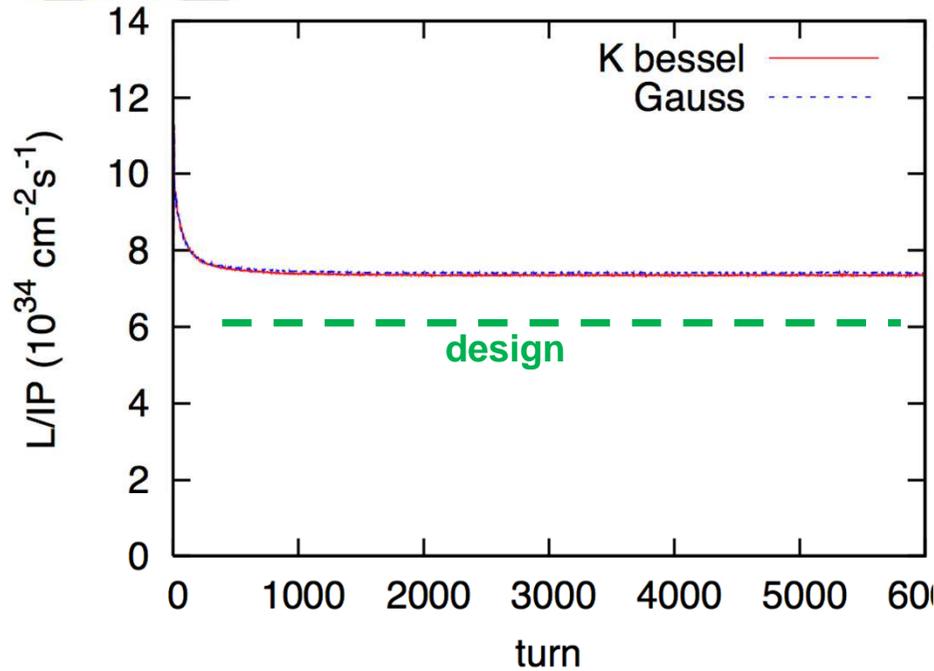


ξ_y and L may be raised significantly (x 4) with Crab-Waist schemes !





Beam-beam simulations



BBSS strong-strong simulation with beamstrahlung

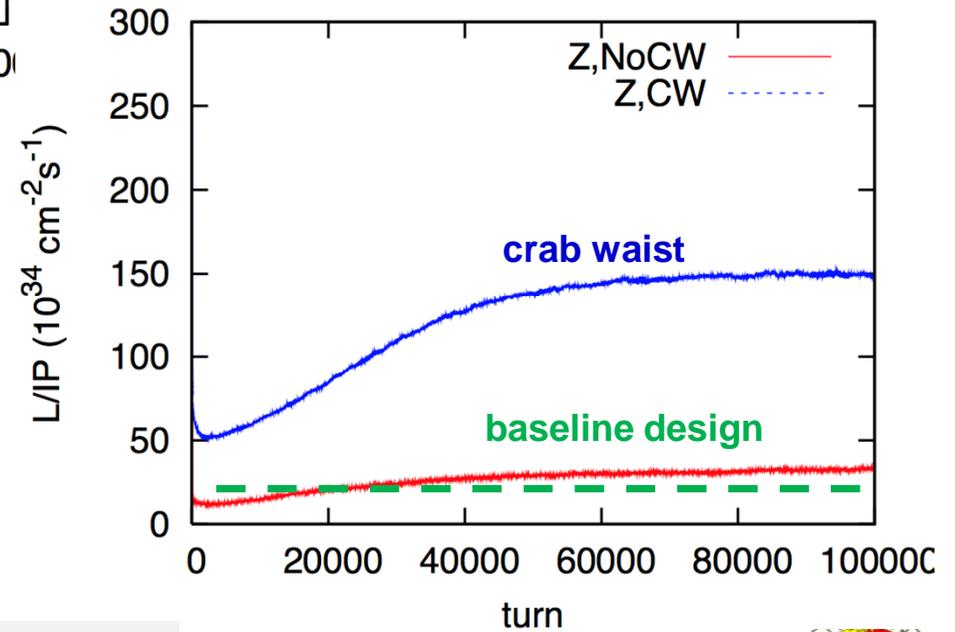
FCC-ee at 120 GeV:

$L \approx 7.5 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ per IP

FCC-ee in crab-waist mode at the Z pole (45.5 GeV):

$L \approx 1.5 \times 10^{36} \text{ cm}^{-2} \text{ s}^{-1}$ per IP

Tracking confirms assumptions!



30/04/2015

K. Ohmi, A. Bogomyagkov, E. Levichev, P. Piminov University of Krakow



Beamstrahlung

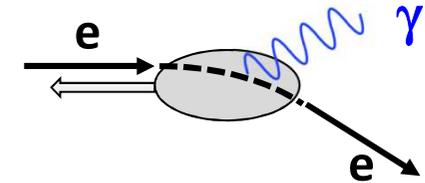
- Hard photon emission at the IPs, '*Beamstrahlung*', can become a lifetime / performance limit for large bunch populations (N), small hor. beam size (σ_x) and short bunches (σ_s).

$$\tau_{bs} \propto \frac{\rho^{3/2} \sqrt{\eta}}{\sigma_s} \exp(A\eta\rho)$$

η : ring energy acceptance

$$\frac{1}{\rho} \approx \frac{Nr_e}{\gamma\sigma_x\sigma_s}$$

$$L = \frac{fkN^2}{4\pi\sigma_x\sigma_y} FH$$



ρ : mean bending radius at the IP (in the field of the opposing bunch)

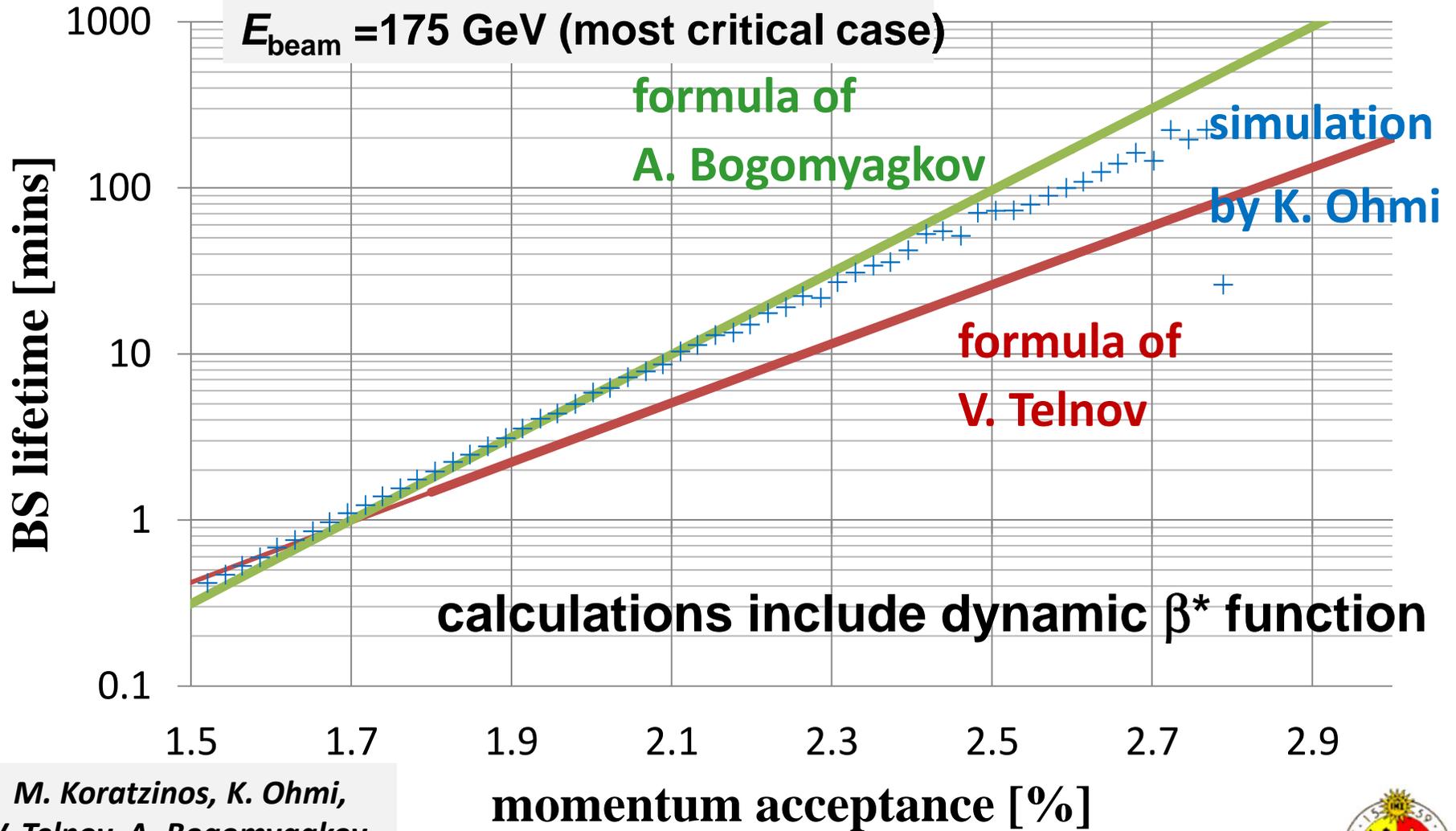
Lifetime expression by V. Telnov

- To ensure an acceptable lifetime, $\rho \times \eta$ must be sufficiently large.
 - Flat beams : large σ_x and small σ_y !
 - Bunch length !
 - Large momentum acceptance of the lattice: 1.5 – 2% required.
 - LEP had < 1% acceptance, SuperKEKB ~ 1-1.5%.



Beamstrahlung lifetime

Reasonable agreement between tracking and analytical estimates.



8th FCC-ee Physics Workshop - Paris - J. Wenninger

30/04/2015

M. Koratzinos, K. Ohmi,
V. Telnov, A. Bogomyagkov,
E. Levichev, D. Shatilov

Blondel FCC-ee Epiphany Conference Krakow





Emittances

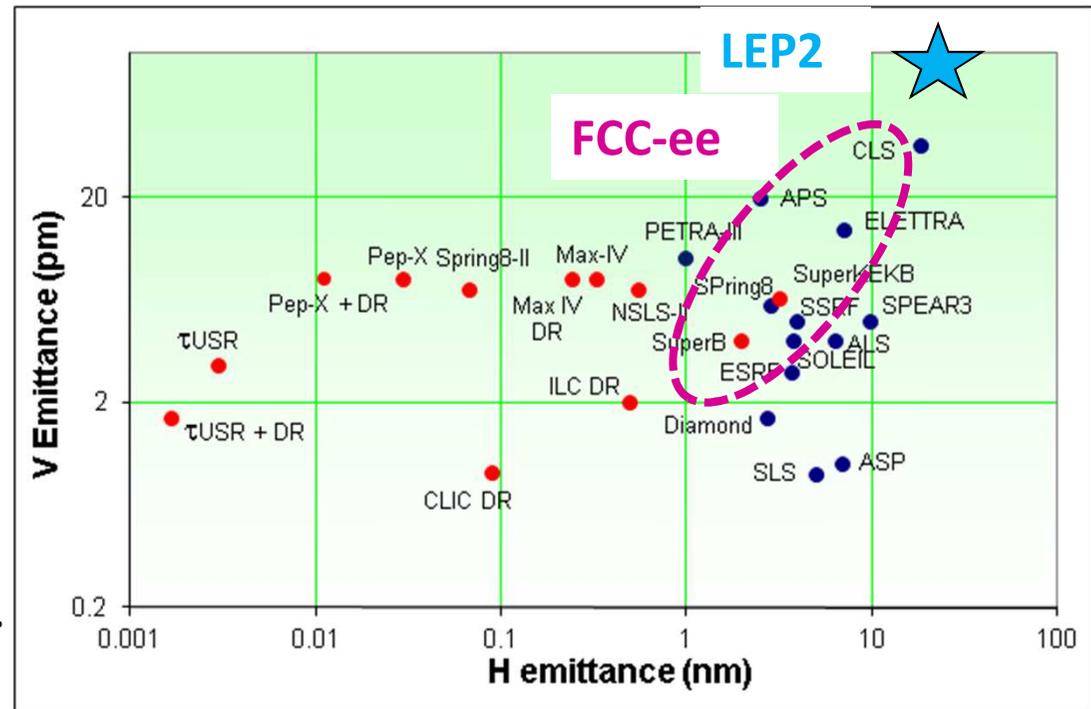
- FCC-ee is a very large machine, scaling of achievable emittances (mainly vertical) is not straightforward.
 - Coupling, spurious vertical dispersion.
- Low emittances tend to be more difficult to achieve in colliders as compared to light sources or damping rings – beam-beam !

□ FCC-ee parameters:

- $\epsilon_y/\epsilon_x \geq 0.001$,
- $\epsilon_y \geq \approx 2 \text{ pm}$

with a ring ~50-100 larger than a typical light source.

- Very challenging target for a ring of that size!
- LEP2 achieved routinely 0.004 beam corrections are much better now.

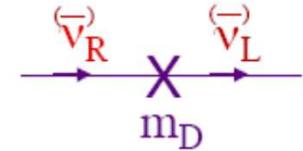


R. Bartolini, DIAMOND



Adding masses to the Standard model neutrino 'simply' by adding a Dirac mass term (Yukawa coupling)

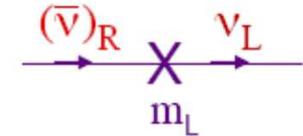
$$m_D \nu_L \bar{\nu}_R \quad m_D \bar{\nu}_L \nu_R$$



implies adding a right-handed neutrino (new particle)

No SM symmetry prevents adding then a term like

$$m_M \bar{\nu}_R^c \nu_R$$



**and this simply means that a neutrino turns into a antineutrino
(the charge conjugate of a right handed antineutrino is a left handed neutrino!)**

It is perfectly conceivable ('natural'?) that both terms are present → 'see-saw'

